

Halo current studies with self-consistent MHD simulations for ITER 15 MA plasmas

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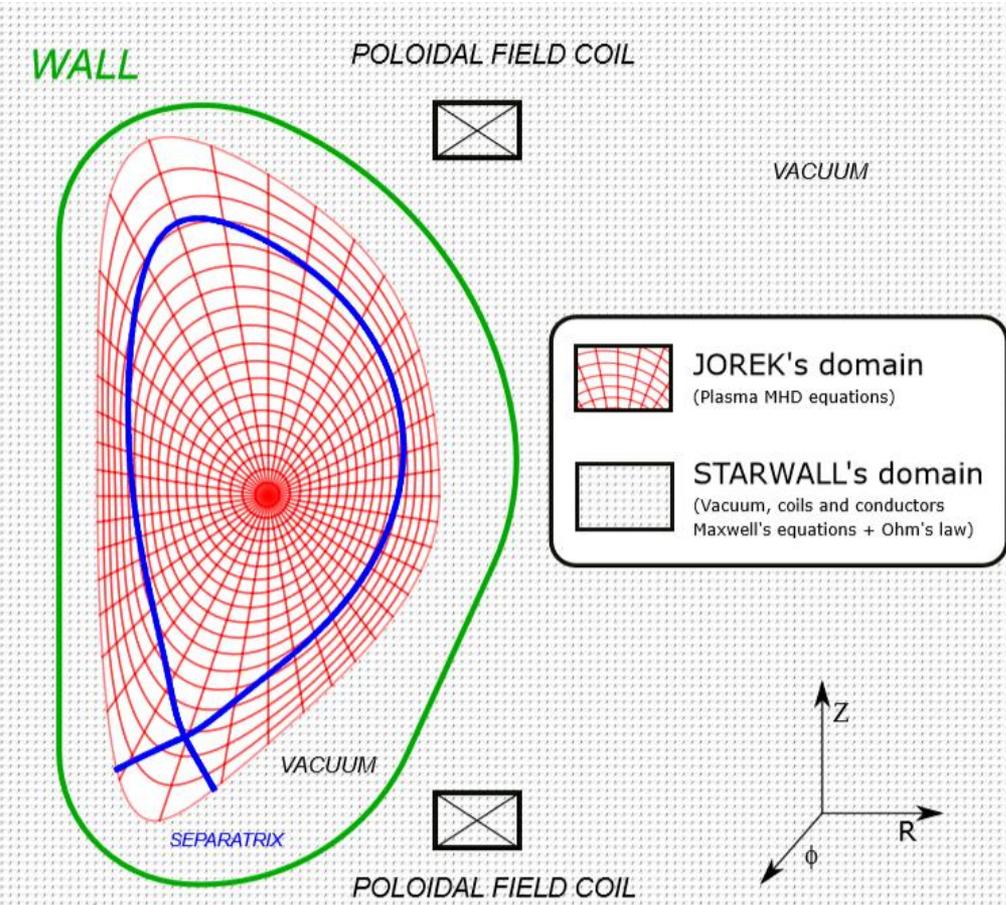
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Outline

- The JOREK-STARWALL code for halo current modelling
- 2D VDE benchmark with M3D-C1 and NIMROD
- Understanding the 2D halo currents at ITER (15 MA/5.3T)
- Prediction of the halo properties and B.C.s
- Conclusions and future work

The JOREK-STARWALL code for halo current modelling

The JOREK-STARWALL code for halo current modelling



JOREK [Huysmans, NF2007]

- 3D non-linear MHD equations
- Toroidal geometry
- C1 Finite elements in poloidal plane
- Fourier harmonics for ϕ direction
- Fully implicit time evolution

STARWALL [P. Merkel, 2015]

- Solves Maxwell's equations and Ohm's law
- Green's function method
- Thin wall approximation

Implicit coupling through B.C.s for \vec{B}

[Hoelzl, 2012]

$$\vec{B}_{\text{tan}} = \vec{f}(B_n, \mathbf{Y})$$

Tangential field Normal field Wall currents

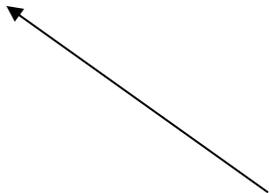
EM boundary conditions

Reduced MHD, the E-field is

$$\mathbf{E} = -\partial_t \psi \nabla \phi - F_0 \nabla_{\text{pol}} u$$

- $\partial_t \psi$ has resistive wall free-boundary conditions
- Ideal wall BCs for poloidal E-field ($u = 0$)
- Poloidal currents calculated from

force balance $\mathbf{J} \times \mathbf{B} = \nabla p$

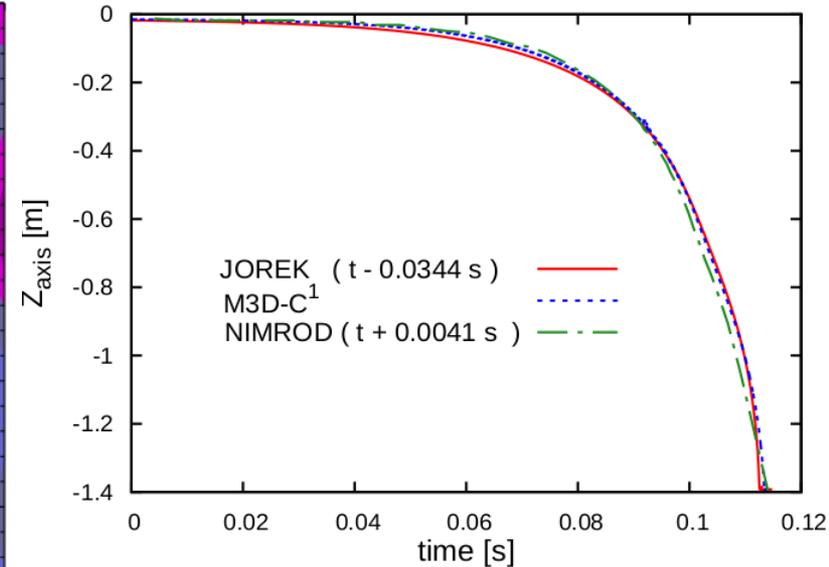
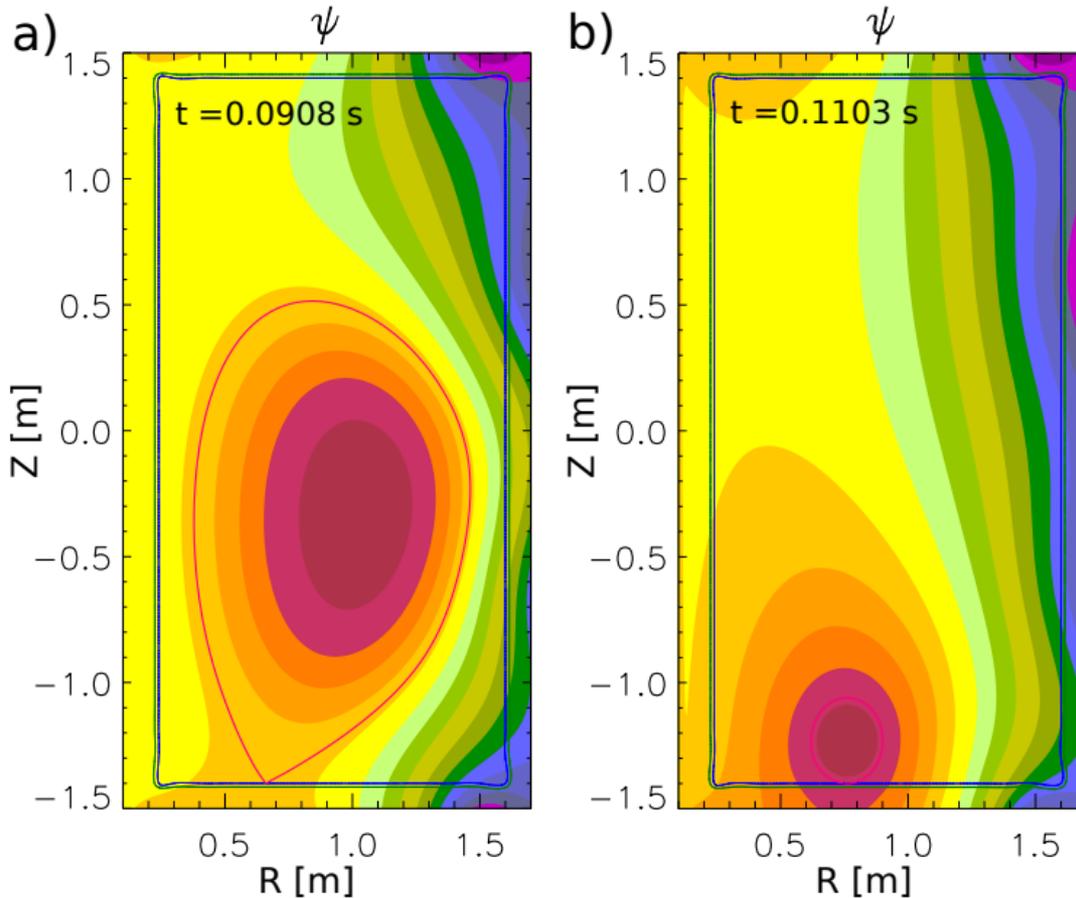


Strong assumption. Poloidal currents do not decay in the wall (infinite conductivity in the poloidal direction)

2D VDE benchmark with M3D-C1 and NIMROD

2D VDE benchmark with M3D-C1 and NIMROD

VDE based on the NSTX #139536 discharge

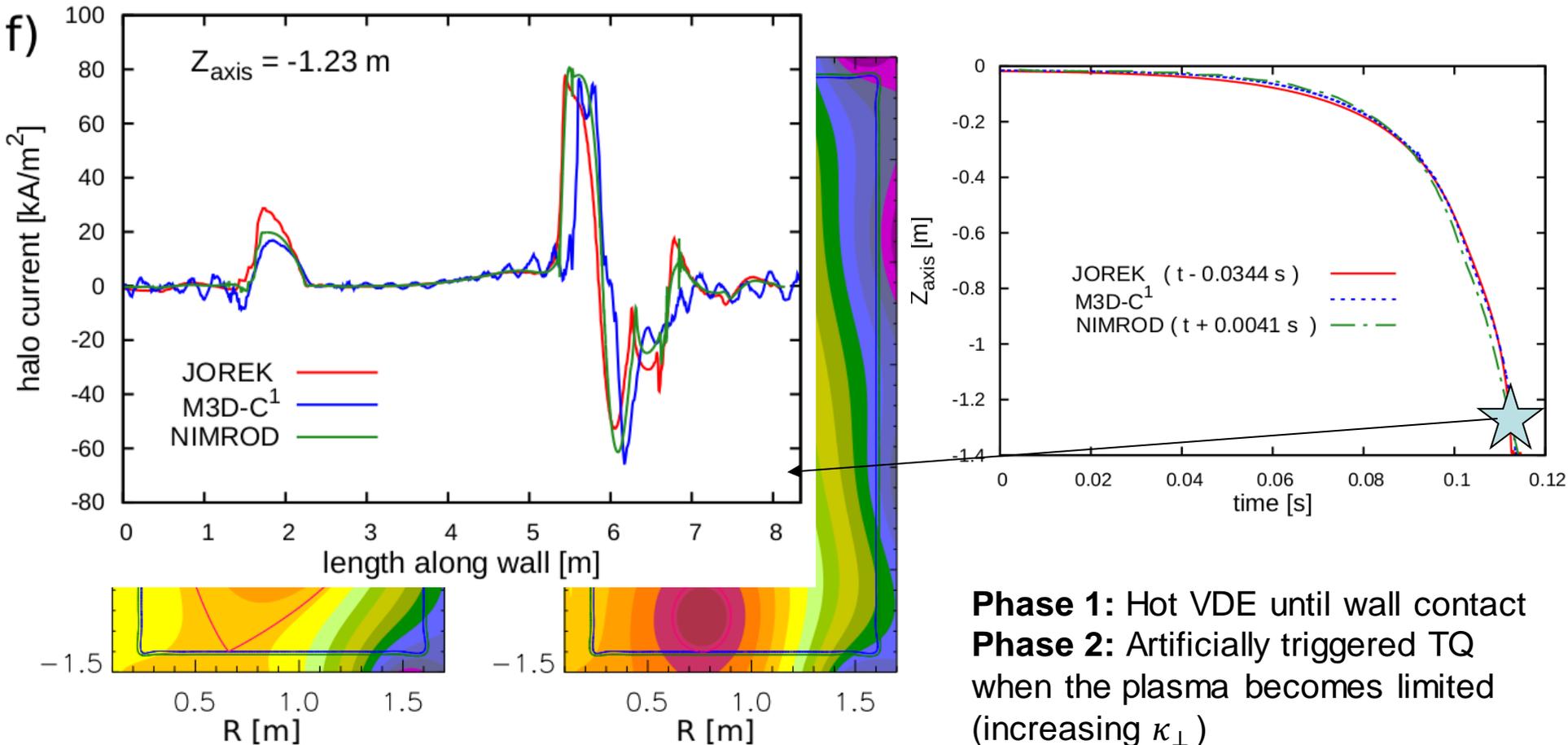


Phase 1: Hot VDE until wall contact
Phase 2: Artificially triggered TQ
when the plasma becomes limited
(increasing κ_{\perp})

[I. Krebs, F.J. Artola, C. Sovinec 2019]

2D VDE benchmark with M3D-C1 and NIMROD

VDE based on the NSTX #139536 discharge



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Understanding 2D halo currents at ITER (15 MA / 5.3T)

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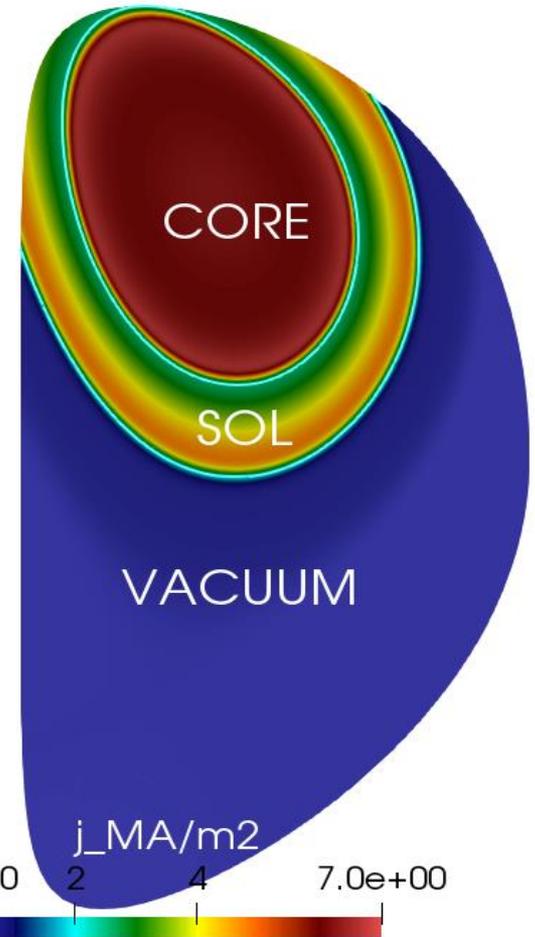
Halo current different regimes

- **Hot VDE regime** ($\tau_w \ll \tau_p$)

$I_p \sim \text{cte}$ during VDE

- **Ideal wall regime** ($\tau_p \ll \tau_w$)

$Z_{axis} = f(I_p)$ [D. Kiramov 2017 PoP]

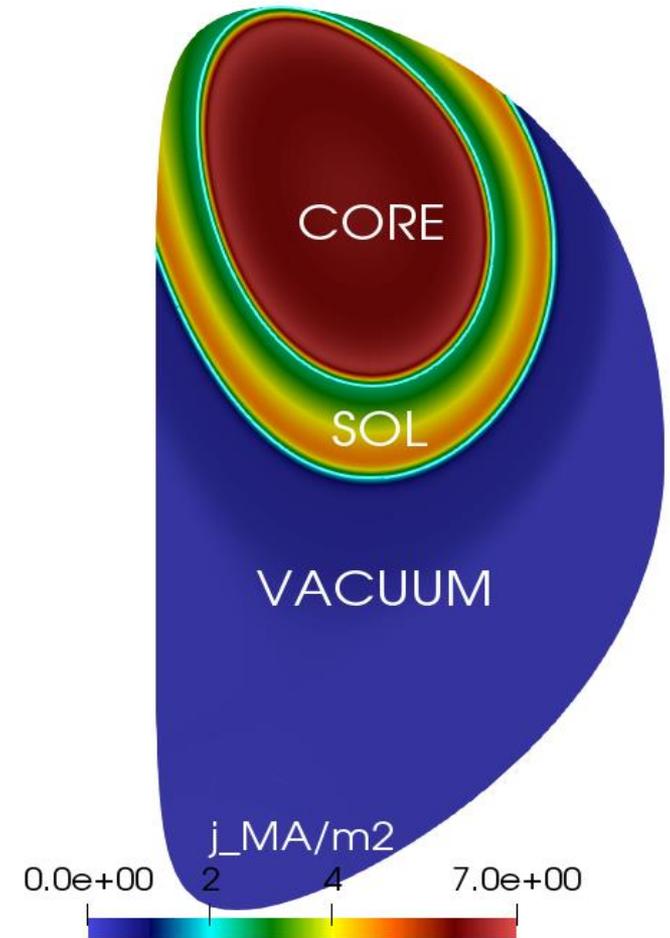
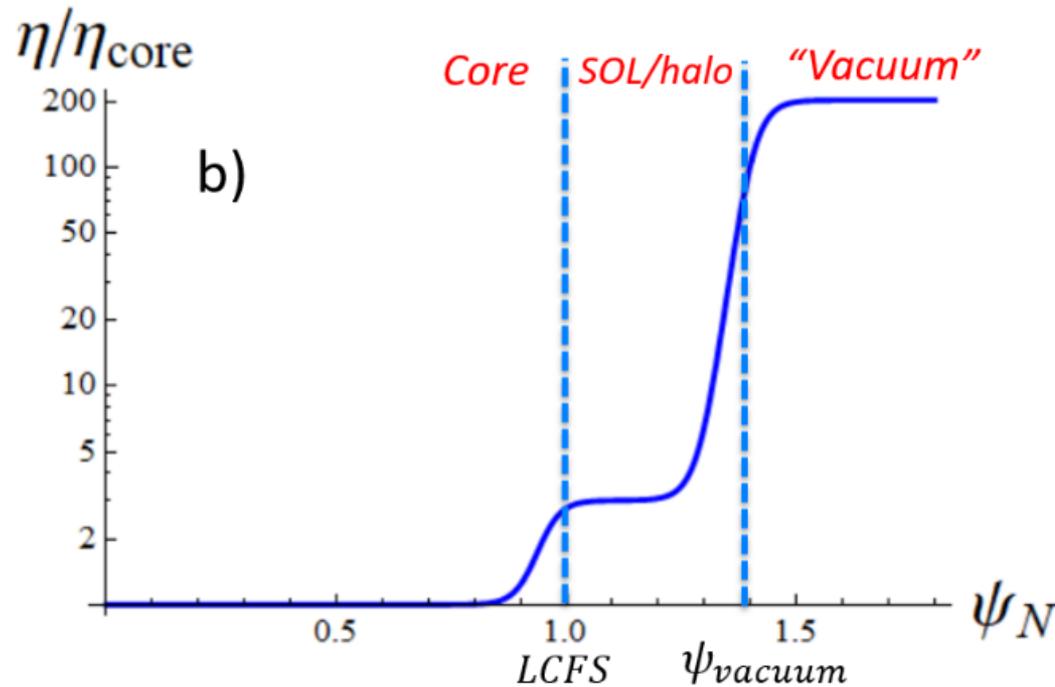


$\tau \equiv$ Current resistive decay time

Understanding 2D halo currents at ITER (15 MA / 5.3T)

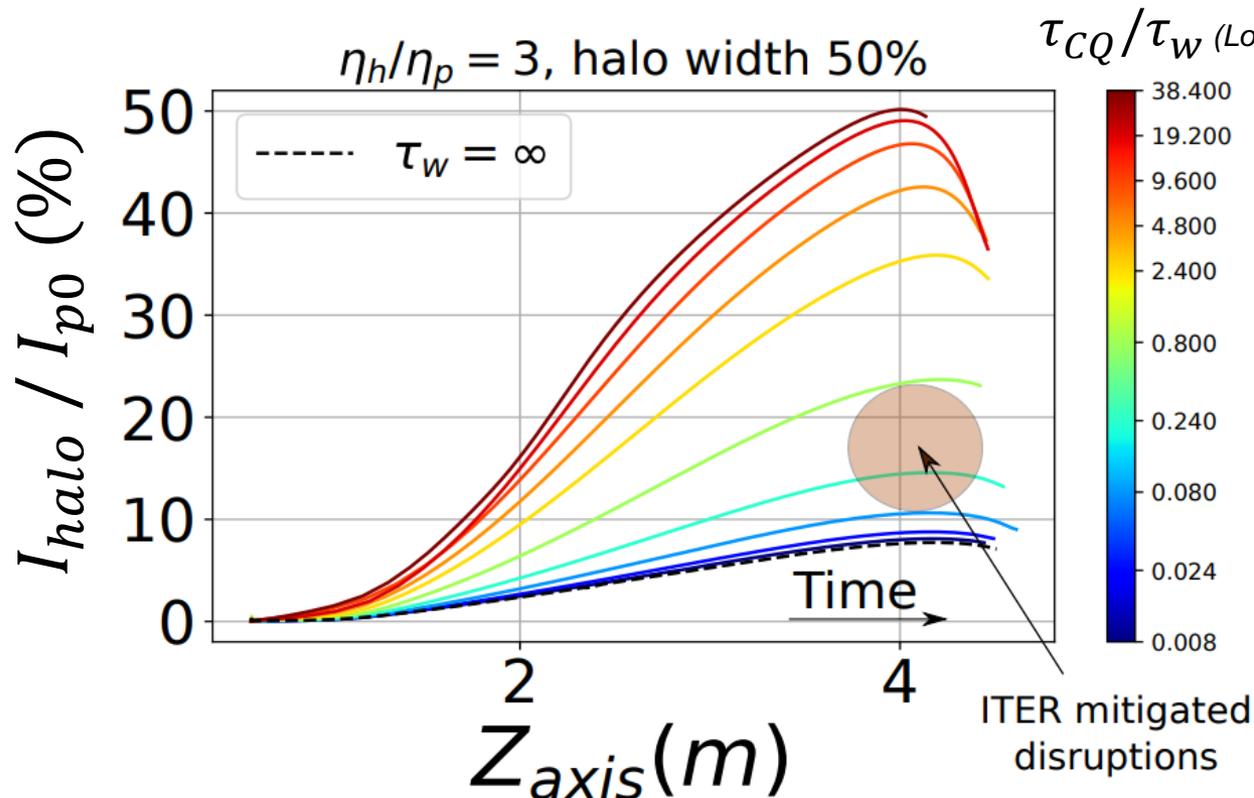
Parametric scan in CQ time

Plasma resistivity is prescribed as a flux function



Understanding 2D halo currents at ITER (15 MA / 5.3T)

Scan in CQ time (scaling plasma resistivity profile)



- **Strong dependence** on CQ time to wall time ratio (τ_{CQ}/τ_w)

- Maximum at $\tau_{CQ}/\tau_w \rightarrow \infty$

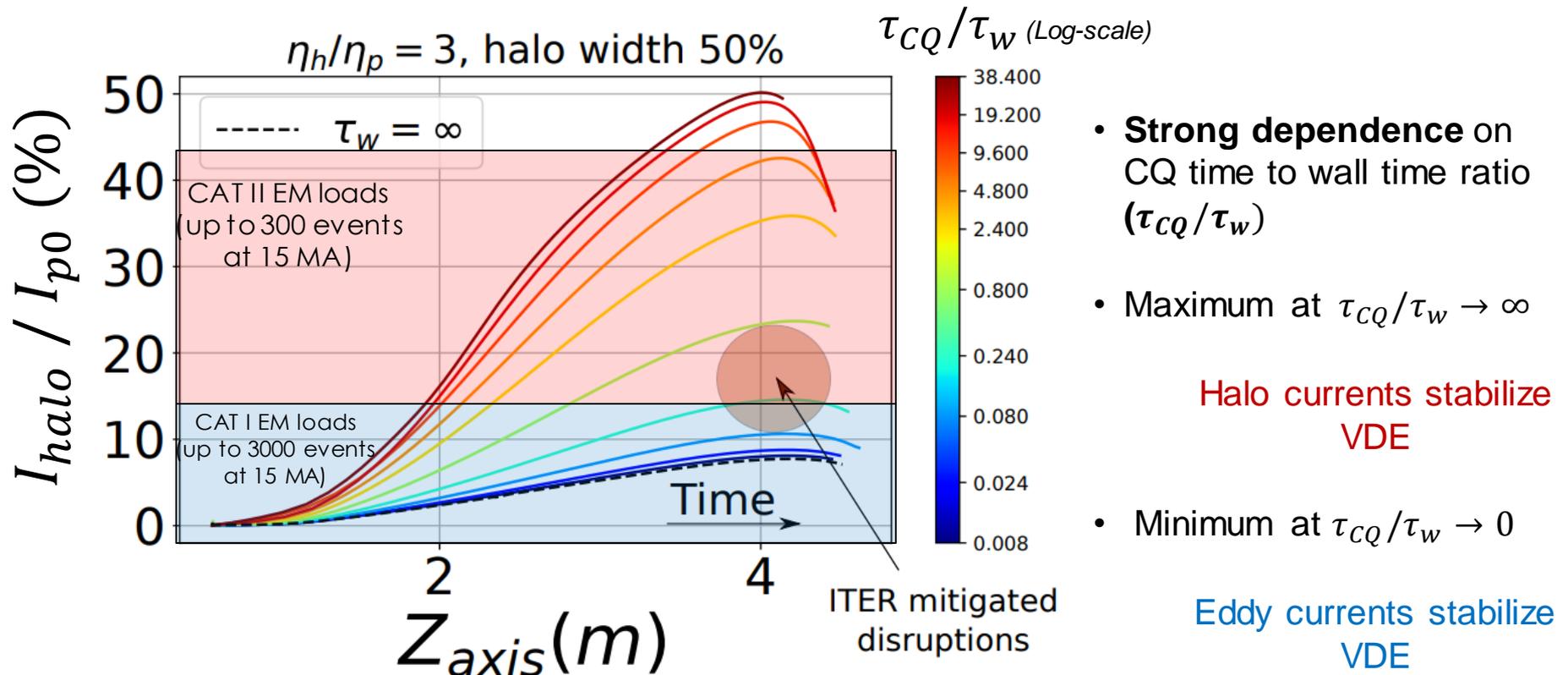
Halo currents stabilize VDE

- Minimum at $\tau_{CQ}/\tau_w \rightarrow 0$

Eddy currents stabilize VDE

Understanding 2D halo currents at ITER (15 MA / 5.3T)

Scan in CQ time (scaling plasma resistivity profile)



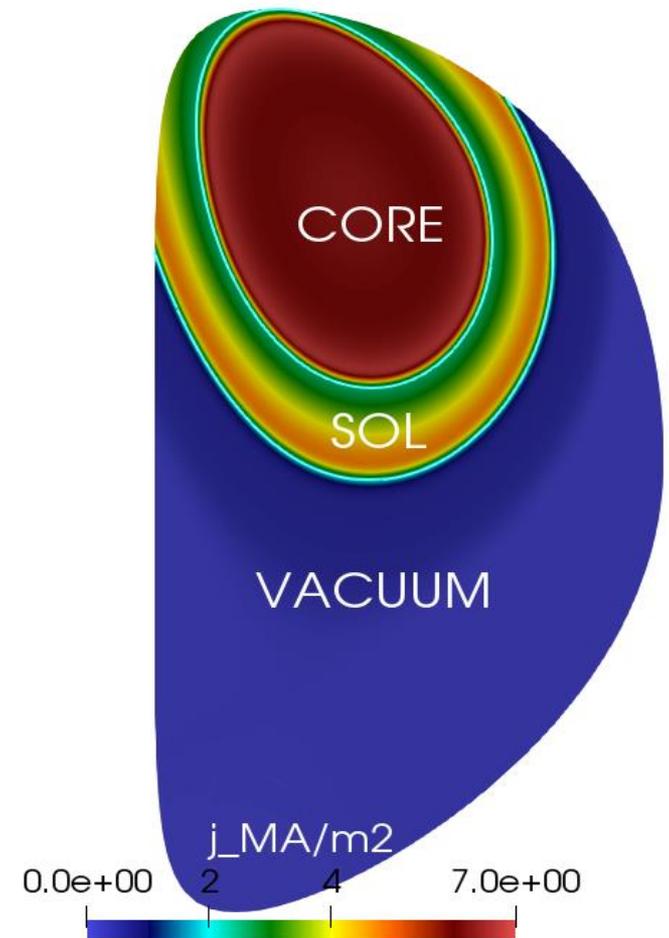
Understanding 2D halo currents at ITER (15 MA / 5.3T)

Halo current different regimes

- **Hot VDE regime** ($\tau_w \ll \tau_p$)
 - **Cold halo** ($\tau_h \ll \tau_w$)
 - **Hot halo** ($\tau_w \ll \tau_h$)
- **Ideal wall regime** ($\tau_p \ll \tau_w$)
 - **Cold halo** ($\tau_h \ll \tau_p$)
 - **Hot halo** ($\tau_p \ll \tau_h$)

also discussed in [Boozer PoP 2013]

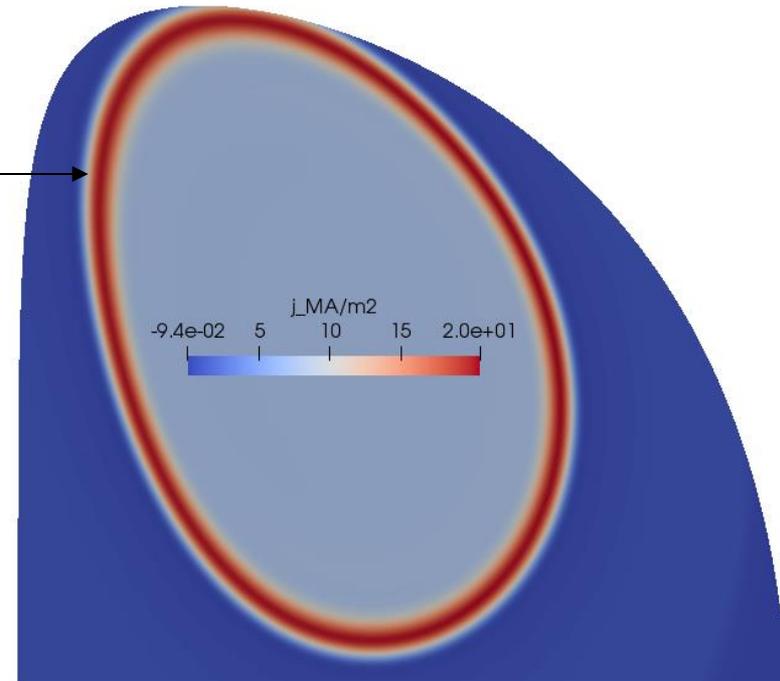
$\tau \equiv$ Current resistive decay time



Understanding 2D halo currents at ITER (15 MA / 5.3T)

Hot VDE + cold halo ($\tau_h \ll \tau_w \ll \tau_p$)

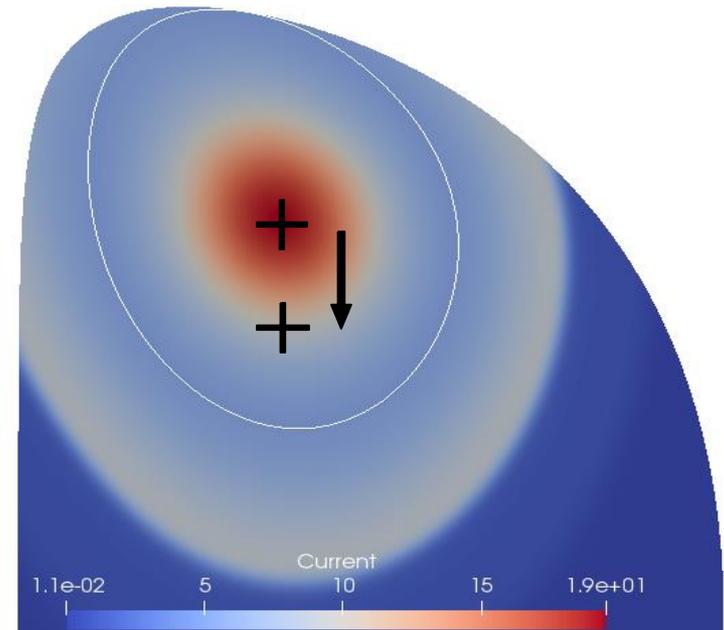
- Currents are lost in wall and halo faster than in plasma core
- Currents are re-induced in plasma edge (large current densities) →
- **Big drop of edge safety factor**
 I_p is largely conserved ($q_a \propto a^2$)
- **Potential destabilization of external kink modes**



Understanding 2D halo currents at ITER (15 MA / 5.3T)

Hot VDE + hot halo $(\tau_w \ll \tau_h \sim \tau_p)$

- Toroidal current is transferred into the halo region as the plasma moves vertically
- Halo currents stabilize vertical motion
- After stabilization, motion is given by resistive decay of core + halos

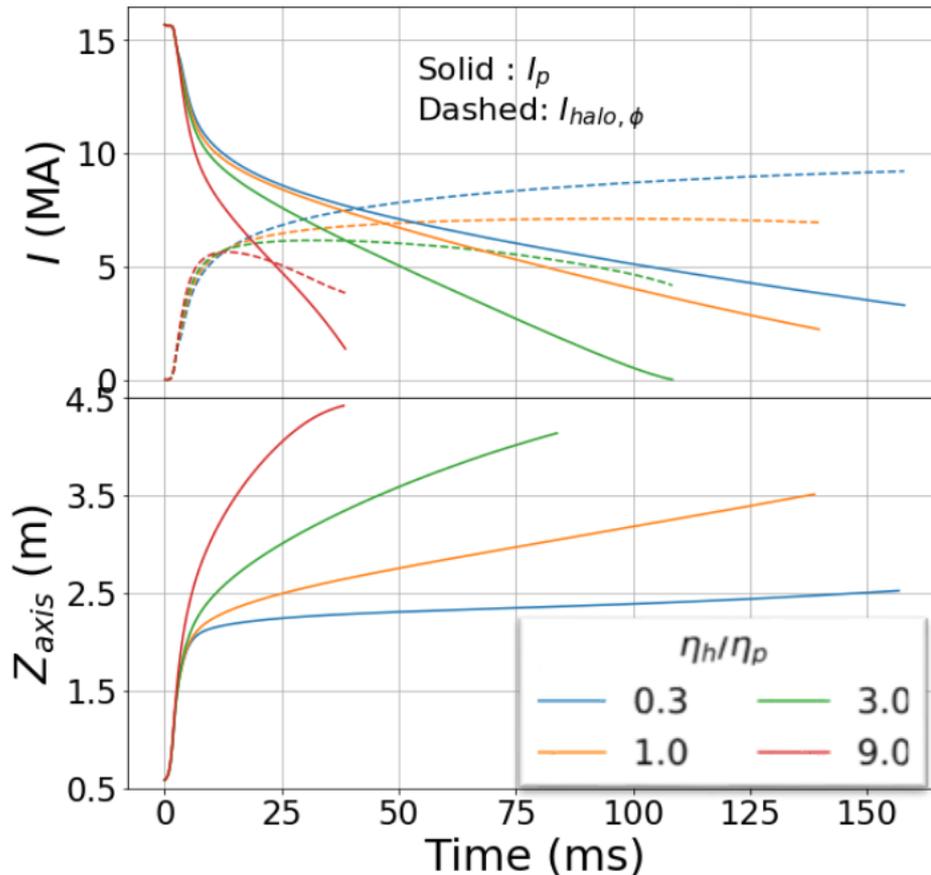


Understanding 2D halo currents at ITER (15 MA / 5.3T)

Hot VDE + hot halo ($\tau_w \ll \tau_h \sim \tau_p$)

Halo resistivity scan

$$\tau_{CQ}/\tau_w = 64$$

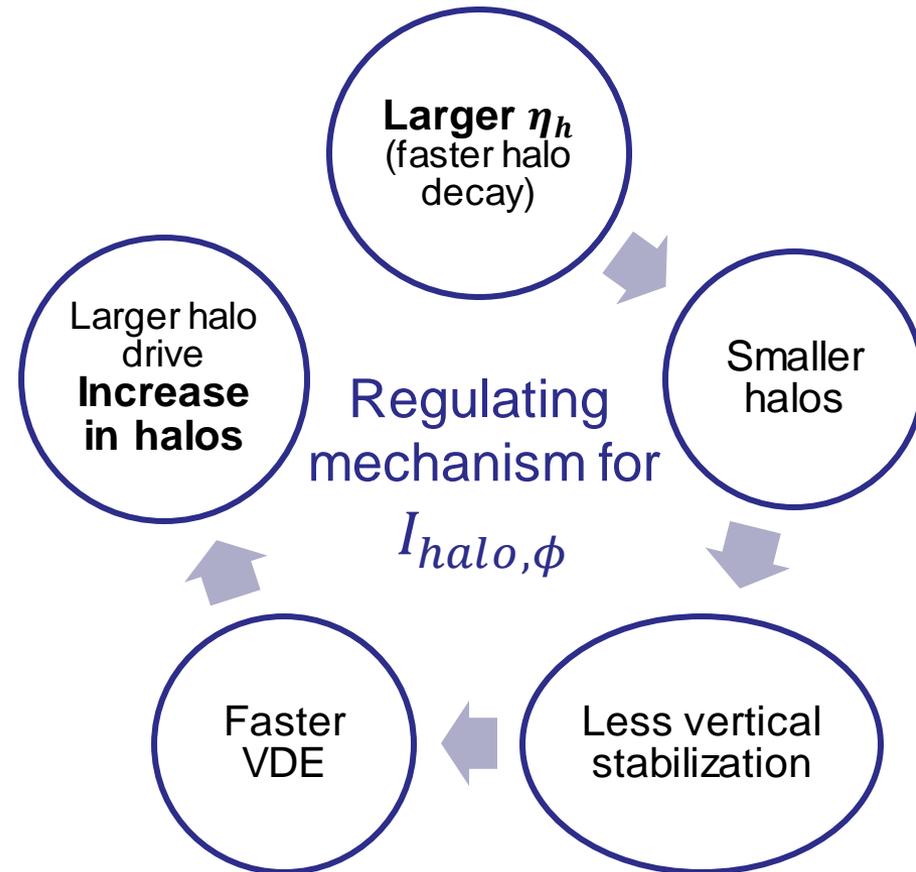
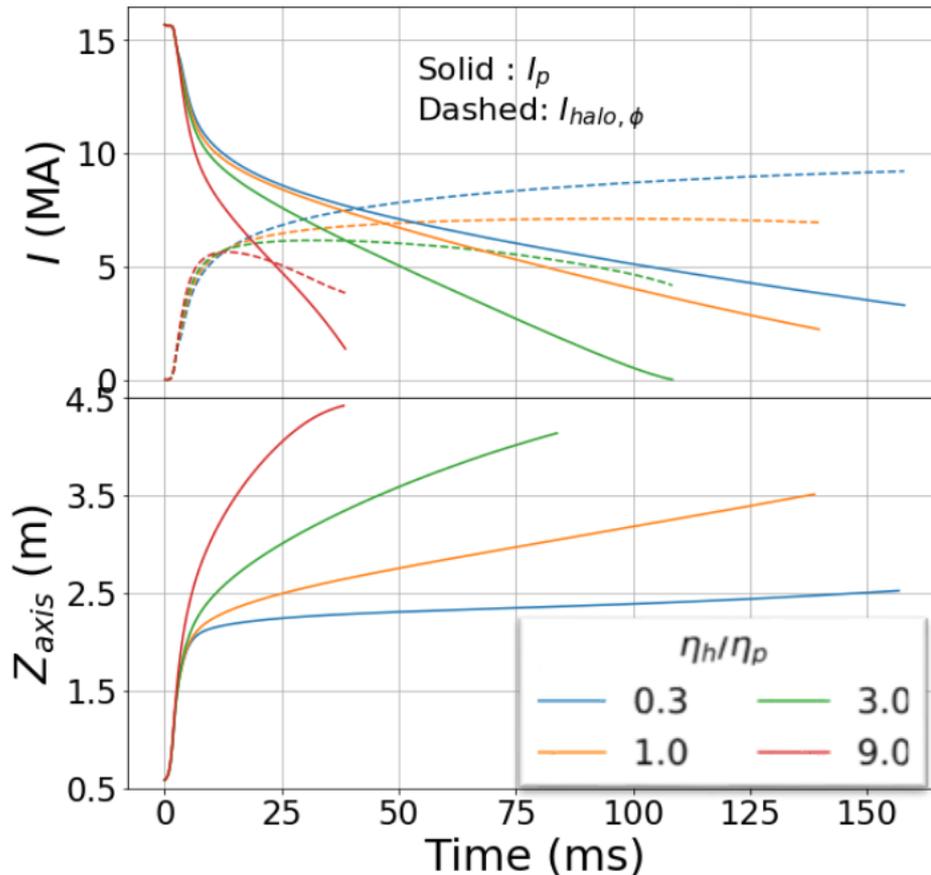


Understanding 2D halo currents at ITER (15 MA / 5.3T)

Hot VDE + hot halo ($\tau_w \ll \tau_h \sim \tau_p$)

Halo resistivity scan

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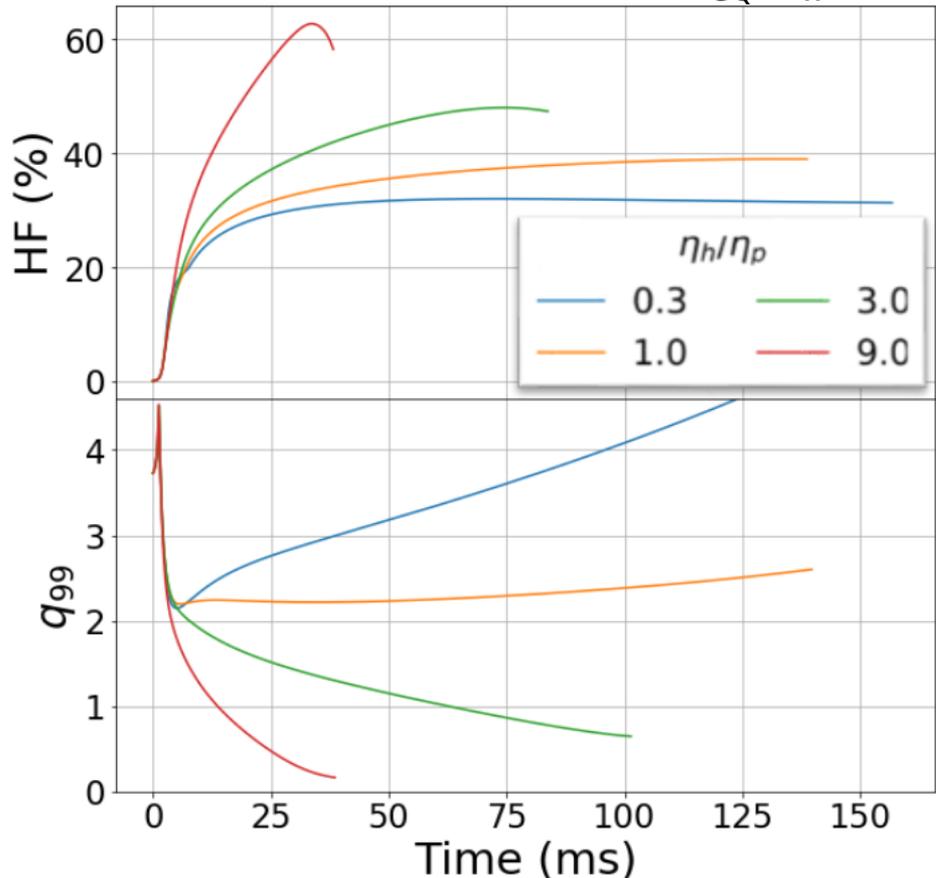
$I_{halo, \phi}$ depends weakly on η_h

Understanding 2D halo currents at ITER (15 MA / 5.3T)

Hot VDE + hot halo ($\tau_w \ll \tau_h \sim \tau_p$)

Halo resistivity scan

$\tau_{CQ}/\tau_w = 64$



But $I_{halo,pol}$ depends strongly on η_h through q_a

$$I_{halo,pol} \sim \frac{I_{halo,\phi}}{q_a}$$

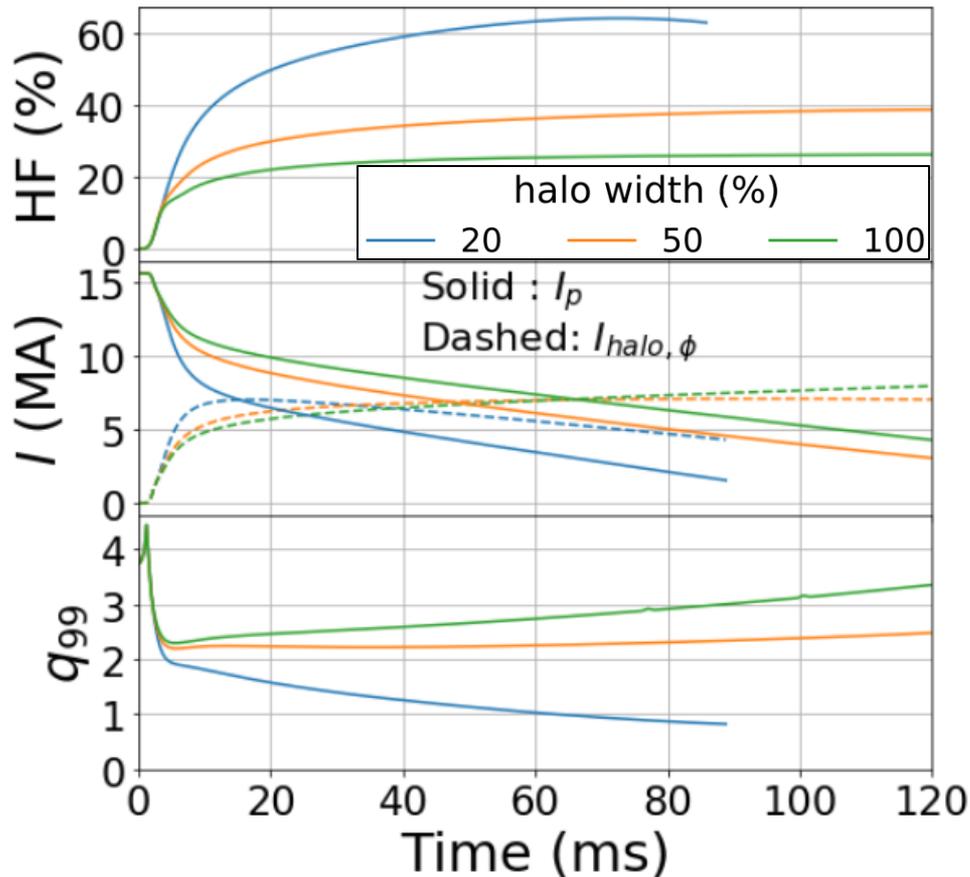
Finally increasing the halo resistivity gives larger poloidal halo currents

Understanding 2D halo currents at ITER (15 MA / 5.3T)

Hot VDE + hot halo ($\tau_w \ll \tau_h \sim \tau_p$)

Halo width scan

$$\tau_{CQ}/\tau_w = 64$$



$I_{halo,\phi}$ also has a weak dependence on the halo width

Weak dependence through

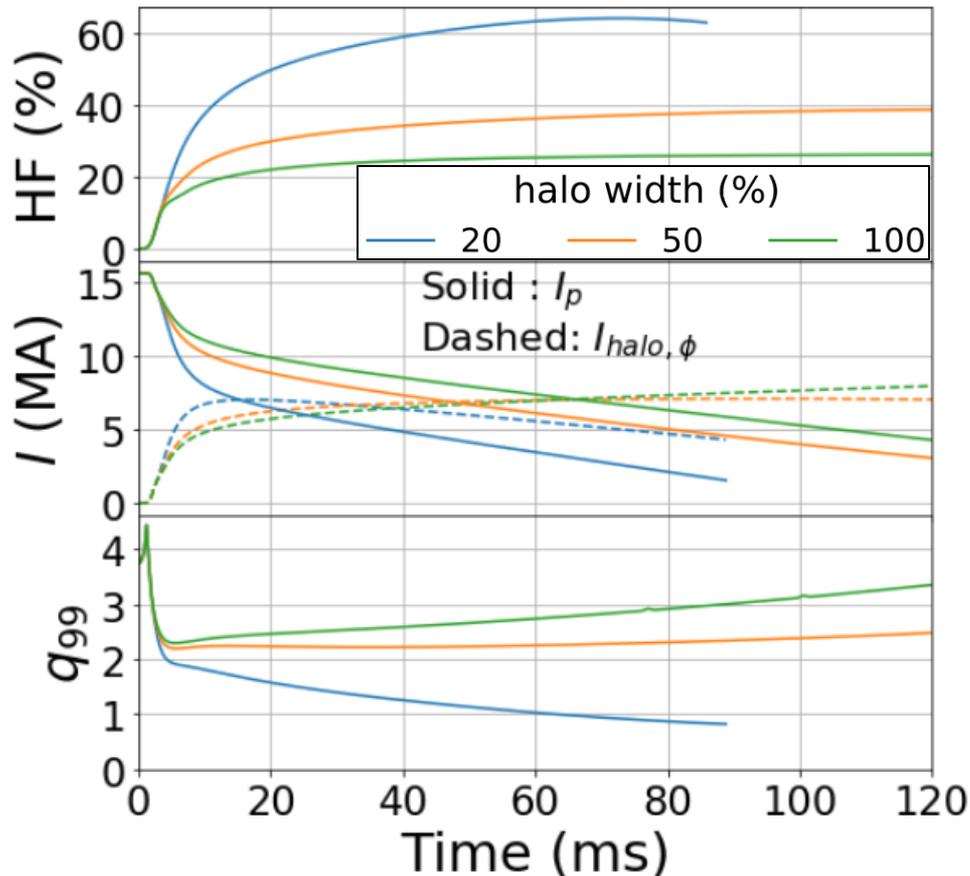
$$\tau_h \propto w_h/\eta_h$$

Understanding 2D halo currents at ITER (15 MA / 5.3T)

Hot VDE + hot halo ($\tau_w \ll \tau_h \sim \tau_p$)

Halo width scan

$$\tau_{CQ}/\tau_w = 64$$



$I_{halo,\phi}$ also has a weak dependence on the halo width

Weak dependence through

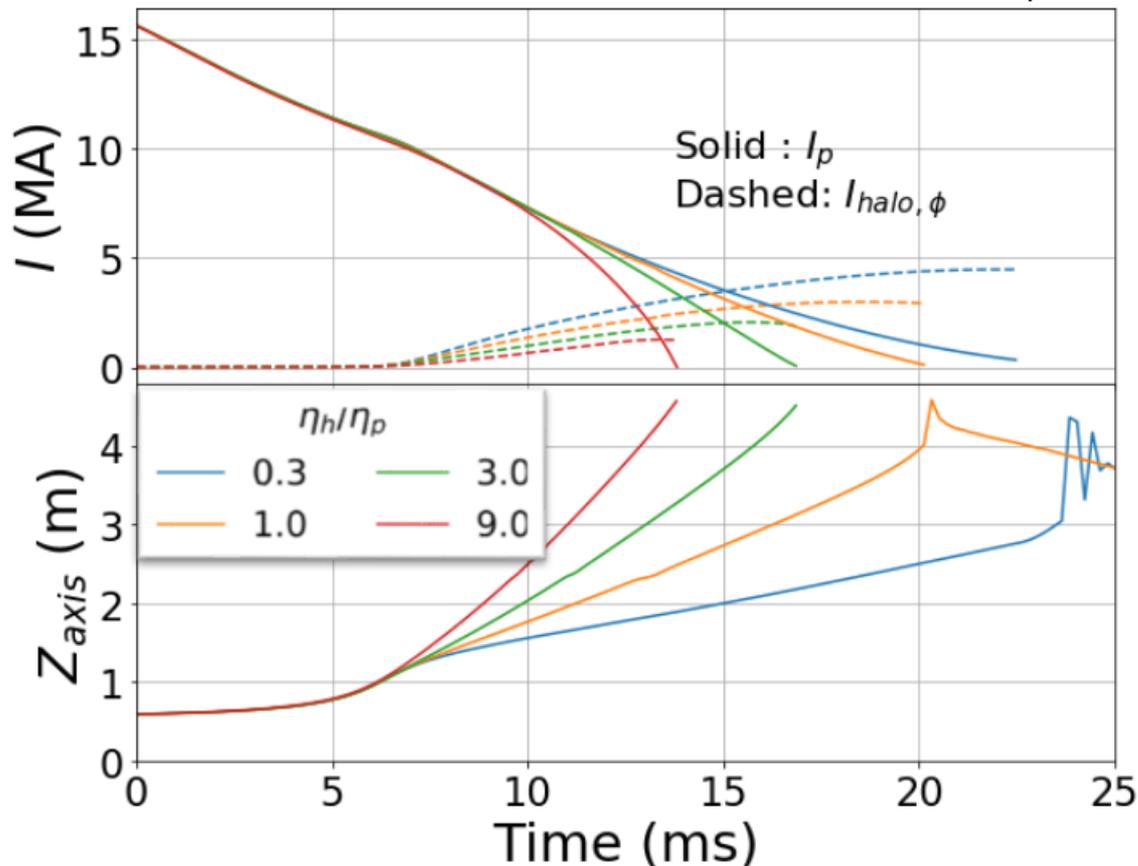
$$\tau_h \propto w_h/\eta_h$$

But $I_{halo,pol}$ has a much stronger dependence through q_a . **More effective at narrow halos.**

Understanding 2D halo currents at ITER (15 MA / 5.3T)

Ideal wall regime ($\tau_p \ll \tau_w$)

Halo resistivity scan $\tau_{CQ}/\tau_w = 0, \eta_h/\eta_p = 1$



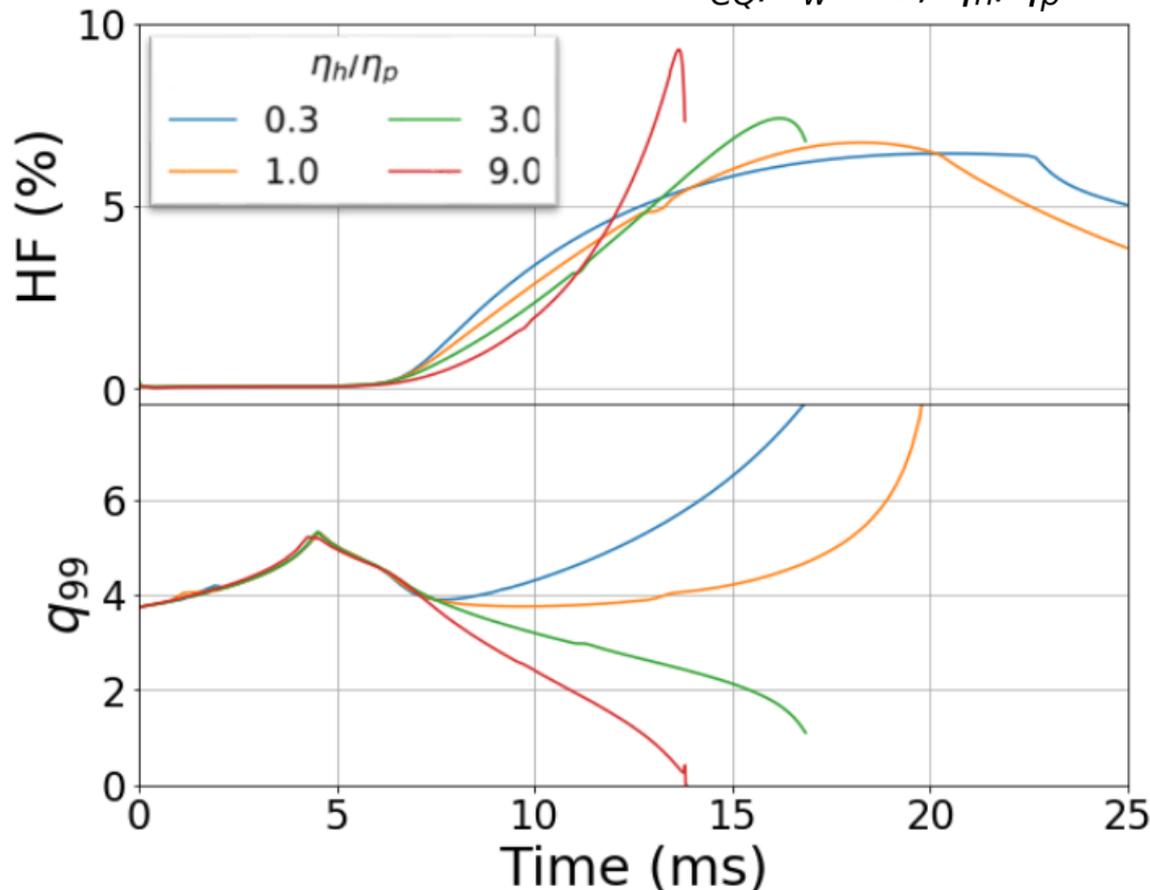
Halo currents also slow down vertical motion

$I_{halo, \phi}$ is also not linear in η_h
(self-regulating mechanism)

Understanding 2D halo currents at ITER (15 MA / 5.3T)

Ideal wall regime ($\tau_p \ll \tau_w$)

Halo resistivity scan $\tau_{CQ}/\tau_w = 0, \eta_h/\eta_p = 1$



q_a and $I_{halo,pol}$ depend strongly on η_h

The poloidal halo fraction (HF) is a factor 5-10 smaller than in the hot VDE regime

Prediction of the halo properties and B.C.s

Prediction of the halo properties and B.C.s

Currently working VDE model

- Temperature dependence for resistivity and parallel conductivity

$$\eta = \eta_0 \left(\frac{T}{T_0} \right)^{-3/2} \quad \kappa_{\parallel} = \kappa_{\parallel,0} \left(\frac{T}{T_0} \right)^{5/2}$$

- Ohmic heating term in energy equation
- Bohm's boundary condition ($v_{\parallel} = c_s$)
- Sheath heat flux B.C.

$$-(\kappa_{\perp} \nabla_{\perp} T + \kappa_{\parallel} \nabla_{\parallel} T) \cdot \mathbf{n} + nT \mathbf{v} \cdot \mathbf{n} = \gamma_{sh} nT c_s \frac{|\mathbf{B} \cdot \mathbf{n}|}{|B|}$$

Prediction of the halo properties and B.C.s

Missing ingredients of VDE model

- Neutrals, recycling and atomic processes (key for density evolution, **now $n_e(\psi)$**)
- Impurity evolution and radiation
- Limit on ion saturation current ($J \leq J_{sat}$), **(in progress)**

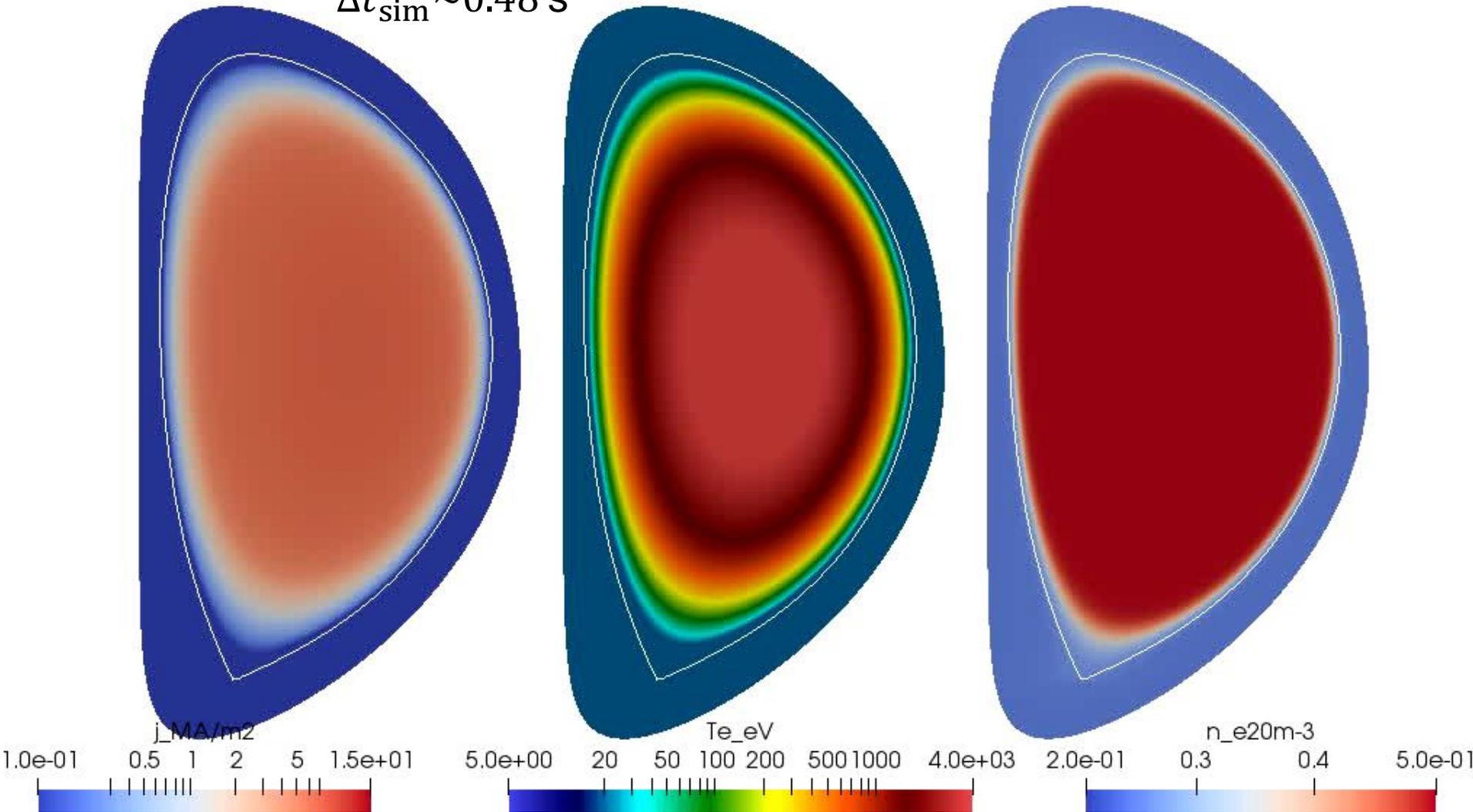
Prediction of the halo properties and B.C.s

Simulation setup

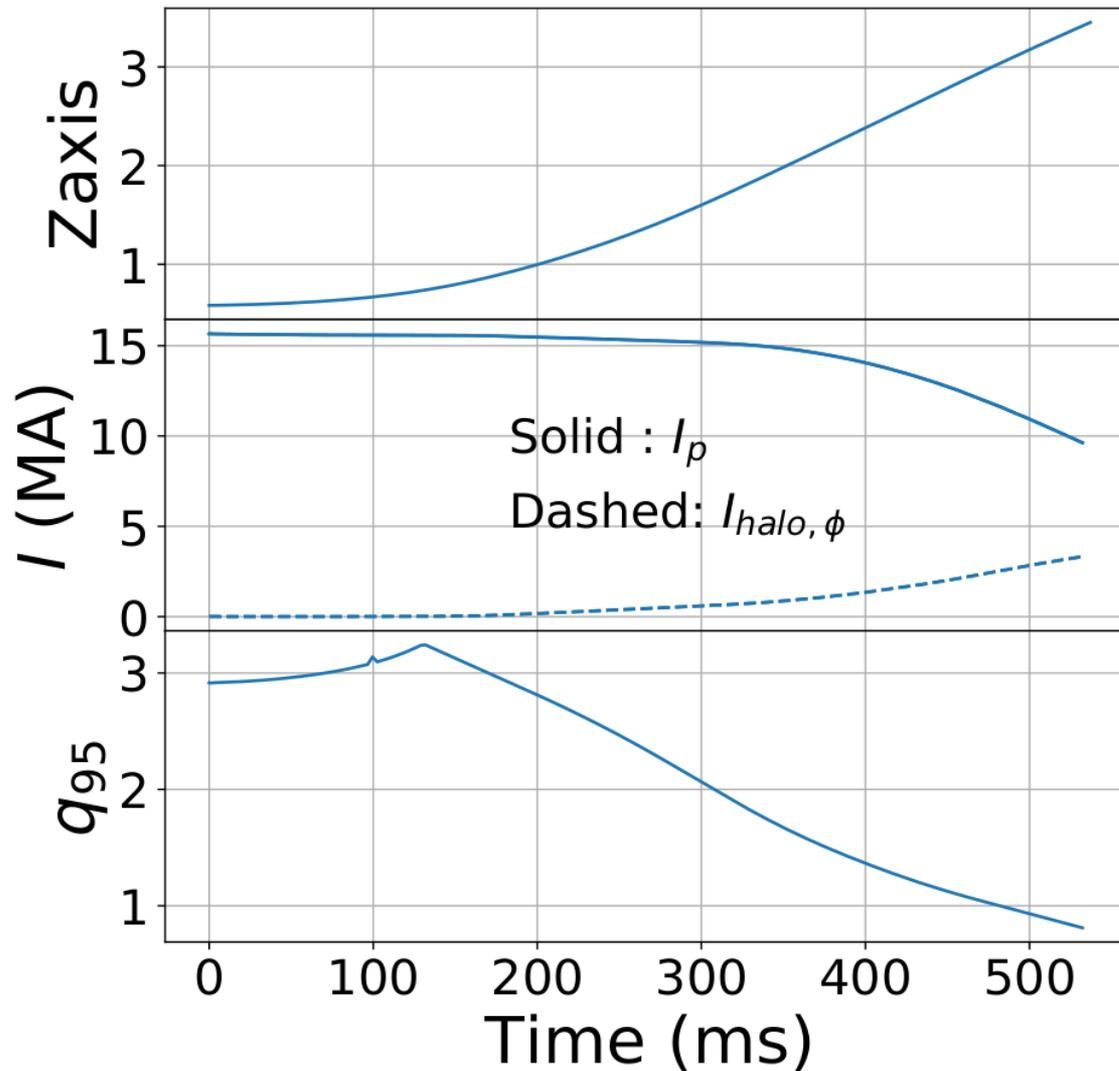
- Upward ITER VDE, 15 MA / 5.3 T
- Post-disruption equilibrium ($\beta_p = 0.05$)
- Flat J-profile after helicity mixing (from DINA)
- No radiation (ohmic heating re-heats the plasma)
- Realistic Spitzer and Braginskii values for η_0 and $\kappa_{\parallel 0}$
- $\kappa_{\perp} = 4 \text{ m}^2/\text{s}$, $\gamma_{sheath} = 8$, $T_e = T_i$, $\tau_w = 0.5 \text{ s}$

Prediction of the halo width and B.C.s

$\Delta t_{\text{sim}} \sim 0.48 \text{ s}$

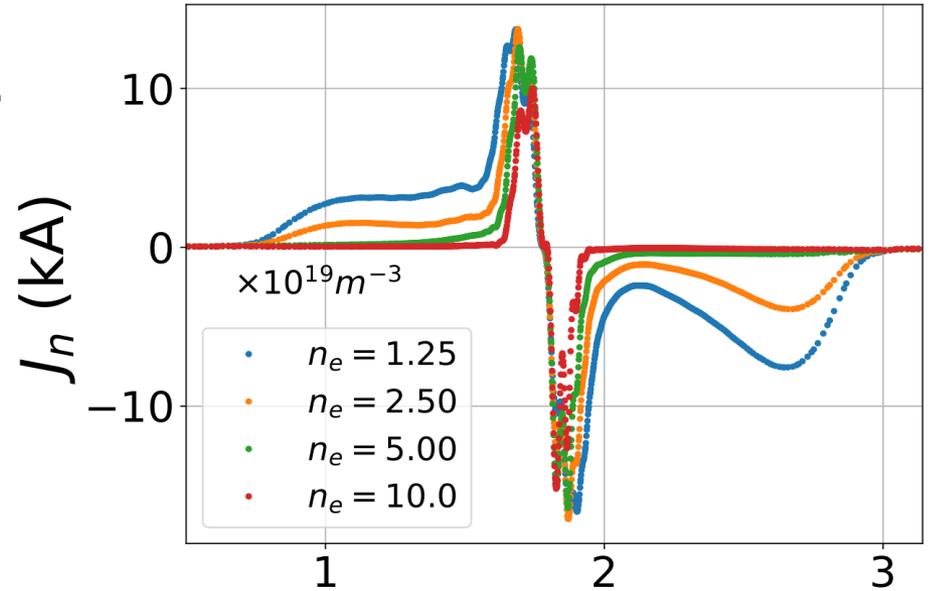
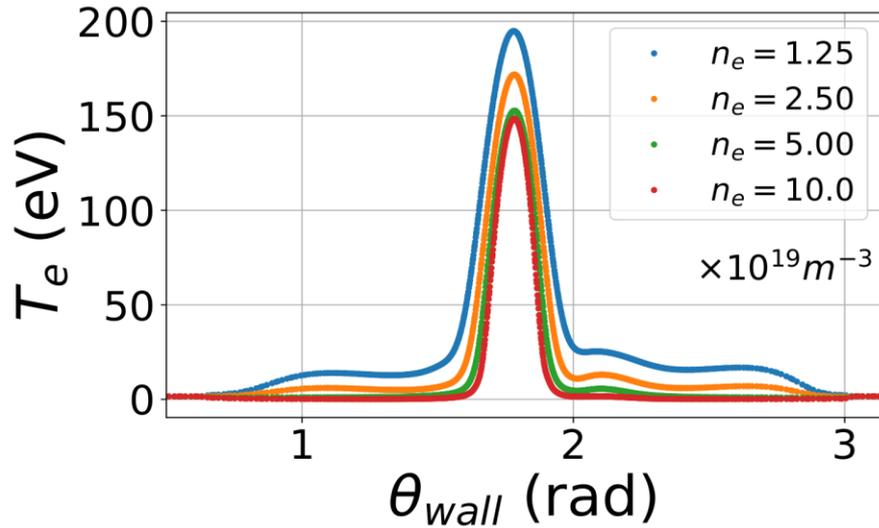


Prediction of the halo width and B.C.s

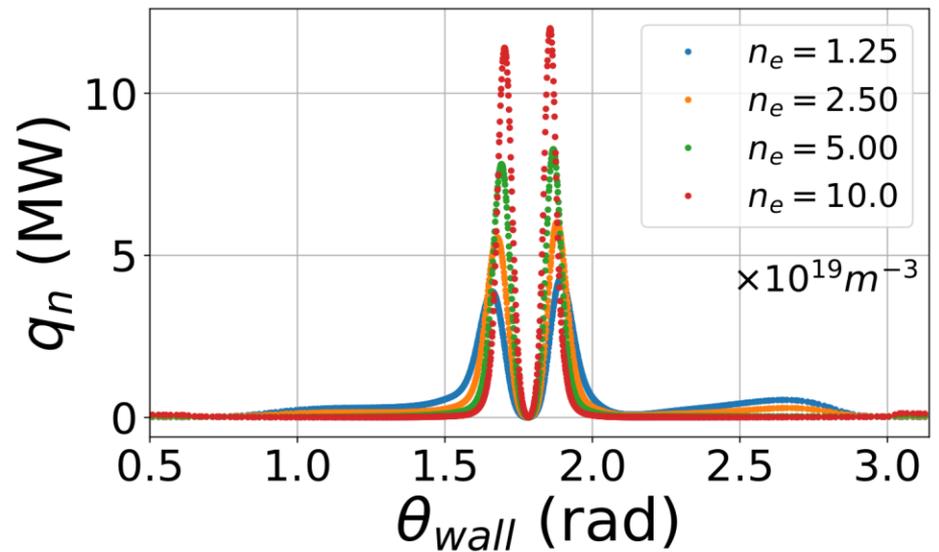
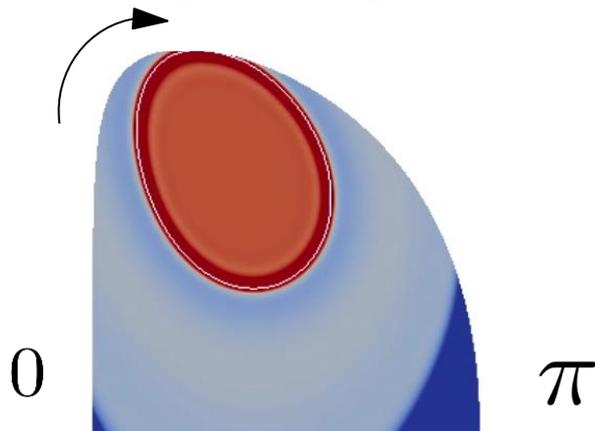


Prediction of the halo width and B.C.s

Density scan (Zaxis=2.0 m)



Poloidal angle along the wall



Conclusions and future work

JOREK / M3D-C1/ NIMROD benchmark: good agreement for 2D halo currents

2D VDEs ITER studies

- **Hot VDE limit** ($\tau_w \ll \tau_p$) **largest halo fractions** ($\text{HF}_{\text{max}} \sim 50\%$)
- **Ideal wall limit** ($\tau_p \ll \tau_w$) **smallest halo fractions** ($\text{HF}_{\text{max}} < 10\%$)
- **ITER mitigated disruptions** ($\text{HF}_{\text{max}} \sim 10 - 25\%$)
- Self-regulating mechanism for $I_{\text{halo},\phi}$, weak dependence on $\tau_h \propto w_h/\eta_h$
- Poloidal halo currents depend strongly on q_a ($I_{\text{halo,pol}} = I_{\text{halo},\phi}/q_a$), which decreases at shorter τ_h
- Maximum HF at ($\tau_w \ll \tau_h < \tau_p$) with small halo widths
- **Influence of initial l_i and core resistivity profile?**

Conclusions and future work

Prediction of halo width and temperature

➤ Current VDE model including

- Realistic resistivities and conductivities
- Energy balance in the halo: sheath losses and ohmic heating
- Sheath B.C.s
- Imposed density $n(\psi)$

➤ Still missing

- Density evolution with neutrals and atomic physics
- Impurity radiation
- Limit on current density (ion saturation current)

➤ Results for hot VDEs show

- Large halo widths (for J_{phi}) at low temperature