

Recent results of analysis of runaway electron experiments at ASDEX Upgrade

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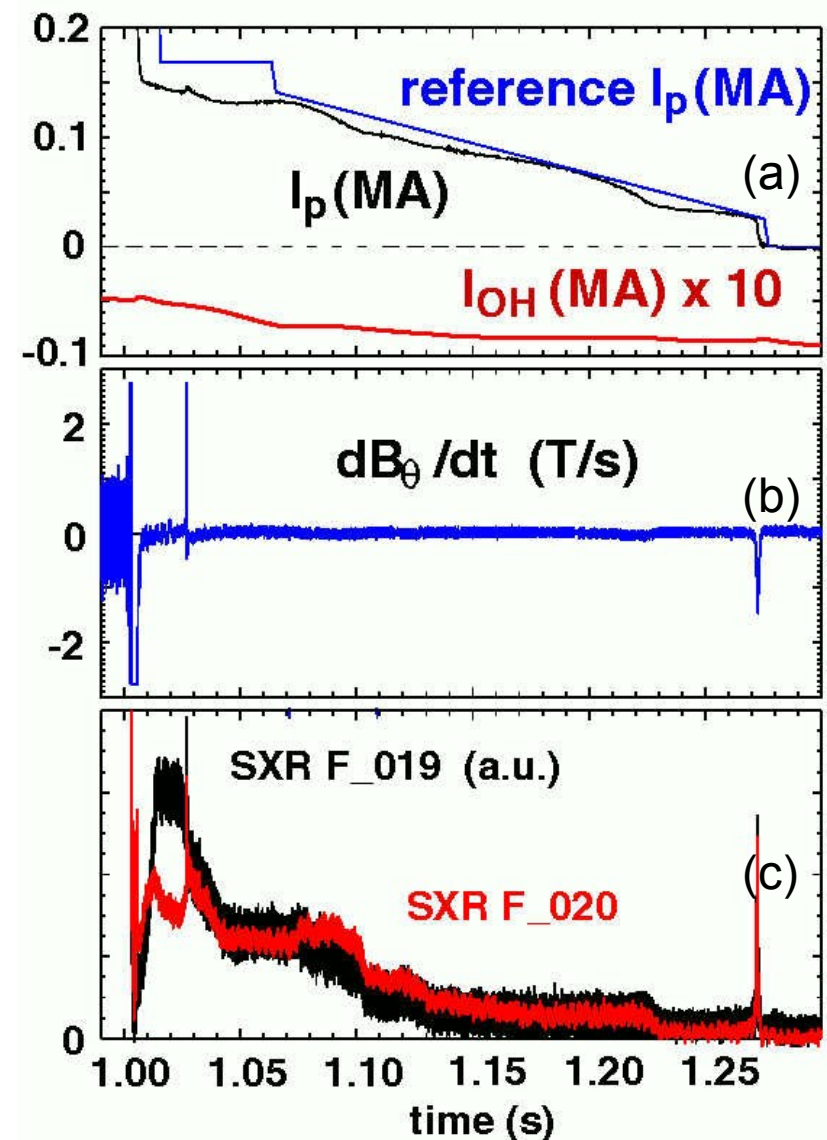
Introduction

- ASDEX Upgrade (AUG): no runaway electron (RE) experiments in last year (water leak in Dec. 2017, restart in July 2018) → further analysis of past exp.s
- Well diagnosed experiments on existing tokamaks are indispensable for validation of RE generation and suppression models, which can then be used for simulation of reactor scenarios
- Quest for RE loss mechanisms during both thermal quench (TQ) and RE beam lifetime, and ways to enhance losses are being pursued
- This talk is about
 - how to pin down origin of RE current (loss or generation) scatter in similar discharges?
 - why 2nd high-Z gas injection is effective for RE suppression in medium-sized tokamaks but not in JET?

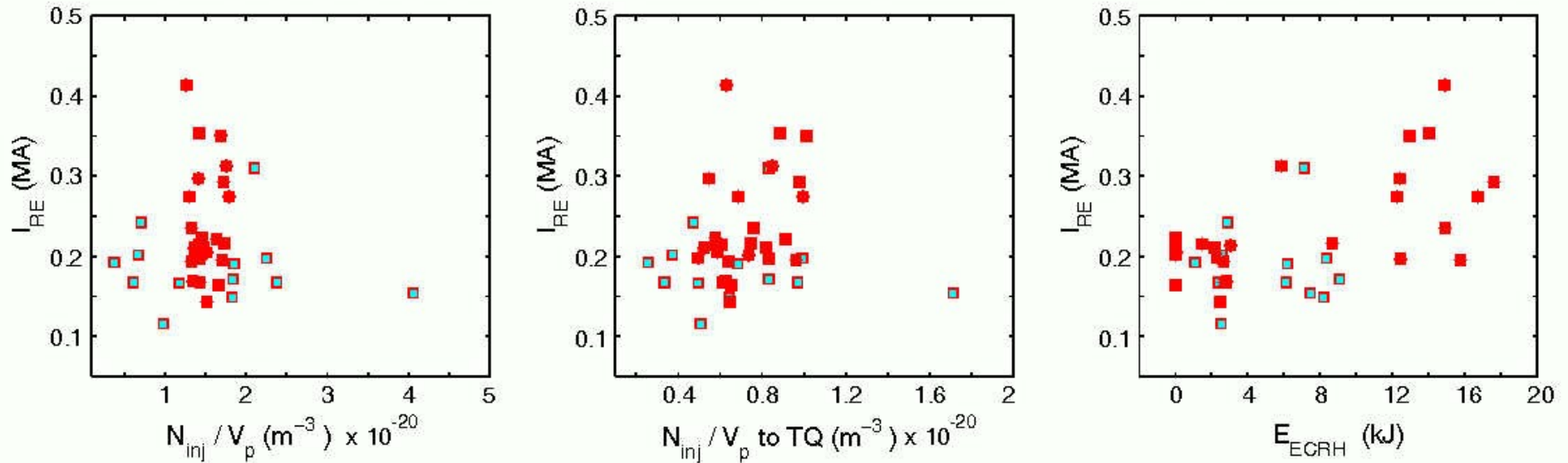
(both questions will remain questions)

AUG scenario for RE experiments

- First exp.s in 2014
(~ 80 discharges by now)
- otherwise uninteresting target plasma: circular, low n_e , $B_t \sim 2.5$ T, $P_{\text{ECRH}} \sim 2$ MW, $I_p = 0.8$ MA
- RE beam ($I_{\text{RE}} < 400$ kA for < 500 ms) is reliably generated with argon puff
- toroidal current is controlled
- plasma has been mostly vertically stable and w/o MHD activity
- use of RMPs not discussed



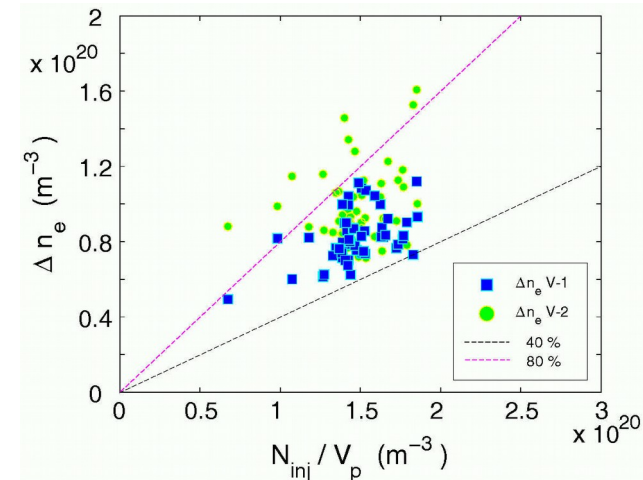
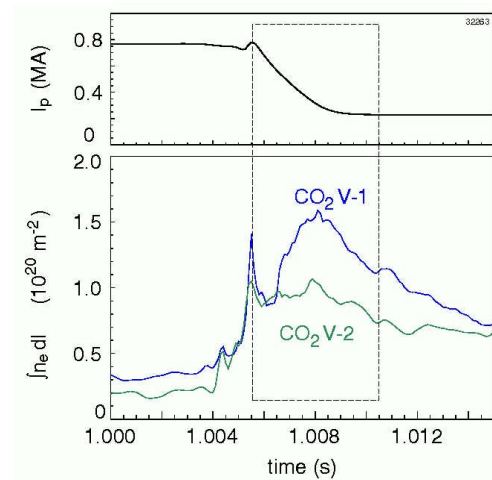
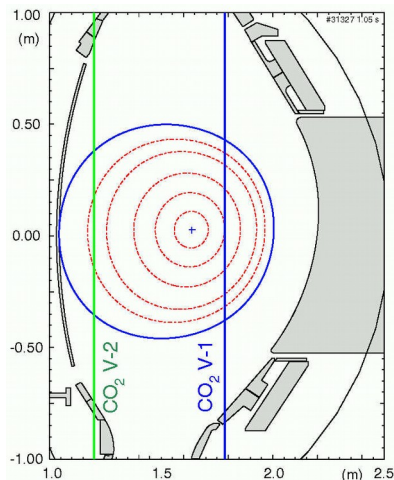
Origin of RE current “scatter”?



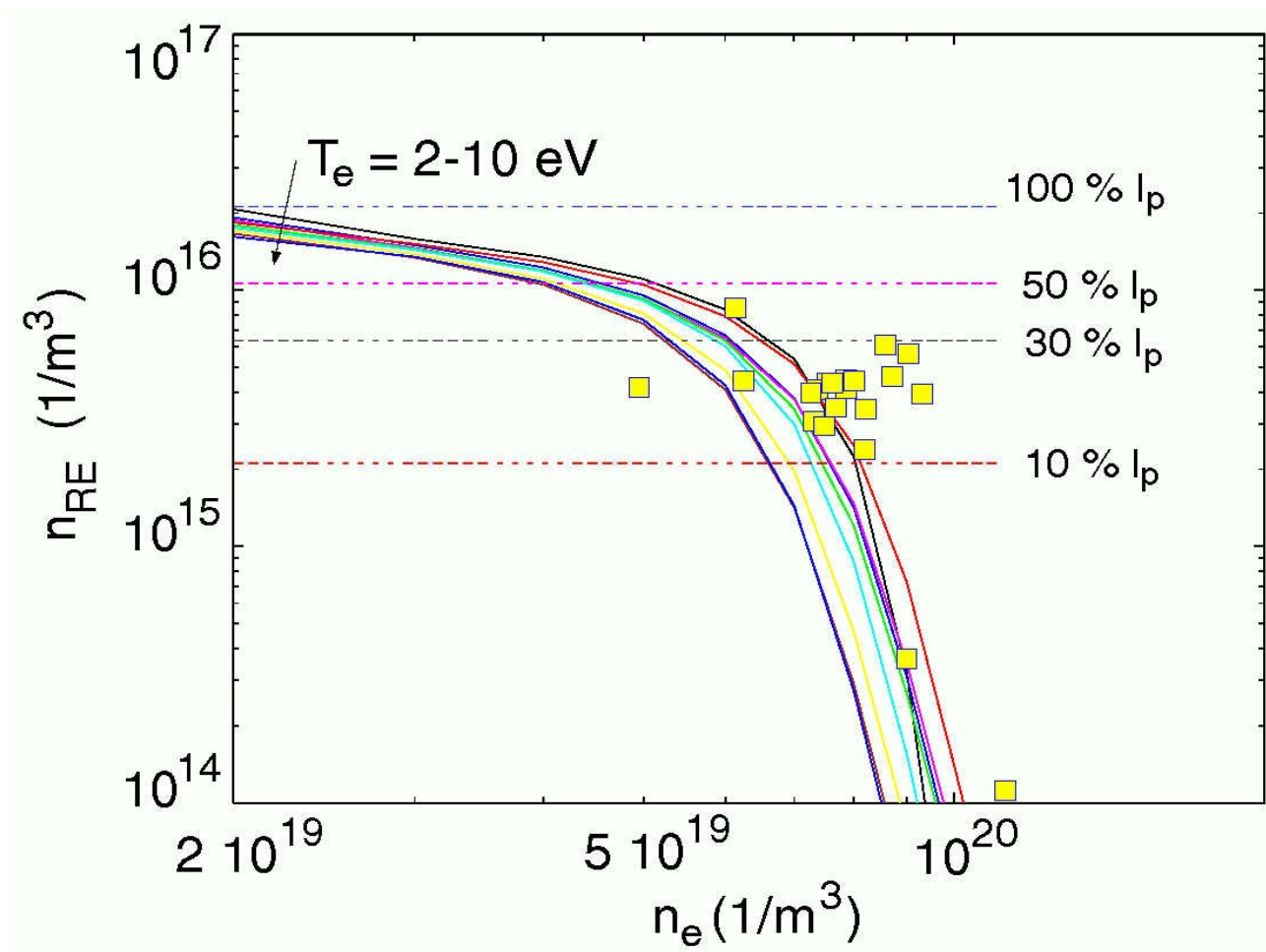
- Set of similar discharges, only 1st MGI, argon amount (N_{inj}) varied
- I_{RE} versus (left) N_{inj} per plasma volume (V_p), (center) N_{inj}/V_p during pre-TQ and (right) ECRH energy injected (and absorbed?) after beginning of pre-TQ
- No other parameter dependence found (pre-TQ time, ΔI_i , ΔI_p , equilibrium ...)

Density measurements

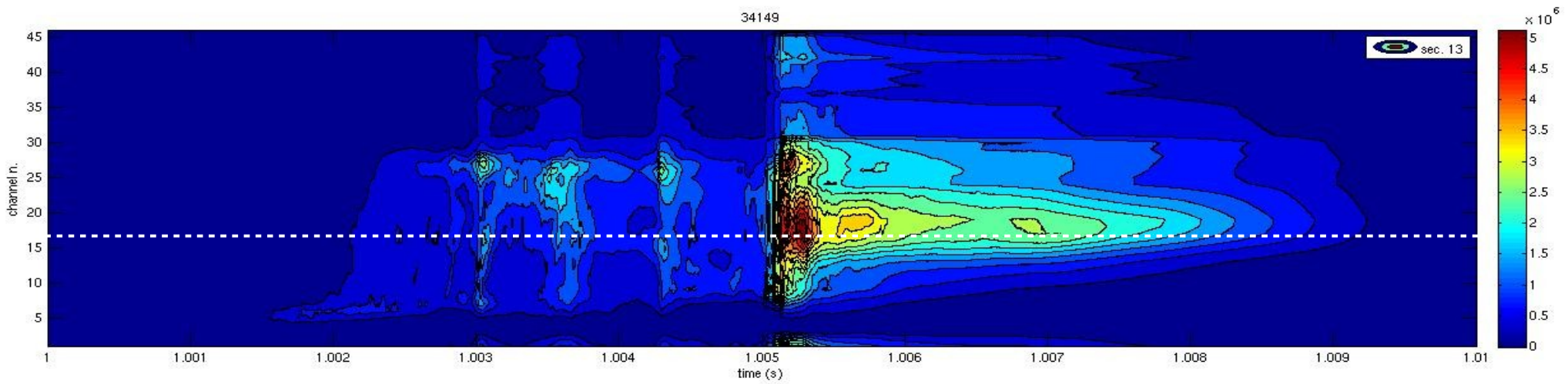
- RE generation/suppression depend strongly on local densities (n_e , n_I , Z , Z_{eff})
- no detailed density profiles during fast CQ (~ 2 ms)
- n_e from CO₂ interferometer V-1 chord on AUG – used as proxy for n_e seen by REs



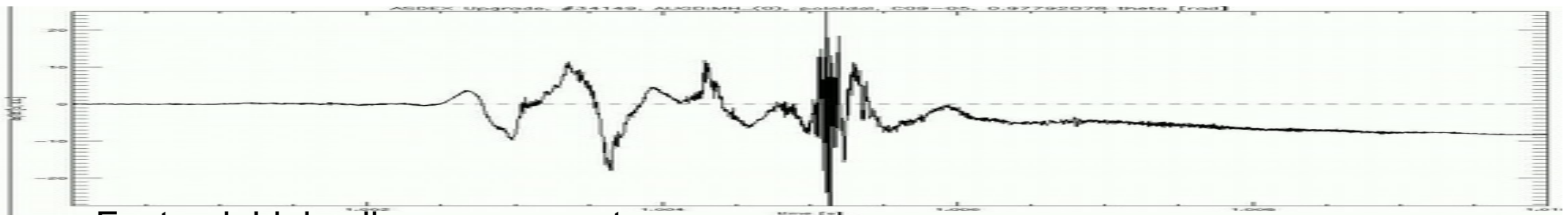
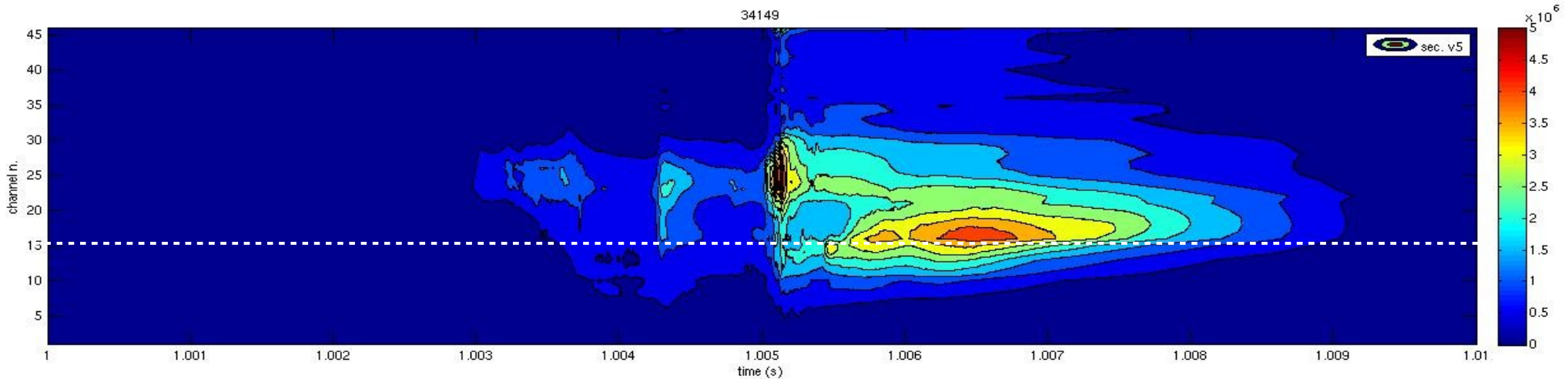
I_{RE} versus n_e and T_e



RE current calculated with 0-D model for several n_e , argon density and T_e , values (lines with different colours) and experimental measurements (yellow squares)



Plasma radiation, top-down view, sectors $\Delta\phi=\pi$

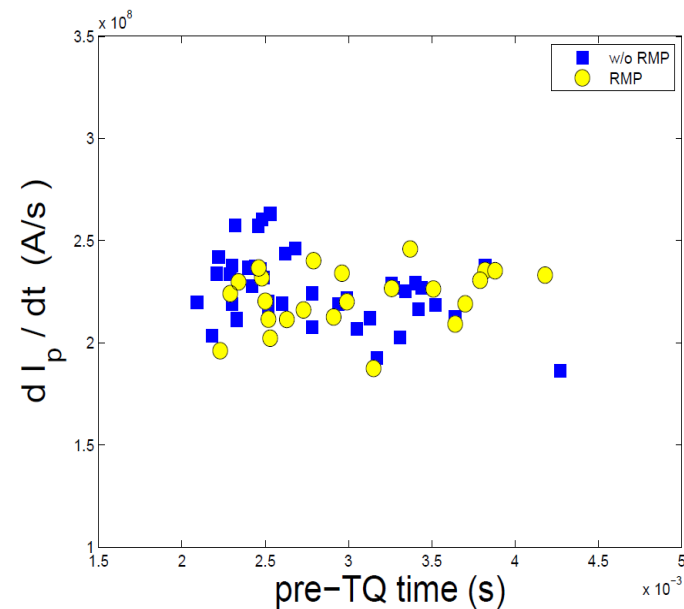
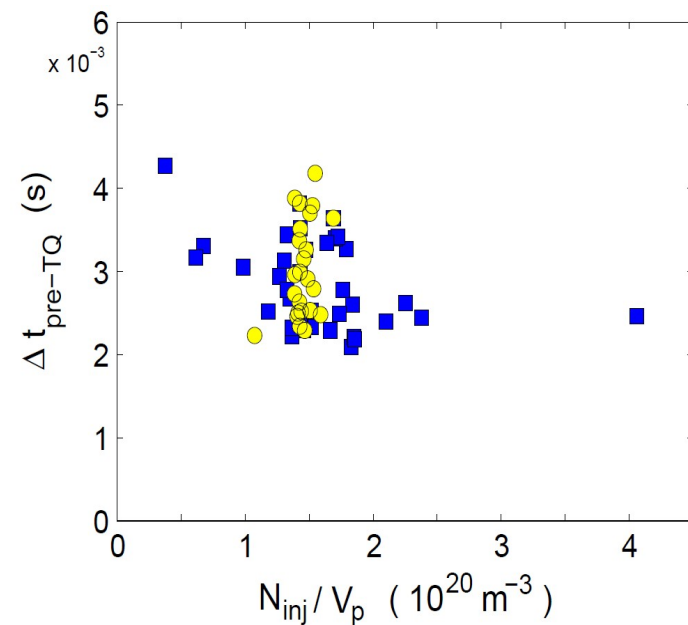
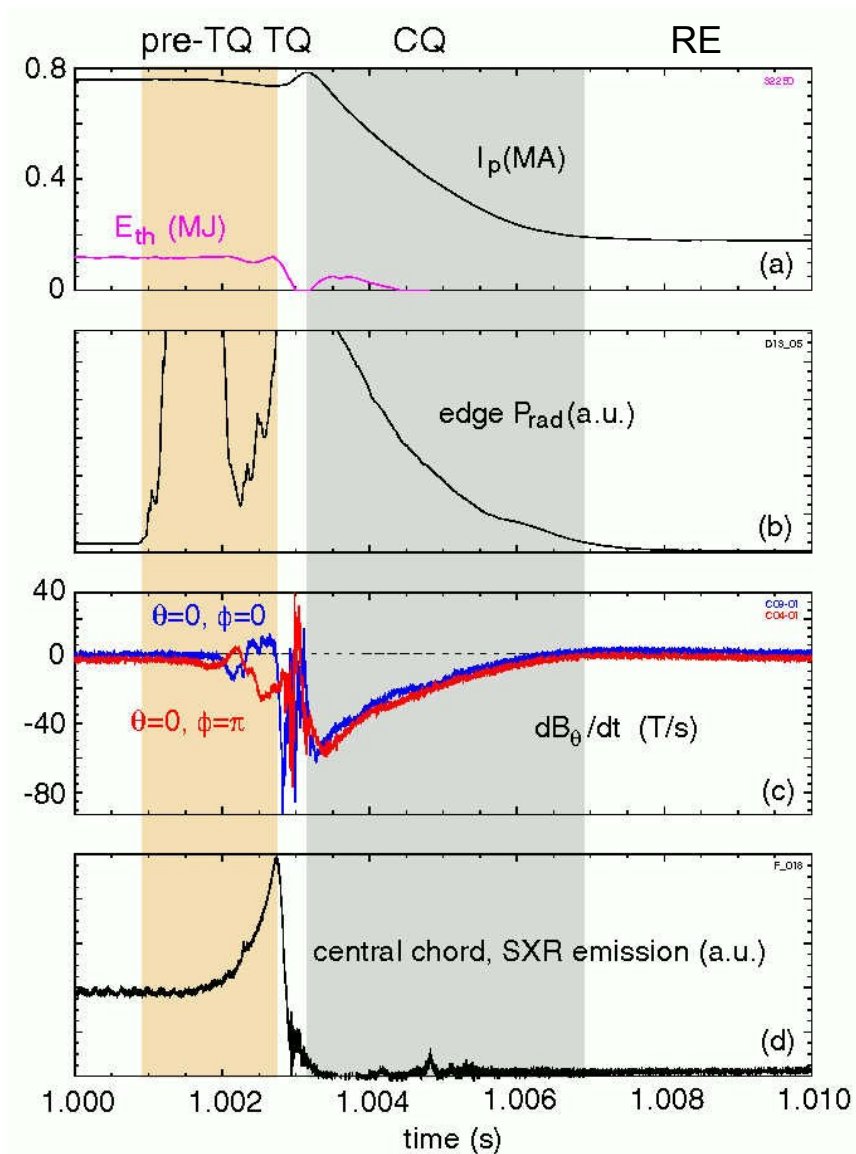


Fast poloidal coil measurement

TQ

fast CQ

Duration of disruption phases, scatter, $\rightarrow I_{RE}$



0D RE generation and suppression model

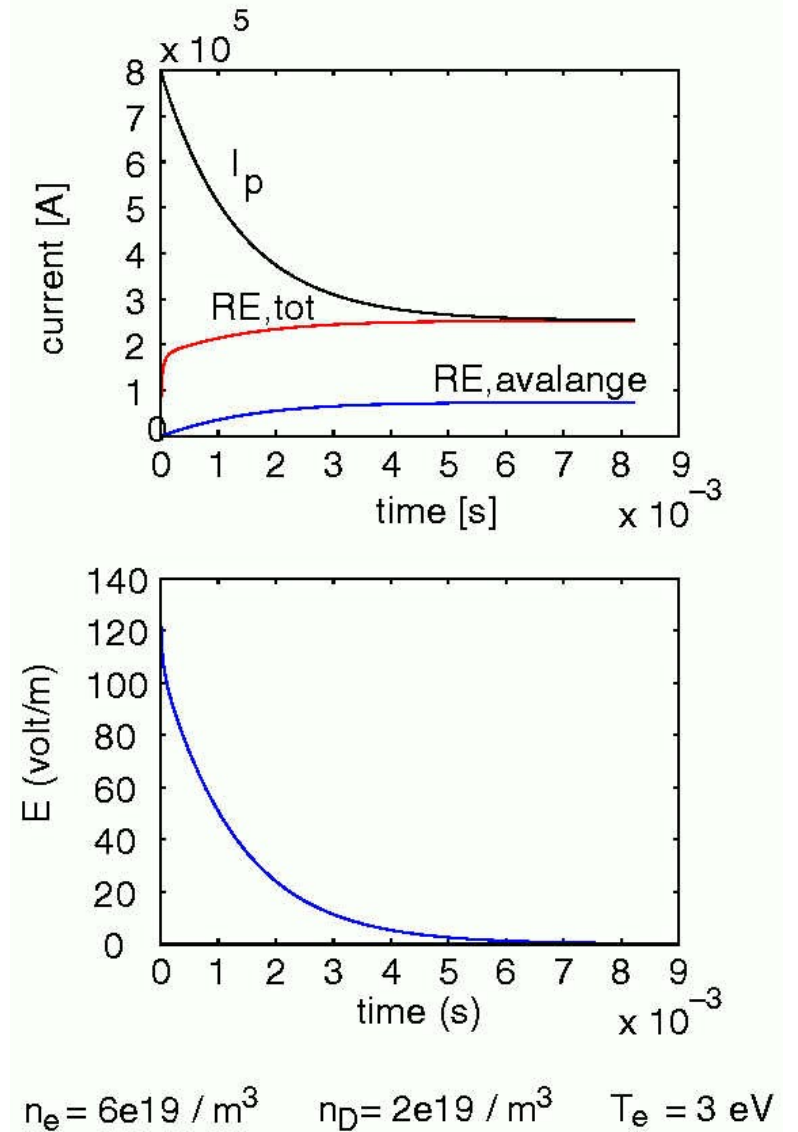
$$E = -\frac{dI_p}{dt} \frac{L}{2\pi R_0} F_{mag}$$

$$E = \rho(j_p - j_{RE})$$

$$\frac{dn_{RE}}{dt} = \left[\frac{dn_{RE}}{dt} \right]_1 + \left[\frac{dn_{RE}}{dt} \right]_2$$

$$j_p = I_p / (\pi a^2 k)$$

$$j_{RE} = I_{RE} / (\pi a^2 k) = c e n_{RE}$$



neither losses nor hot-tail; 1 = primary generation; 2 = avalanche + impurities

T_e and Z_{eff} from 0D model

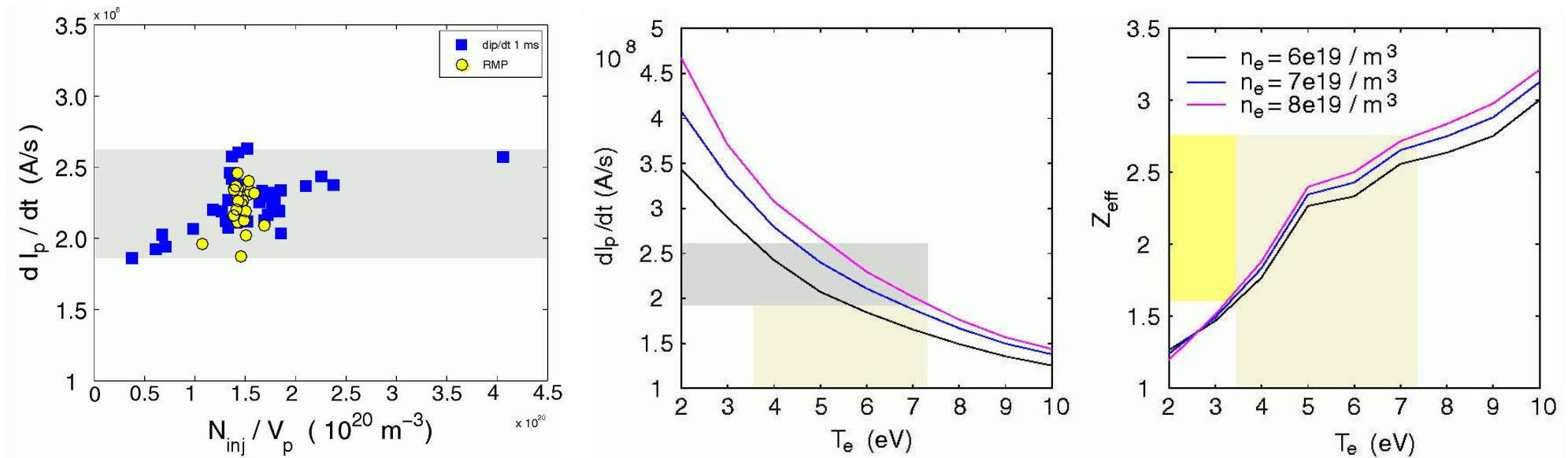
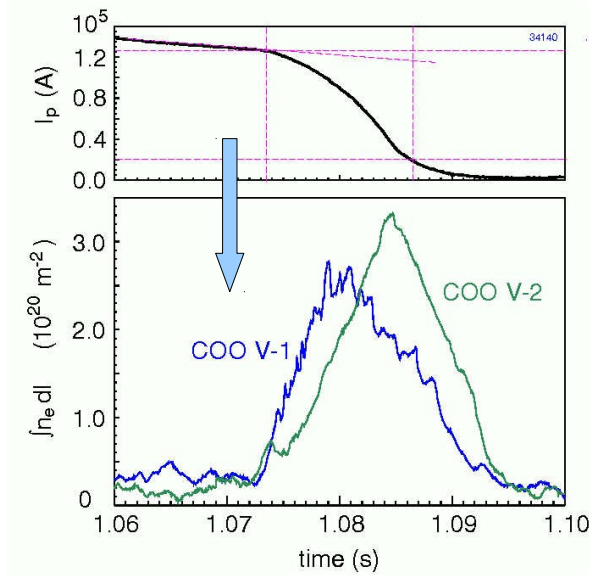


Fig. left: Measured fast I_p decay rate after TQ

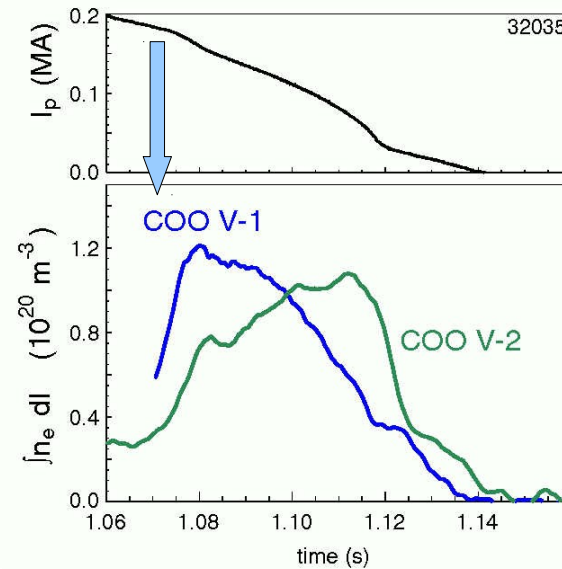
Centre: calculated I_p decay versus T_e

Right: calculated effective charge (Z_{eff}) versus T_e

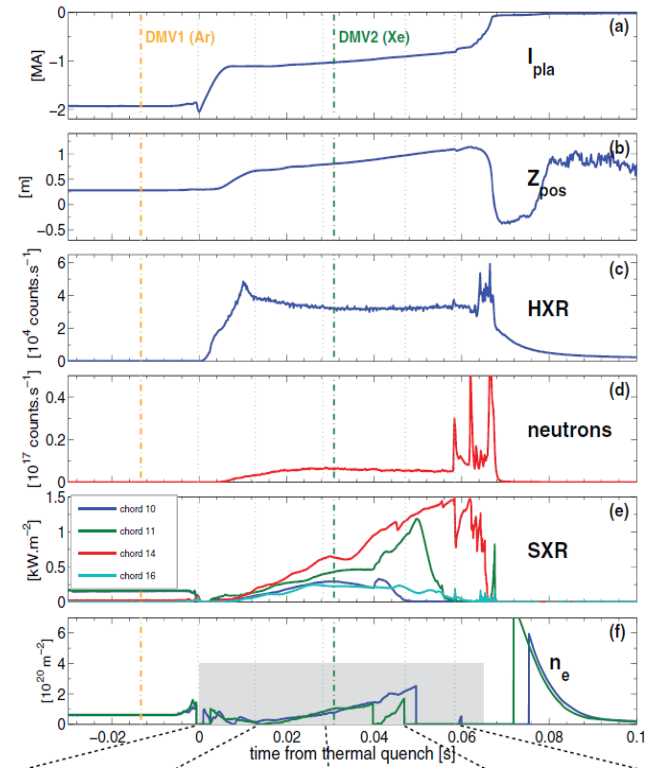
MGI injection (2nd) in RE beam



argon (AUG)



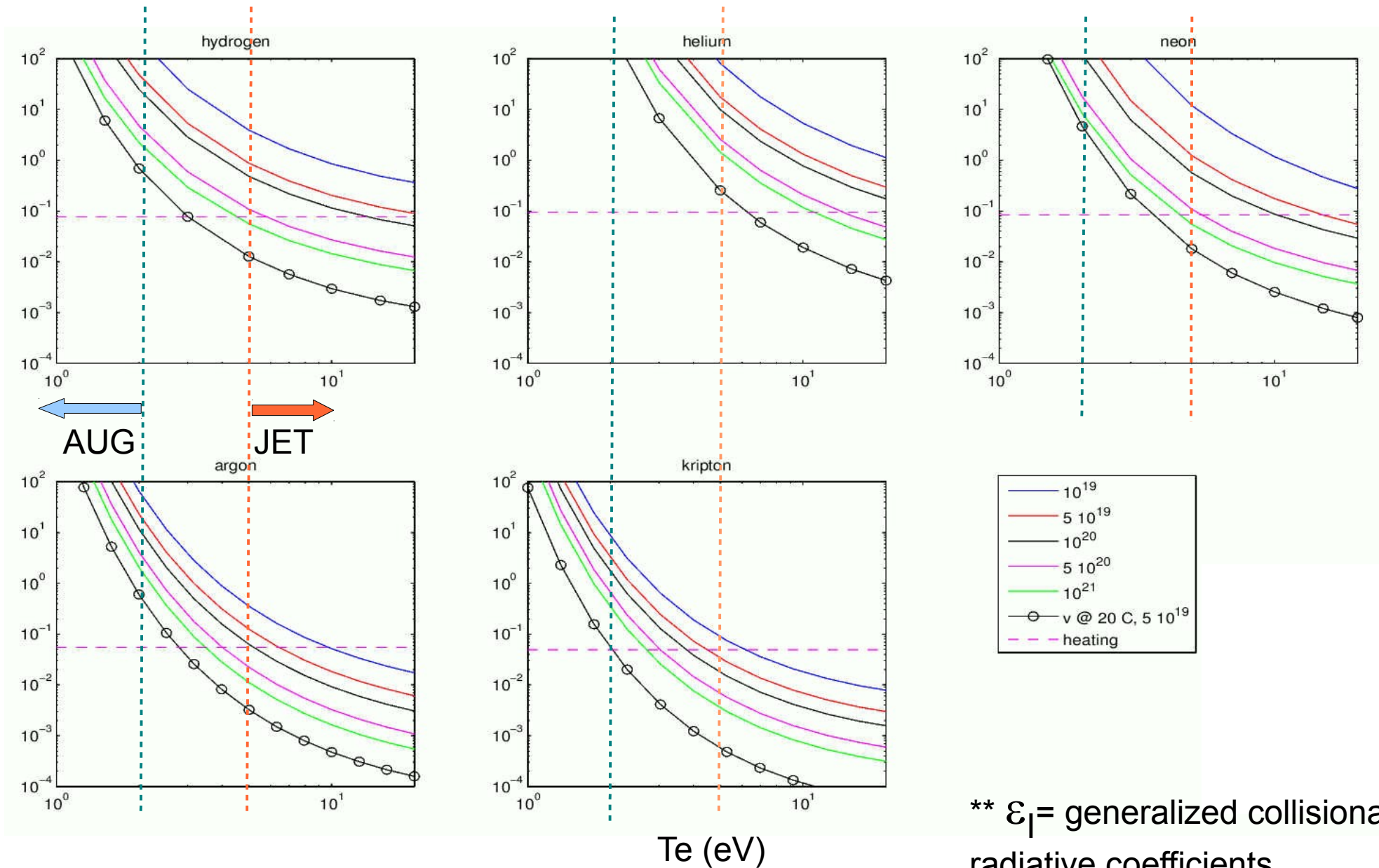
neon (AUG)



xenon (JET) **

- high Z gas dissipates RE current in AUG, immediate effect (ms); effect consistent with collisional dissipation \longrightarrow = trigger time
- no significant effect in JET (30 ms) **C. Reux et al, Nuclear Fusion 55 (2015)
- note: background plasma $T_e < 2 \text{ eV}$ in AUG; $T_e \sim 5\text{-}15 \text{ eV}$ in JET (C. Reux et al, TSDW 2017)
- note: $d(n_e)/dt$ measured reflects particle flow from valve

Ionization mean free path, $\text{mfp}_i = v_{\text{th},i} / (\epsilon_i n_e) \text{ (m)}$ **



** ϵ_i = generalized collisional radiative coefficients

Gas diffusion into RE beam

The diffusion equation in cylindrical coordinates

$$\frac{\partial n(r, t)}{\partial t} = \frac{1}{r} \frac{\partial}{\partial r} r D(r, t) \frac{\partial n(r, t)}{\partial r} + S(r, t) \quad (1)$$

was solved explicitly by the finite difference method and the assumptions

$$D_i = \rho_i^2 \nu_{ii}, \quad \rho_i = \frac{m_i v_{\perp, i}}{e Z_i B_t}, \quad v_{\perp, i} = \sqrt{2} v_{th, i} = \sqrt{2 T_i / m_i}, \quad [14] \quad (2)$$

$$\frac{1}{\nu_{ii}} = \tau_{ii} = 12 \pi^{3/2} \frac{\epsilon_0^2 m_i^{1/2} T_i^{3/2}}{n_i Z_i^4 e^4 \ln \Lambda_{ii}} \quad [14], \quad (3)$$

$$\ln \Lambda_{ii} = 23 - \ln \left(\frac{Z_i^3}{T_i} \sqrt{\frac{n_i}{T_i}} \right), \quad T_i \text{ in eV} \quad [15] \quad (4)$$

and a source of particles at the plasma edge of intensity

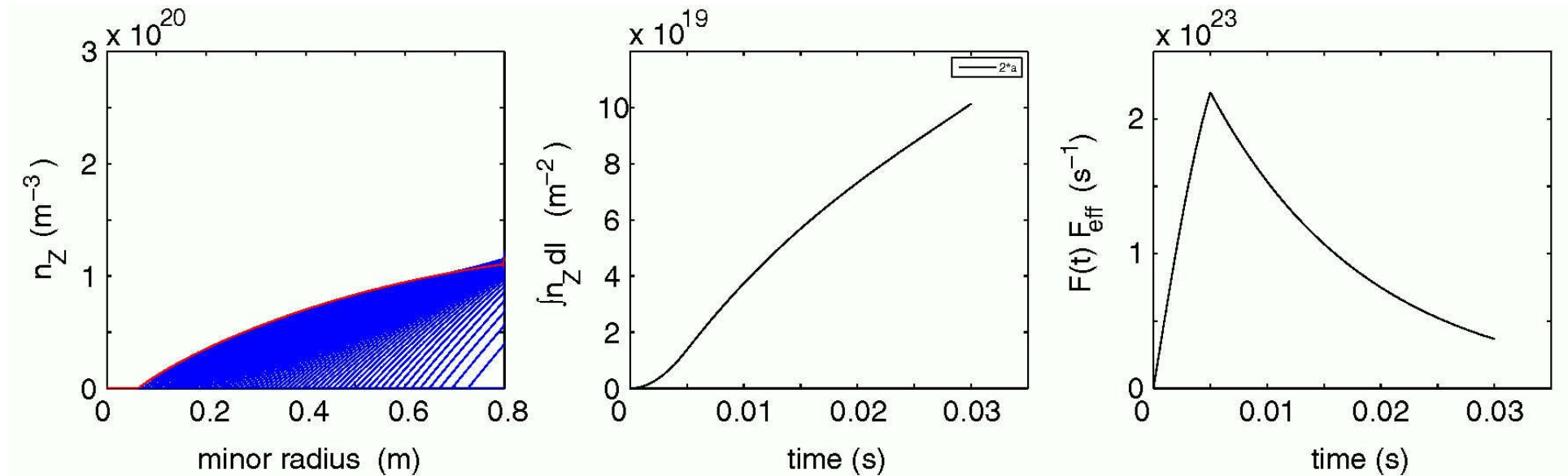
$$S(r, t) = \frac{F(t)}{4 \pi^2 a R_0} \delta(r - a) F_{\text{eff}} \quad (5)$$

were $F(t)$ is the gas flow from the valve, F_{eff} is the fueling efficiency, a and R_0 are the minor and major radius of the circular plasma.

[14] J. Wesson, Tokamaks

[15] NRL plasma formulary

High Z gas diffusion into JET RE beam



- JET RE beam like, injection of $2472 \text{ Pa} \times \text{m}^3$ xenon *C. Reux et al, Nuclear Fusion 55 (2015)*
- $T_e = 5 \text{ eV}$ and $Z_{\text{Xe}} = 3$; 30 ms run
- DMV2 vol. = $9.75 \times 10^{-4} \text{ m}^3$, diameter orifice = 30 mm (Kruezi SOFT 2014)
- **$F_{\text{eff}} = 0.6 \%$** (to approximate n_e measured – from publication)
- small F_{eff} is limiting gas penetration towards RE beam

Summary

- Scatter of I_{RE} after argon-induced disruptions and lack of N_{inj} dependence motivated (1) careful analysis of density measurements and (2) calculation of RE generation with simple 0-D fluid model
- Within uncertainties affecting density profile, experimental measurements fall into ballpark of calculated I_{RE} ; scatter difficult to explain
- In AUG, RE beam background plasma has $T_e < 2$ eV \rightarrow it is transparent to injected impurities
- JET plasma is has hotter plasma (7 eV) \rightarrow gas injected must diffuse from edge into core
- Nevertheless, gas assimilation and not (or not only) diffusion limits effectiveness of 2nd high Z MGI

Additional slides

