

IAEA-PPPL

Alpha particle dynamics and Alfvénic instabilities in ITER post-disruption plasmas

[Lier et al, NF 2021]

**Andrej Lier¹, Gergely Papp¹, Philipp Lauber¹,
O. Embreus², G. J. Wilkie³, S. Braun⁴**



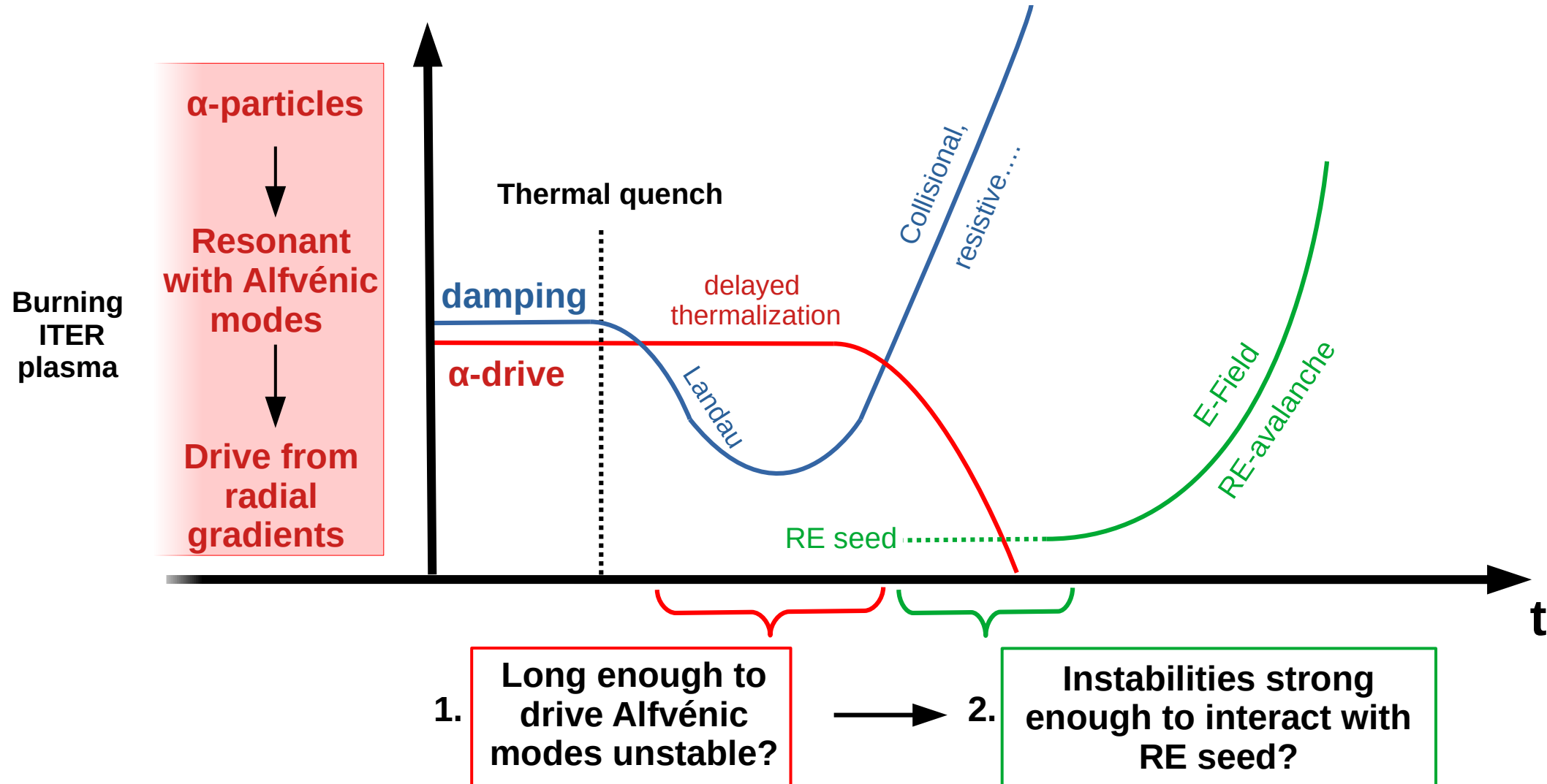
¹IPP, Garching

²Chalmers, Gothenburg

³PPPL, Princeton

⁴RWTH, Aachen

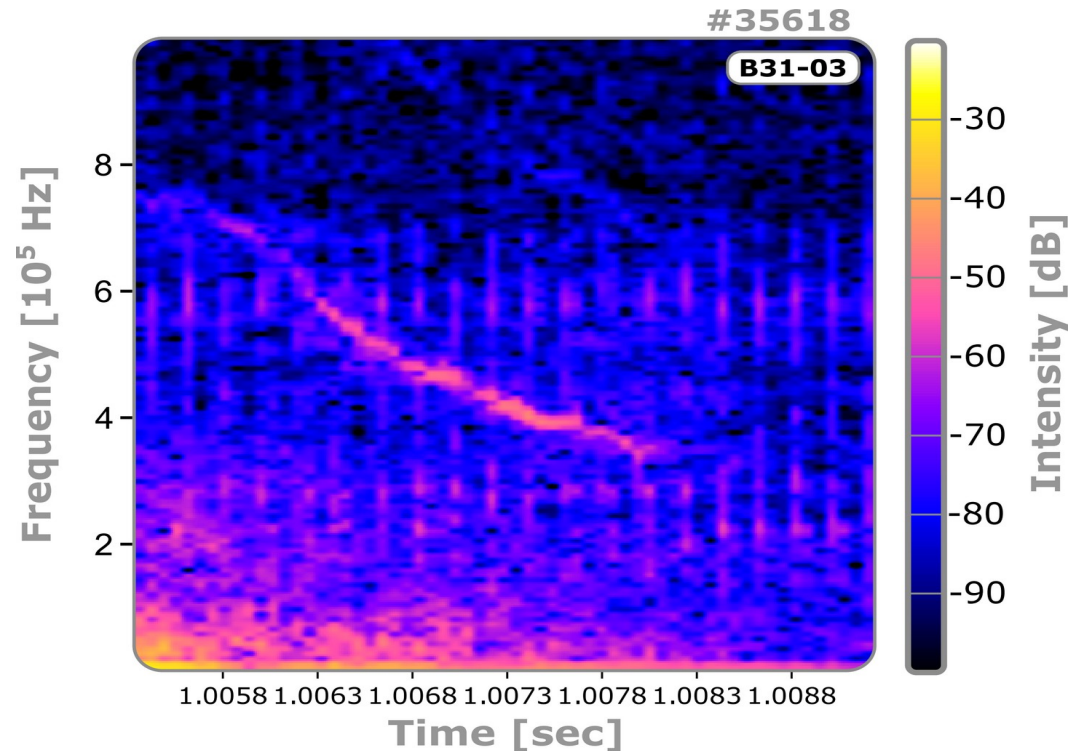
Physical Concept



Context – Experimental observations (1/2)

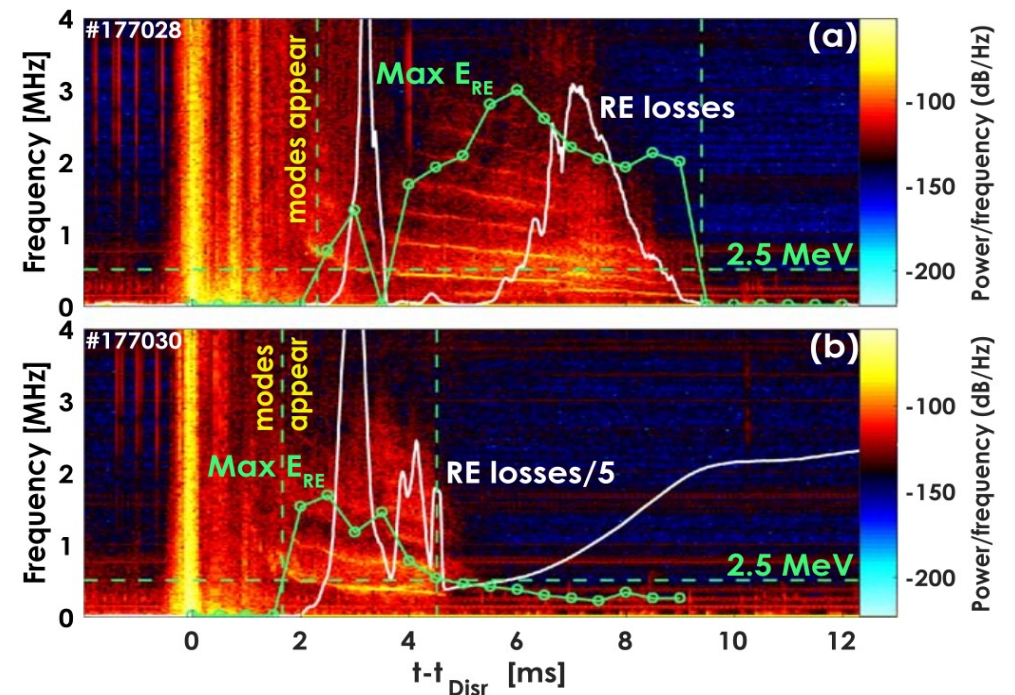
Present day tokamaks already observe
post-disruption **modes**

AUG



Current quench spectrogram
AUG#35618 [P. Heinrich]

DIII-D

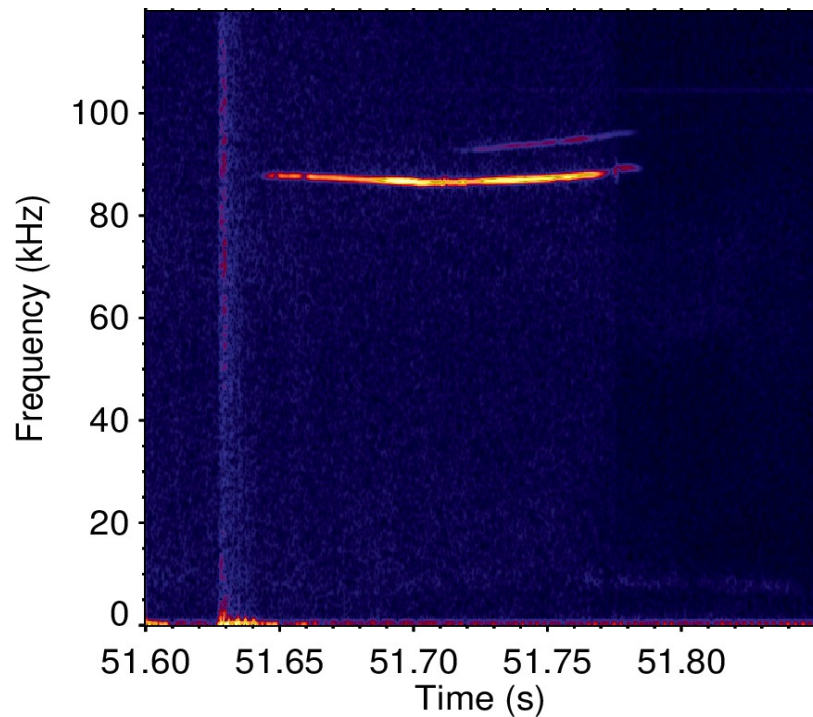


Spectrograms of DIII-D
[Lvovskiy, PPCF, 2018]

Context – Experimental observations (2/2)

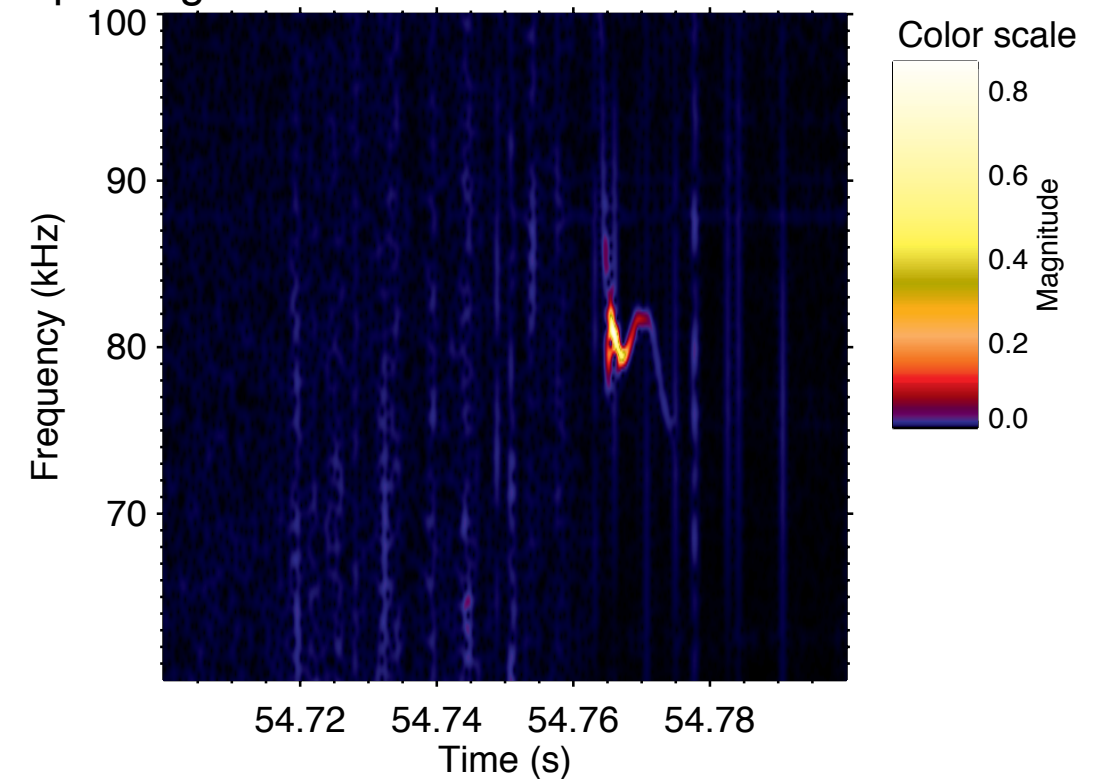
Present day tokamaks already observe
post-disruption **modes**

JET



Current quench spectrogram of JET
#89141. [S. Newton, P. Pölöskei]

Spectrogram of JET 42976-DI/C1F-CHAN8/131



Current quench spectrogram of JET
DT shot #42976. [S. Sharapov]

Alpha particles – velocity distribution

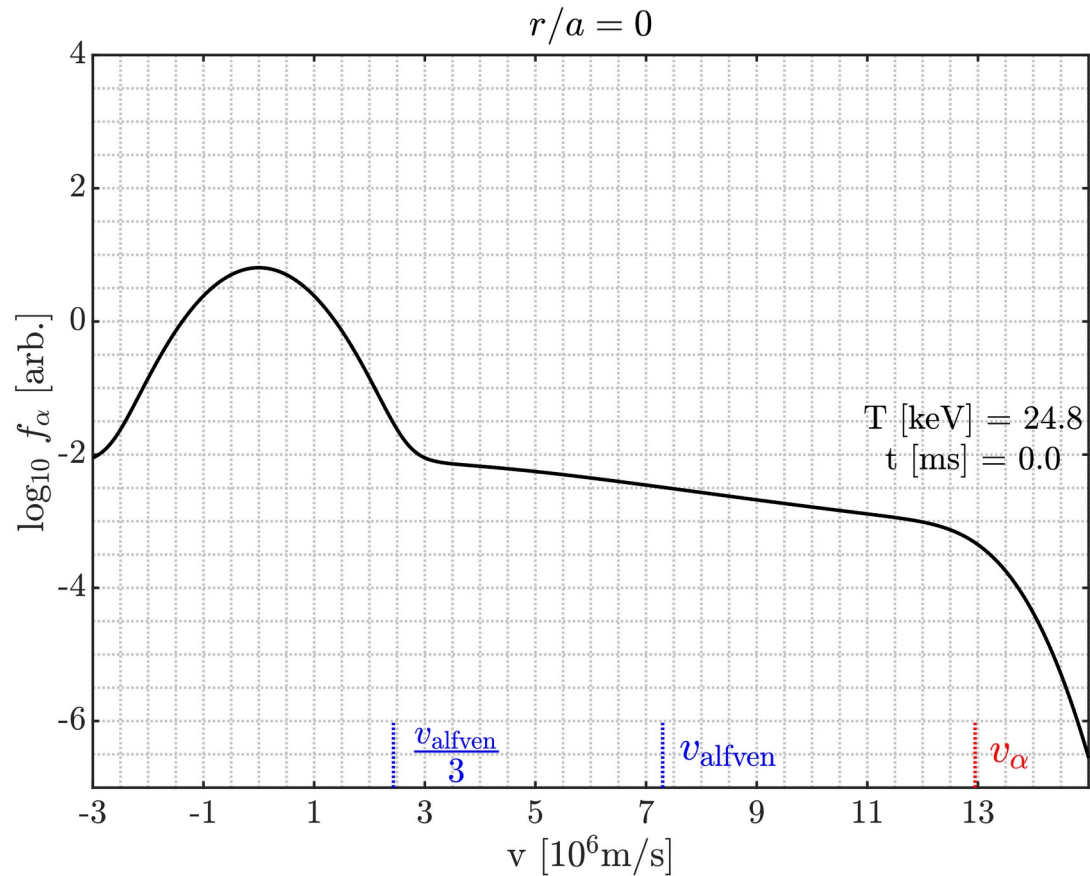


Fig. 1

CODION¹ simulation: Isotropic alpha particle velocity distribution for ITER 15MA scenario² #2

Fusion born alpha population is energetic by nature:

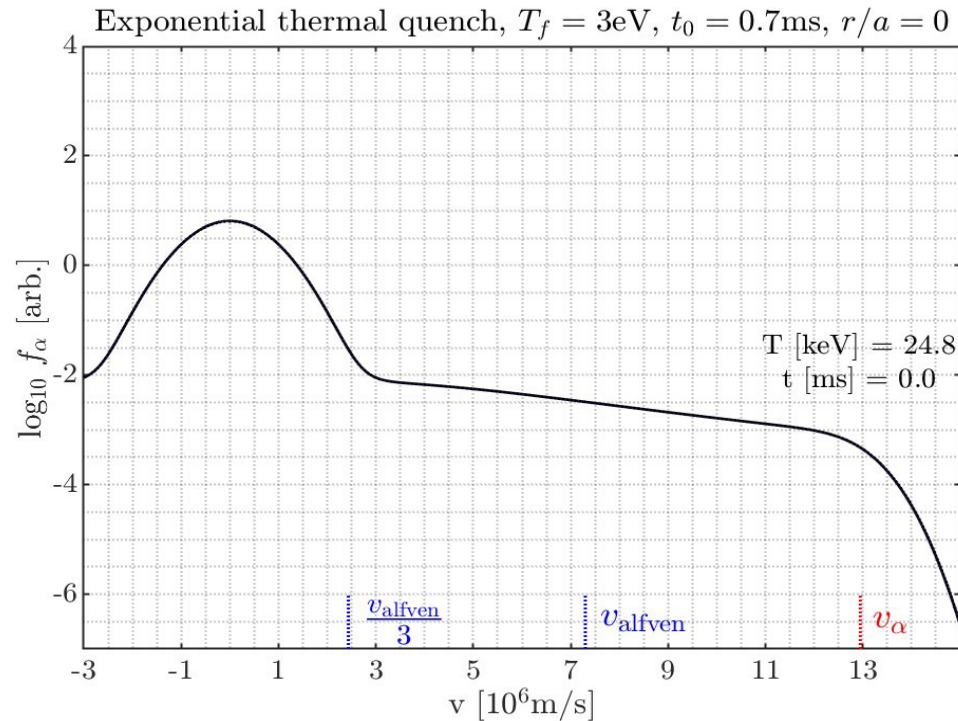
$$\alpha \text{ birth velocity } v_\alpha \approx 13 \cdot 10^6 \text{ m/s} > v_{\text{alfven}} = \frac{B}{\sqrt{\mu_0 m_i n_i}} \text{ Alfvén phase velocity}$$

Simulations³ show weakly unstable modes in ITER quiescent phase

Hypothesis
Post-disruption damping* drops faster than post-disruption alpha drive

*damping dominated by Landau damping $\sim \exp(-T)$

Alpha particles – delayed thermalization



Consider ‘worst case’, unmitigated disruption:

$$T(r, t) = T_f + [T(r, 0) - T_f] \exp(-t/t_0)$$

with Fokker-Planck solver **CODION**¹

Collisional cooling ineffective for **energetic particles**



Resonances possible far into the thermal quench

VIDEO

CODION¹ simulation: initial alpha distribution undergoing thermal quench.

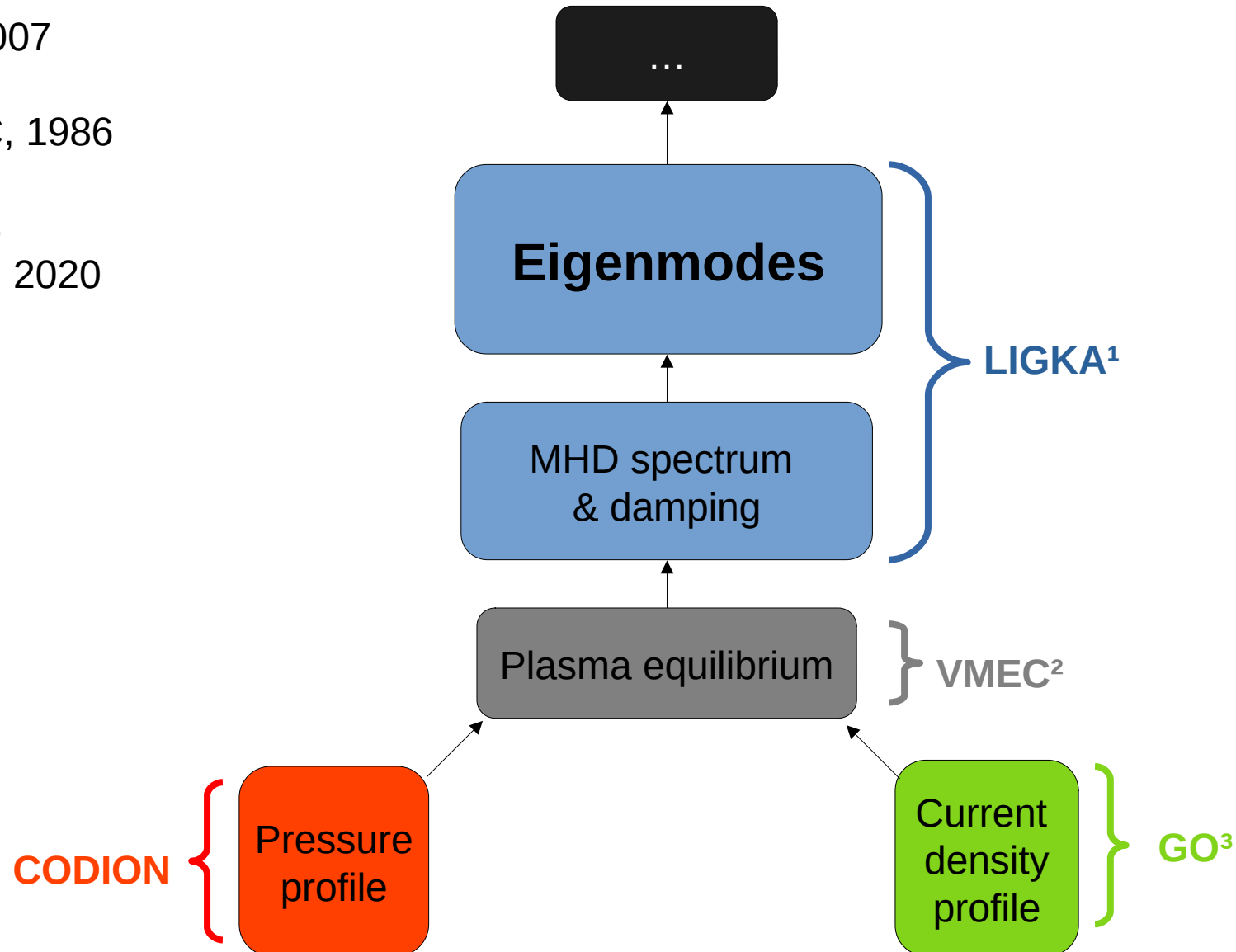
Workchain towards post-disruption Eigenmodes

¹Lauber, JCP, 2007

²Hirshman, CPC, 1986

³Papp, NF, 2013

³Vallhagen, JPP, 2020



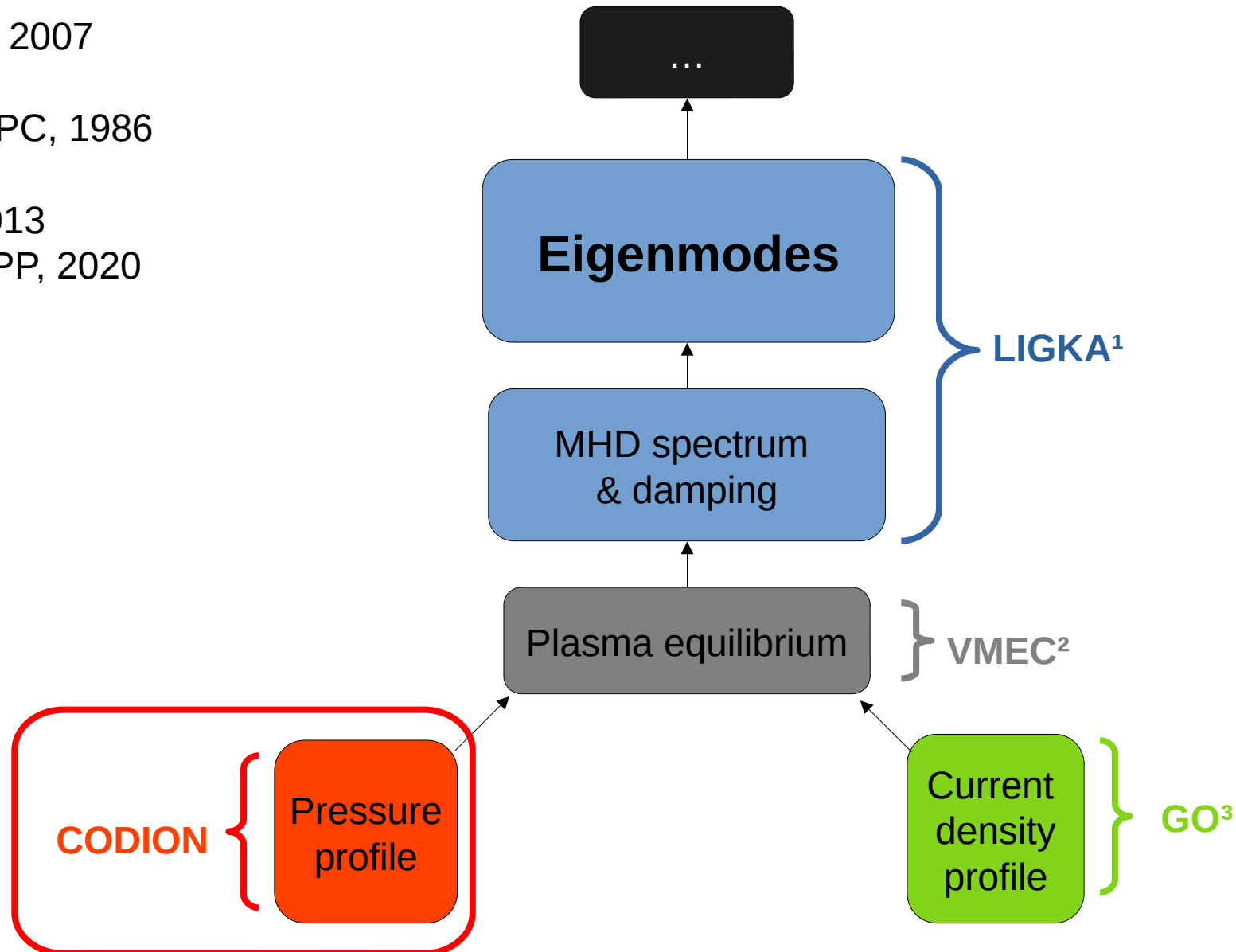
Workchain towards post-disruption Eigenmodes

¹Lauber, JCP, 2007

²Hirshman, CPC, 1986

³Papp, NF, 2013

³Vallhagen, JPP, 2020



Current density profile $j(r,t)$

GO code solves the induction equation in 1D

- Electric field diffusion
 - RE generation
- } $j(r,t)$

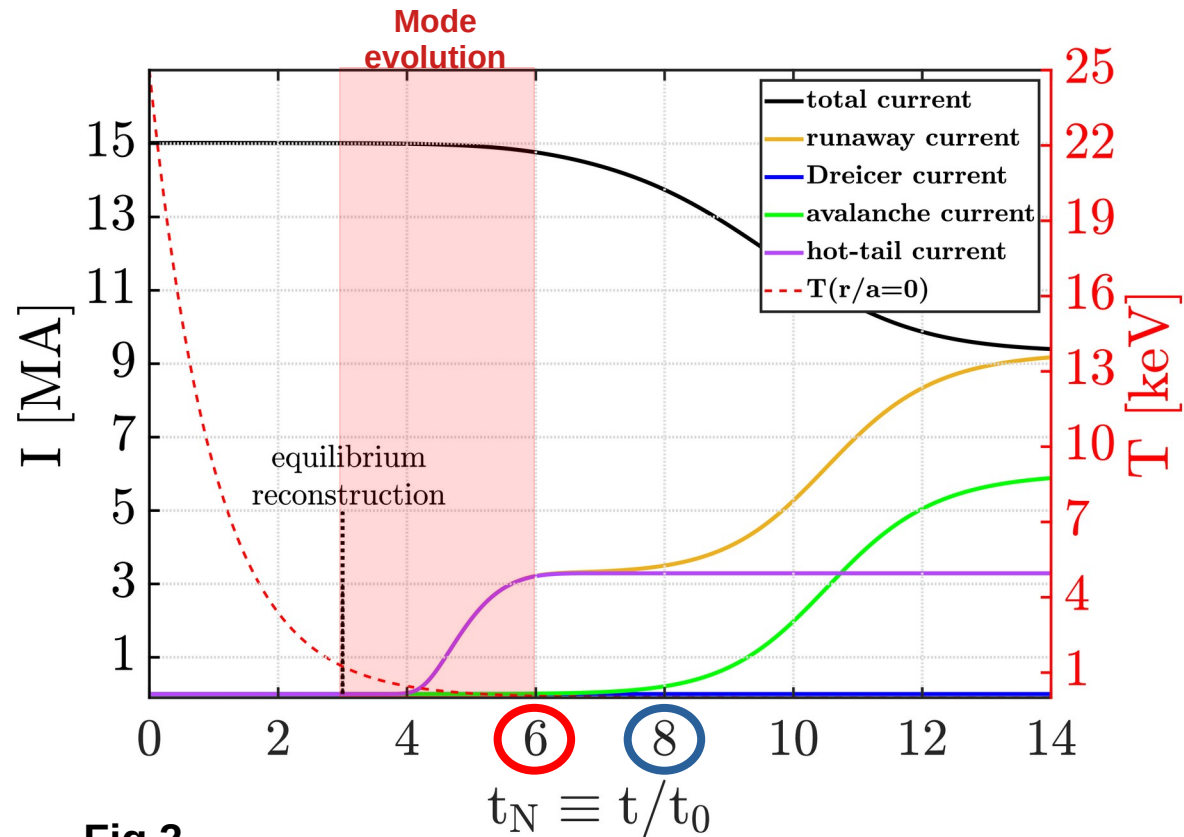


Fig.2

Currents of an unmitigated disruption identified by $T_f=3\text{eV}$ and $t_0=0.7\text{ms}$ and its background temperature

- ⑥ $T \sim 100 \text{ eV}$: Alphas thermalize
- ⑧ $T < 100 \text{ eV}$: Avalanching

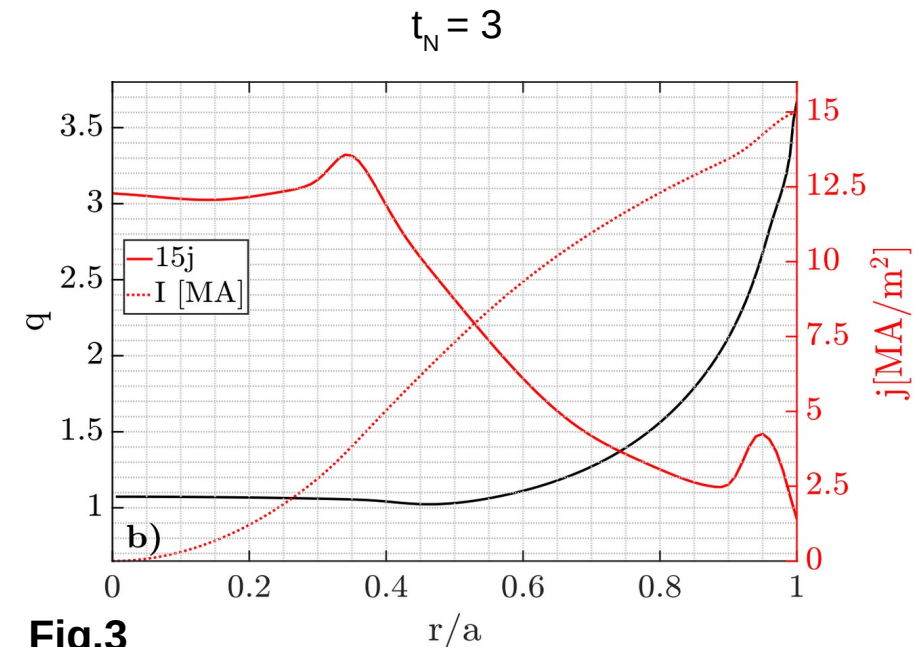
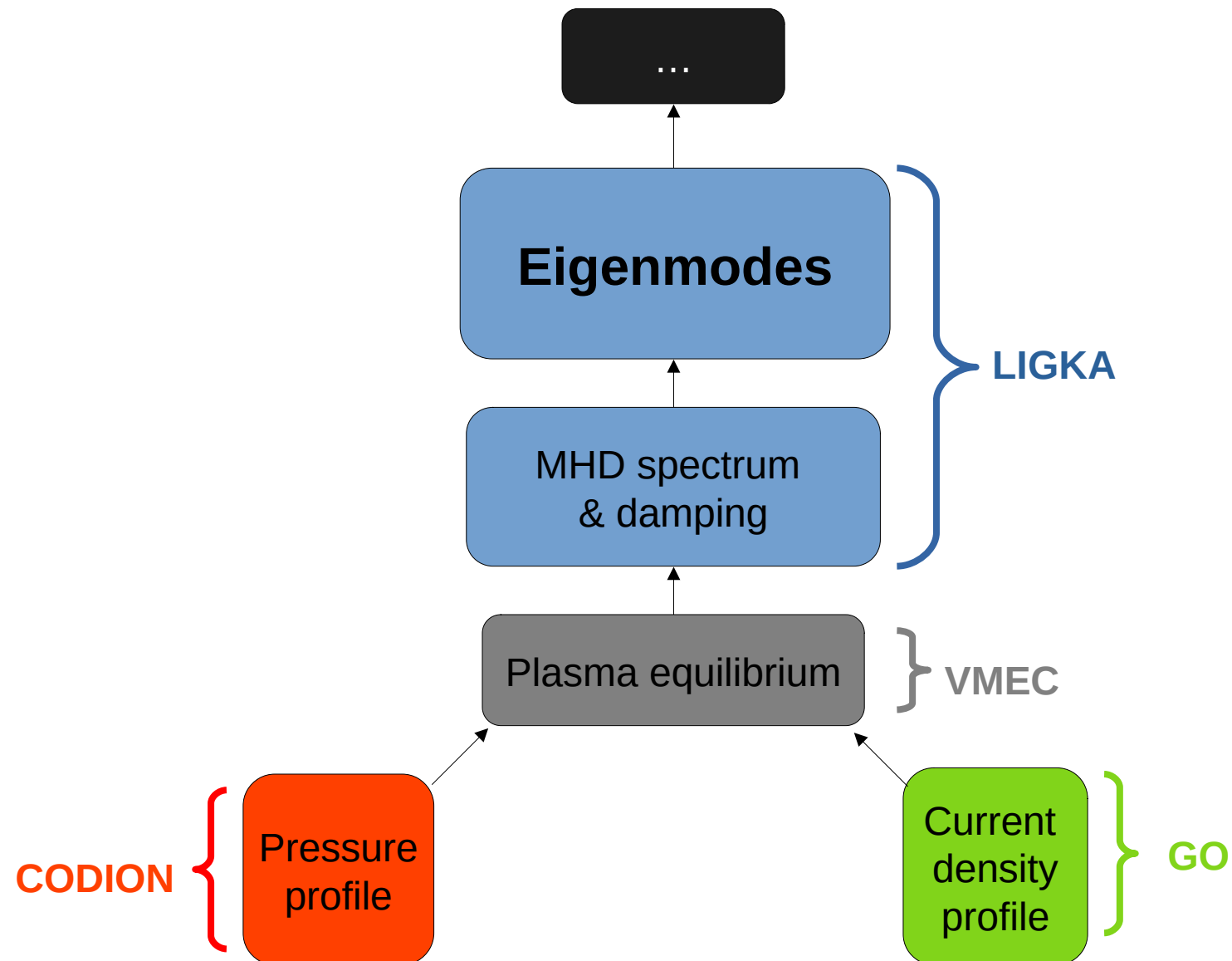


Fig.3

Resulting profiles of safety factor q , current density j and integrated current I .

Workchain towards post-disruption Eigenmodes



LIGKA tool employed:

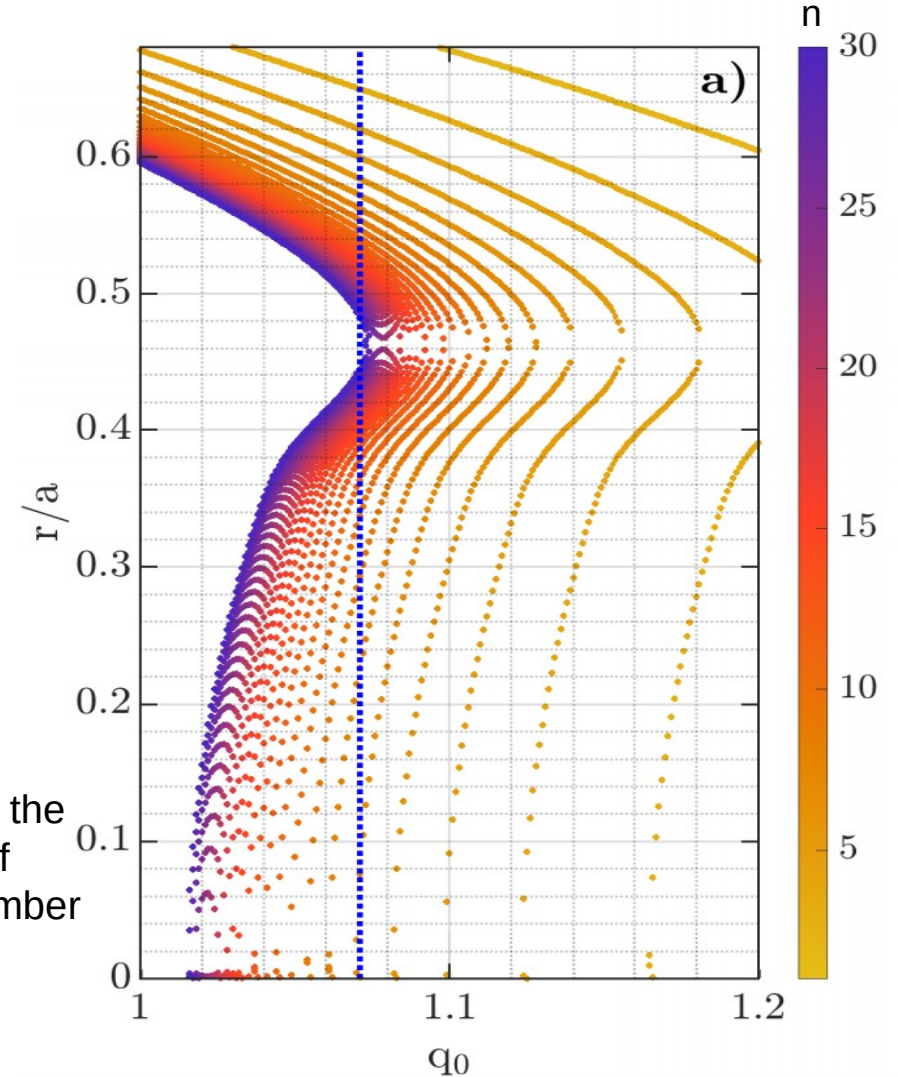
- found frequency gaps for TAEs (and BAEs) in the ideal MHD spectrum
- scan over absolute scaling of q-profile (fig. 6a) shows vast availability irrespective of q_0

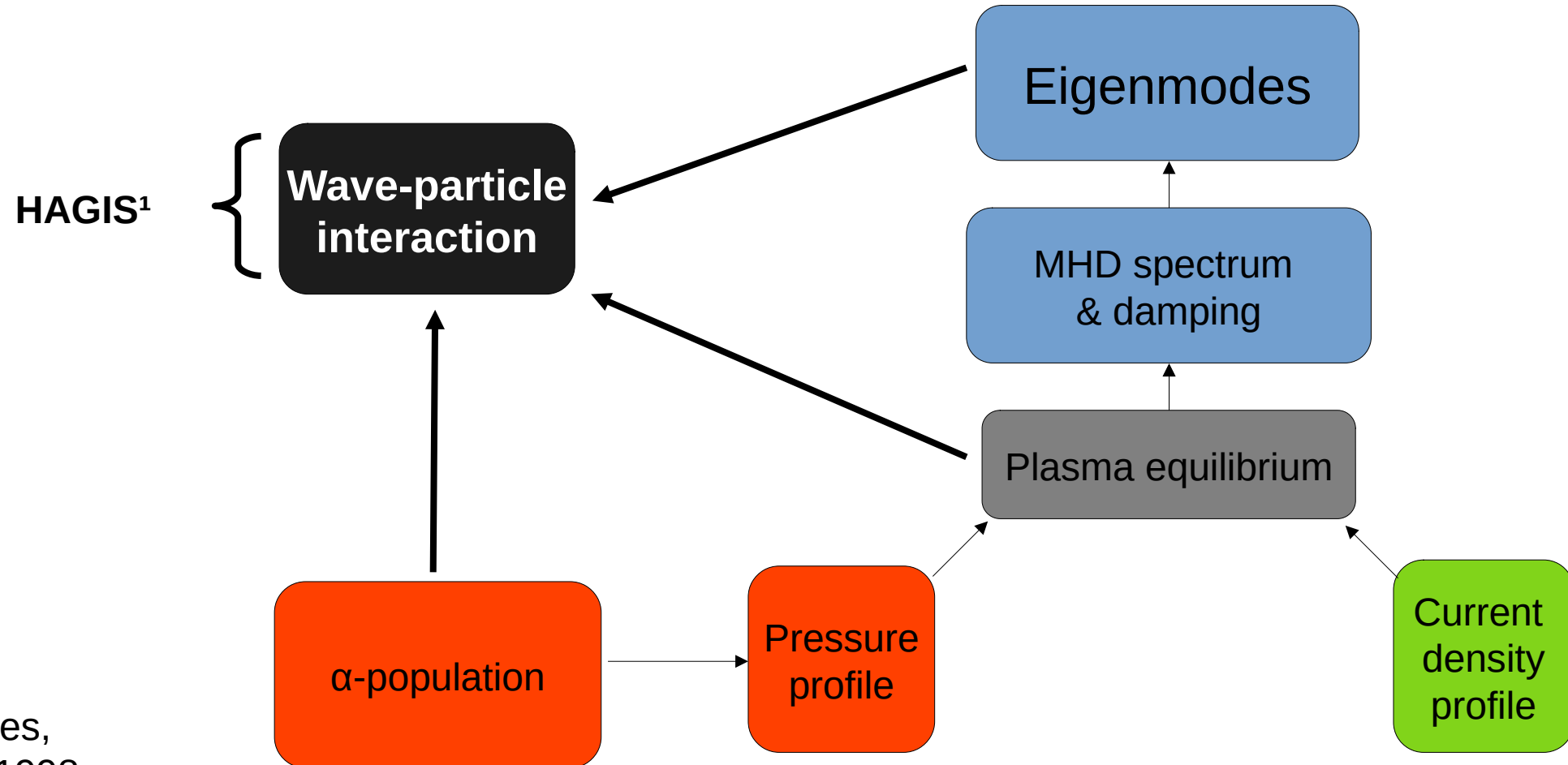
$q_0=1.07$ chosen

$n=[7-15, 22-26]$

$$m = \left[\left\{ (n-2) - (n+4) \right\}, \left\{ (n-2) - (n+6) \right\} \right]$$

Fig.4
Radial location of the frequency gaps of toroidal mode number $n=m$ TAEs as a function of q_0 .





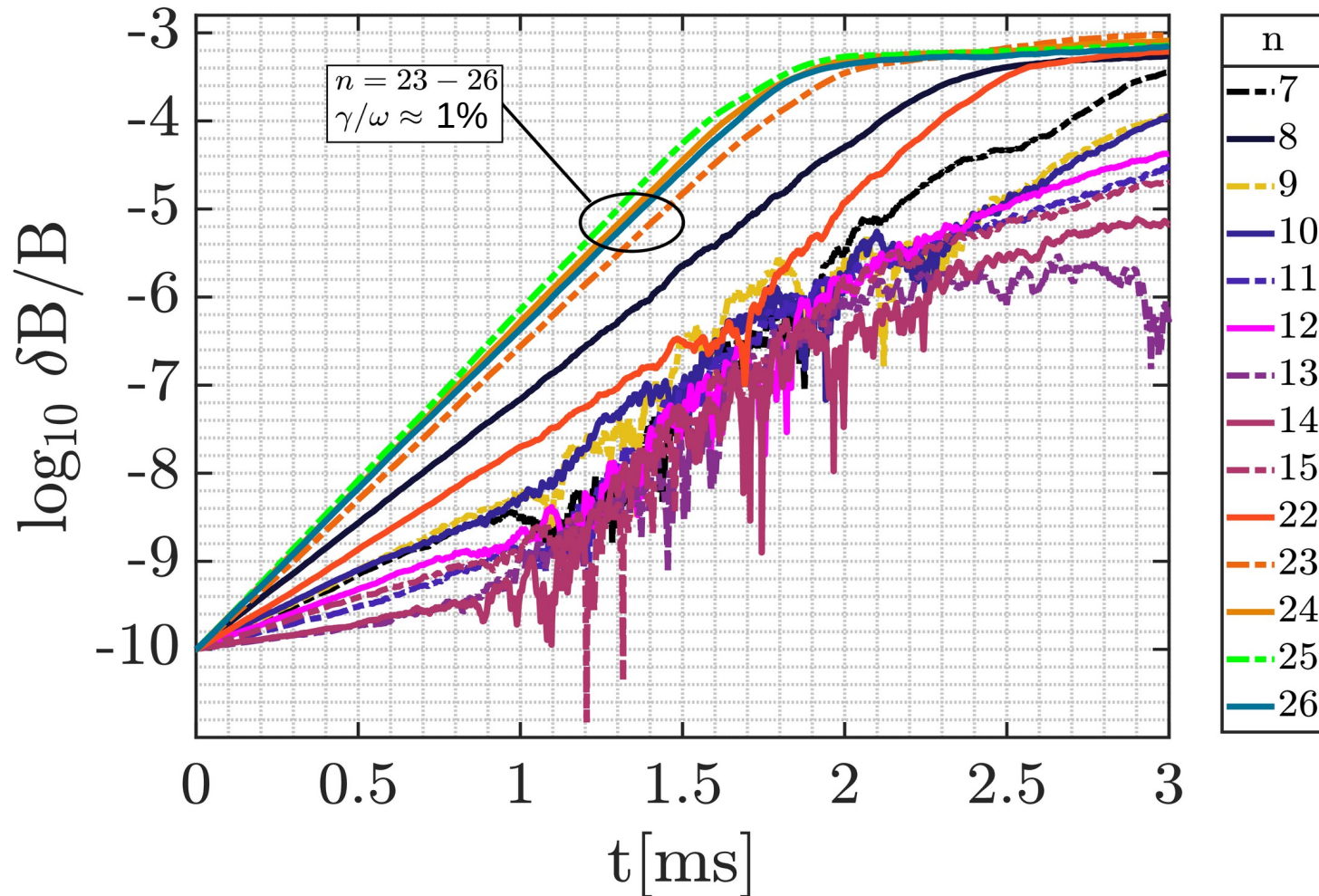
¹Pinches,
CPC, 1998

Active mode evolution

HAGIS calculates non-linear wave-particle interaction



evolves mode through EPs and redistributes EPs through modes

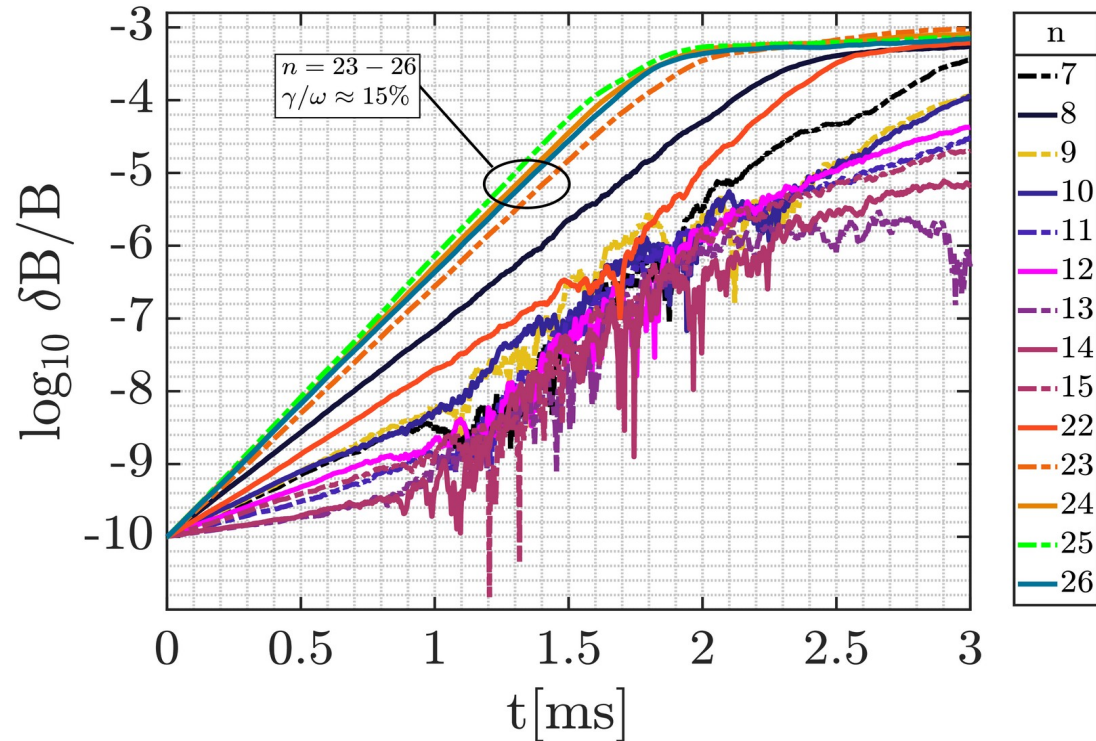


$\max(\delta B/B) \sim 0.1\%$
before RE avalanching

Fig.5

Evolution of mode amplitude $\delta B/B$ as caused by resonant interaction with f_{SD} in multi-mode simulation.

Mode effects on RE dynamics



$\max(\delta B/B) \sim 0.1\%$
before RE avalanching

→ effects on RE dynamics?

- HAGIS simulation indicating RE radial transport¹
- Study² found (stochastic) mag. Perturb ($\sim 0.05\%$) sufficient for RE avalanche suppression
- Further study: use ASCOT³ to determine transport coefficients as a function of (E, λ, r) for REs

¹[Lier, NF, 2021]

²[Svensson, JPP, 2021]

³[Schneider, NF, 2019]

Proof of principle stage¹:

- Unmitigated disruptions
- Perfect alpha particle confinement



Ongoing work:

- (simple) mitigated scenarios
- Alpha particle transport

¹[Lier et al, NF 2021]

1. Mitigated disruptions

To assess the increased parameter space:

Alpha distribution now calculated by analytical model (O. Embreus)

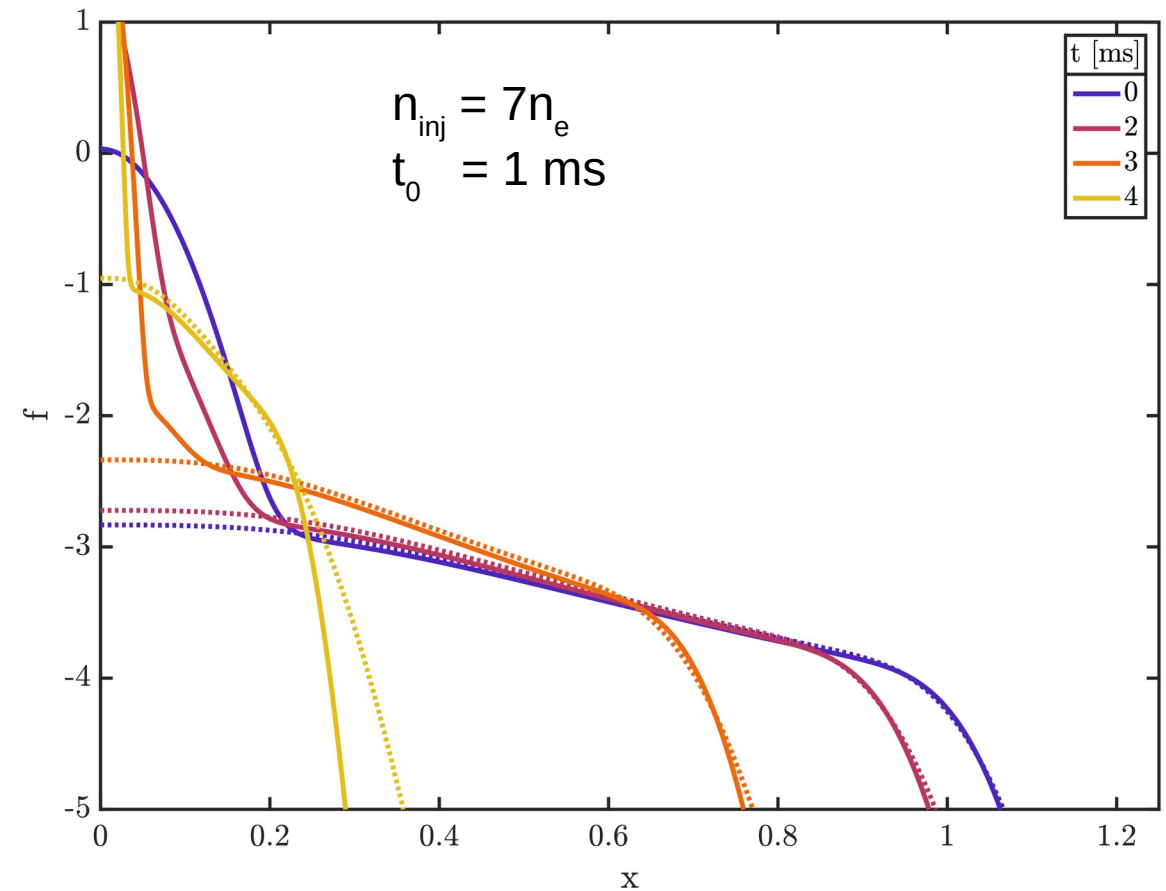
$$f_{\alpha}(r, t, v)$$



$$T(r, t), n_e(r, t)$$

Ion composition secondary to collision dynamics

→ use ions to tune v_A ?



Fig

Analytical (dashed) vs CODION (solid) results for alpha population in mitigated disruption

2. Alpha particle transport

Avoid MHD simulations for the entire parameter space:

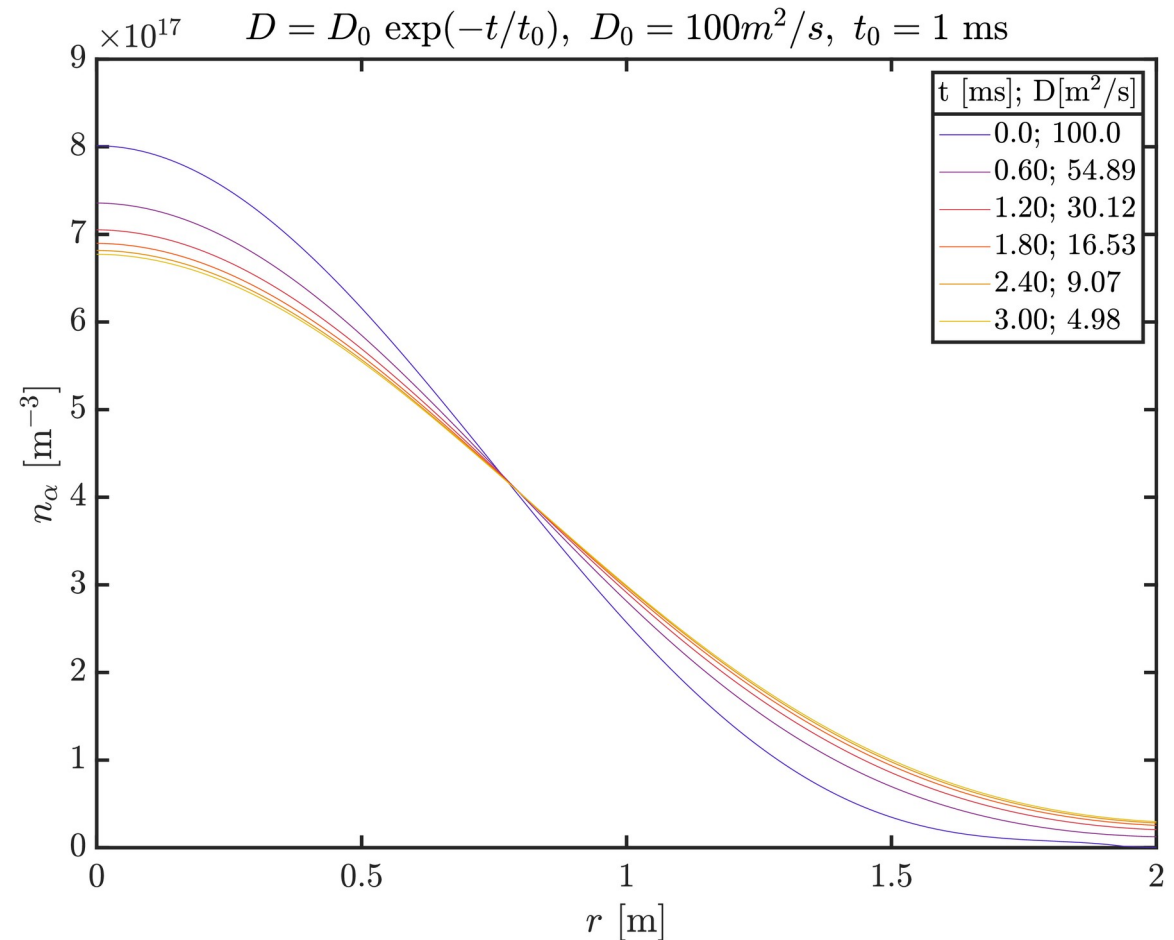
Solve diffusion equation

$$\frac{\partial u(r, t)}{\partial t} = D(t) \frac{\partial^2 u(r, t)}{\partial r^2}$$

with

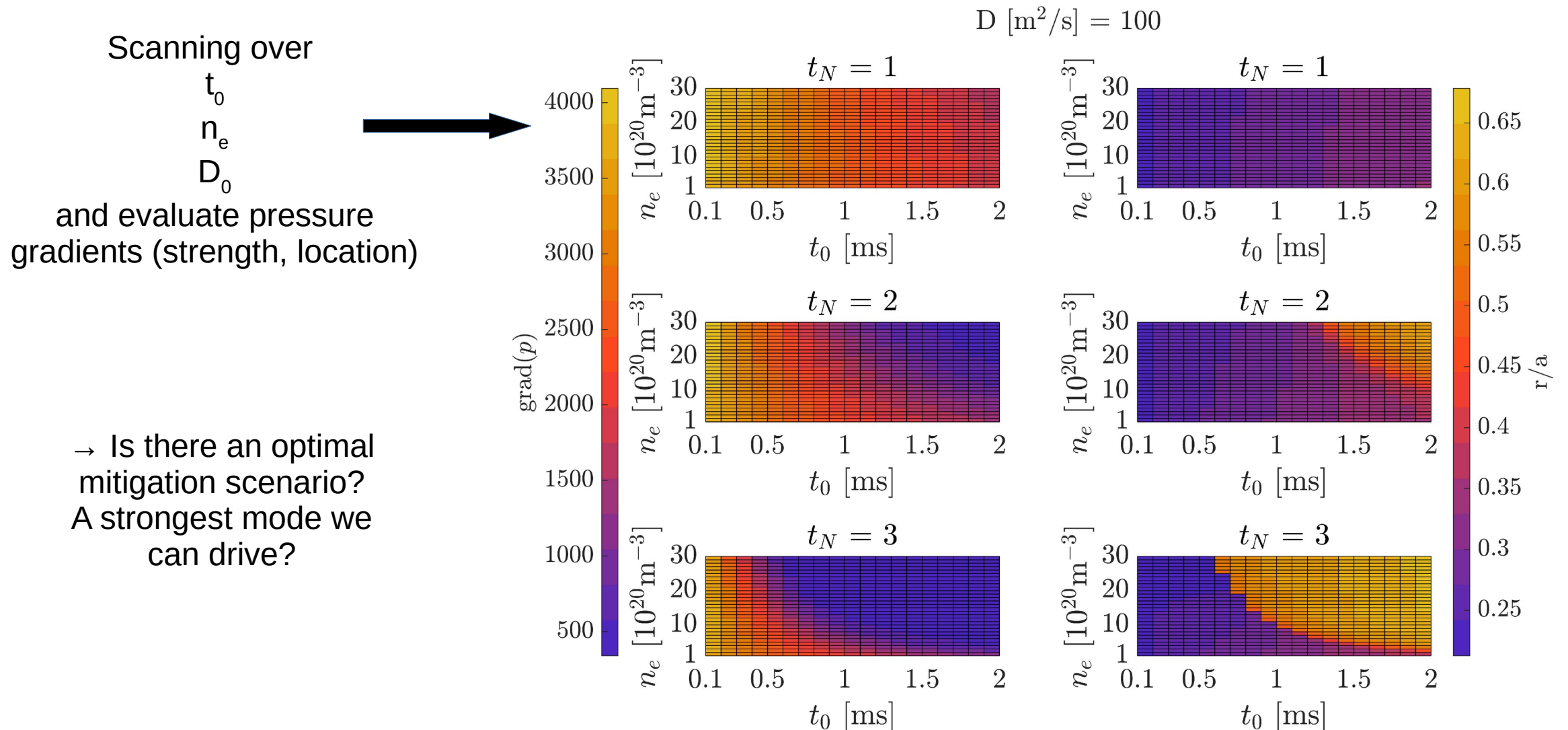
$$D(t) = D_0 \exp(-t/t_0)^1$$

and scan D_0 for stochasticity



¹[Linder, 2020, NF]

Ongoing work

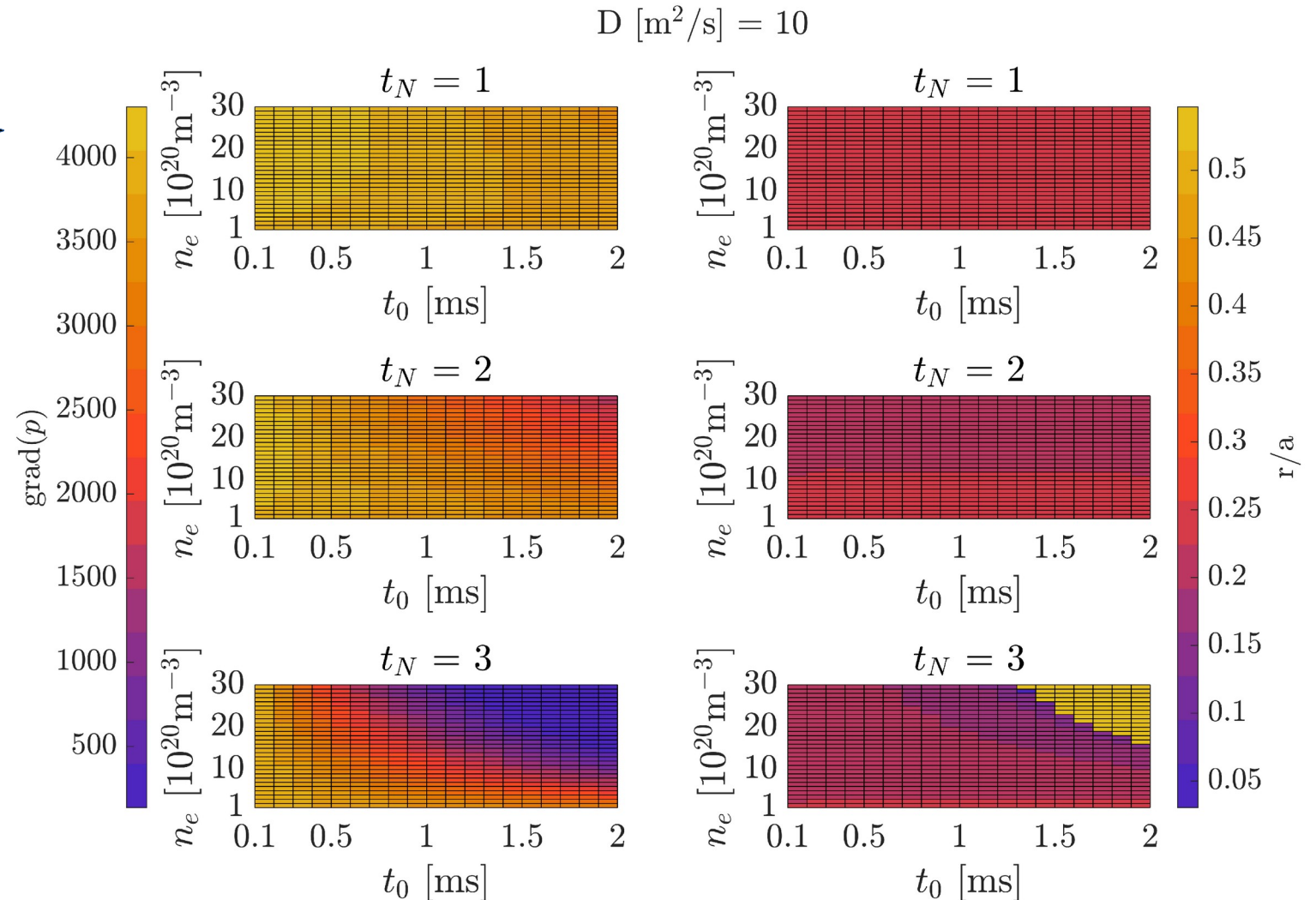


- Showed survivability of the energetic tail of a fusion-born alpha population far into the thermal quench
- The post-disruption MHD spectrum shows availability of a wide range of Toroidal Alfvén Eigenmodes which experience low damping
- Wave-particle interaction showed those TAEs to be driven unstable by the alpha population up to amplitudes of $\delta B/B = 0.1\%$
- The modes driven indicate a capability to enhance RE transport → effect on RE dynamics (suppression?)
- **Ongoing work**: Analytical model to scan big parameter space (mitigated disruptions) and search for optimum for this mechanic

Backup – Ongoing work

Scanning over
 t_0
 n_e
 D_0
and evaluate pressure
gradients (strength, location)

→ Is there an optimal
mitigation scenario?
A strongest mode we
can drive?



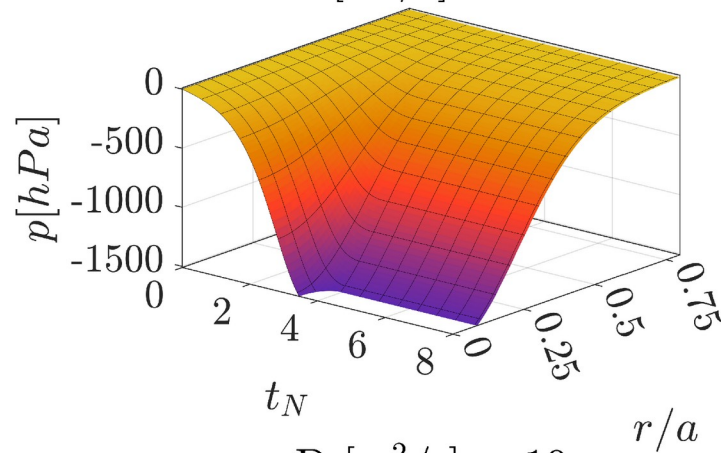
Backup – Ongoing work

Spatio-temporal pressure evolution

$$n_e [10^{20} \text{m}^{-3}] = 20; t_0 [\text{ms}] = 1$$

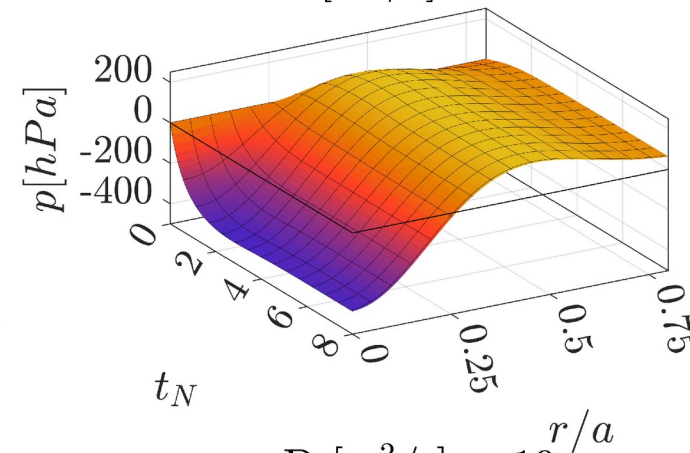
only slowing-down losses

$$D [\text{m}^2/\text{s}] = 100$$

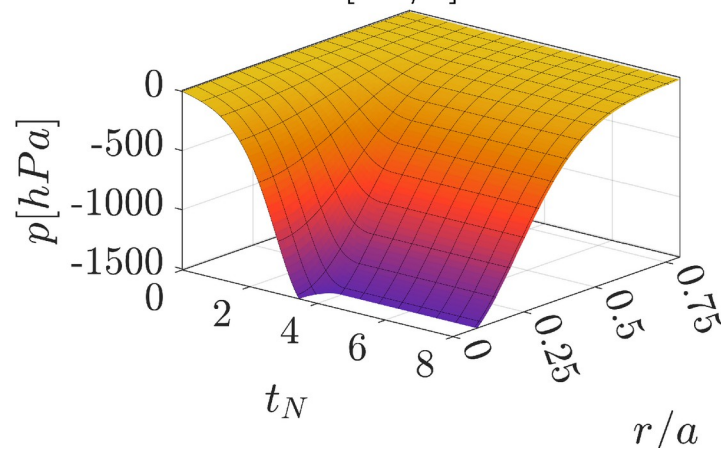


only diffusive process

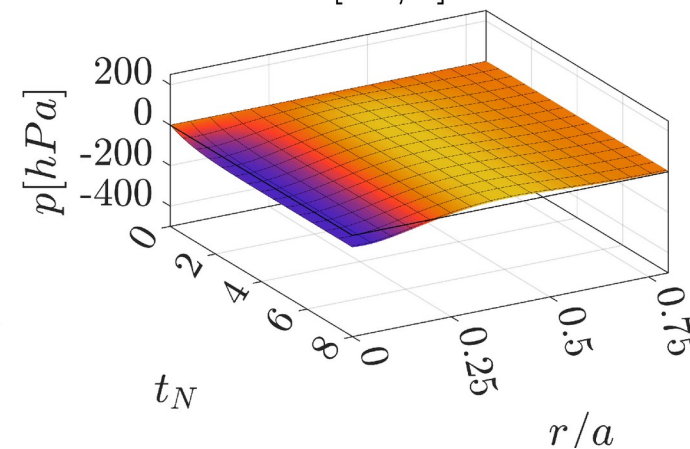
$$D [\text{m}^2/\text{s}] = 100$$



$$D [\text{m}^2/\text{s}] = 10$$



$$D [\text{m}^2/\text{s}] = 10$$



Backup - Mode effects on a RE seed

Which mode amplitudes to choose?

We are already $3t_N$ into the disruption and at $6t_N$ damping and avalanching becomes significant.

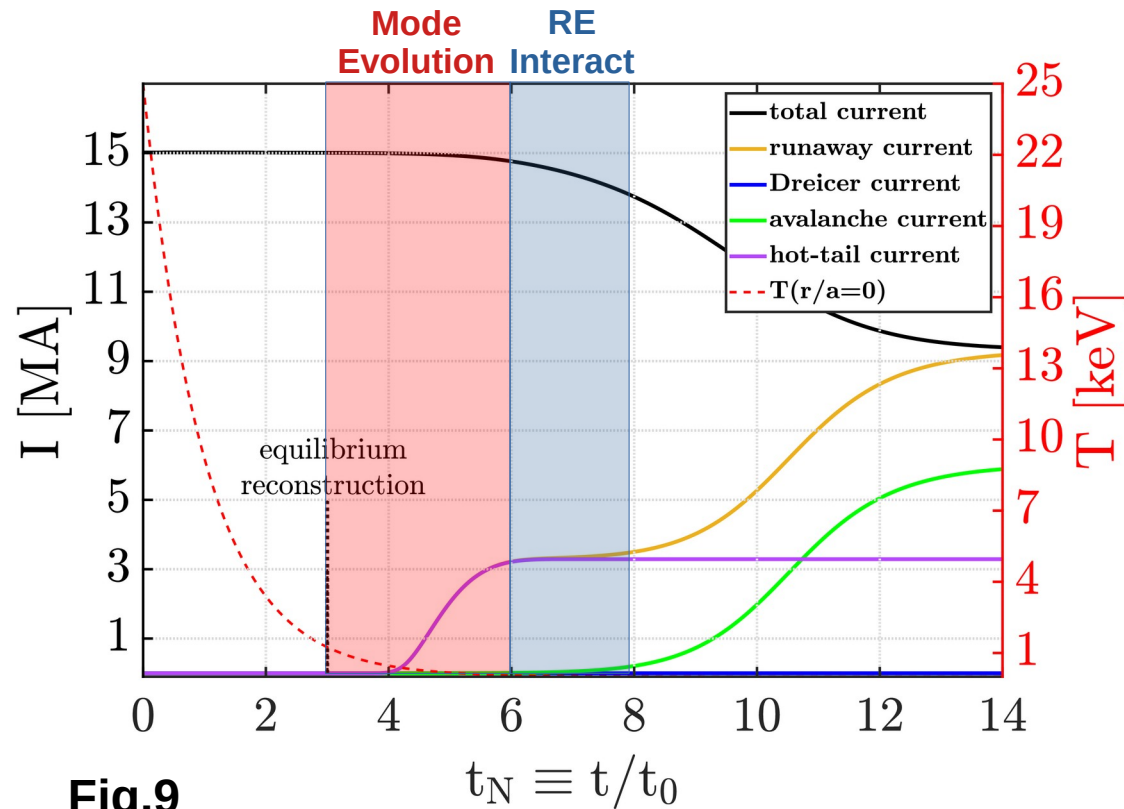
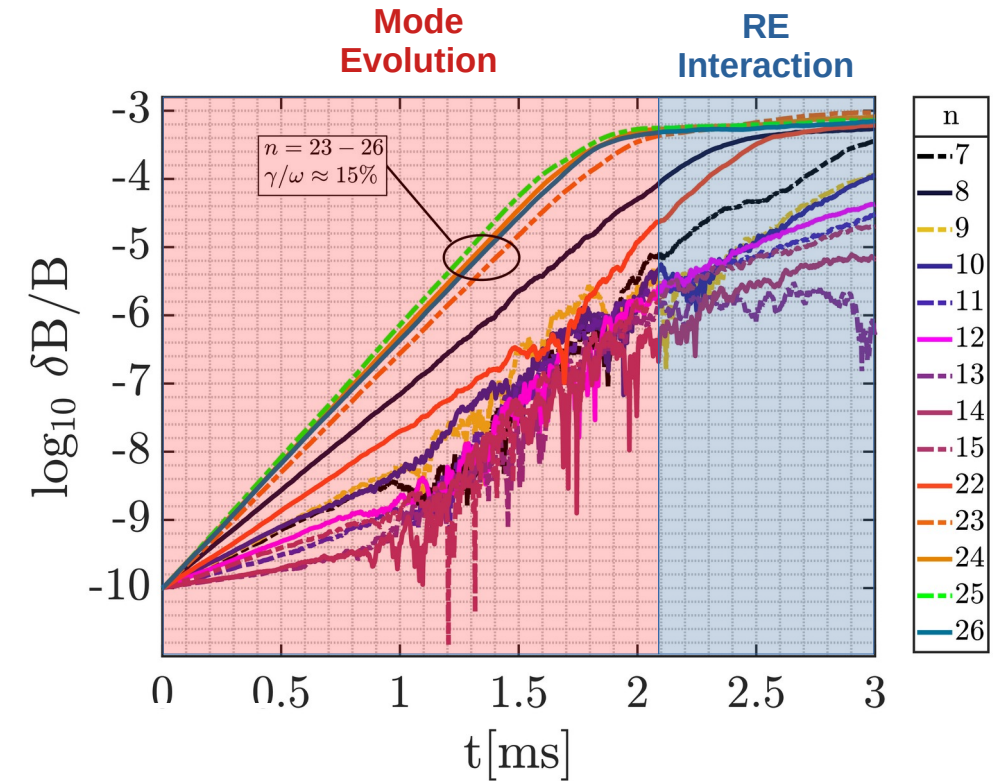


Fig.9

Currents of an unmitigated disruption identified by $T_f=3\text{eV}$ and $t_0=0.7\text{ms}$ and its background temperature



Backup - Mode effects on a RE seed

RE seed initialized:

$$\left. \begin{aligned} E_{kin} &= [10 \text{ keV} - 30 \text{ MeV}] \\ \lambda \equiv v_{\parallel} / v &= [0 - 1] \\ r/a &= [0.05 - 0.45] \end{aligned} \right\} \begin{aligned} &\text{Each triple combination} \\ &\text{represented by 25 REs,} \\ &\text{uniformly distributed} \\ &\text{along torus.} \\ &\Sigma \# \text{REs} = 10000 \end{aligned}$$

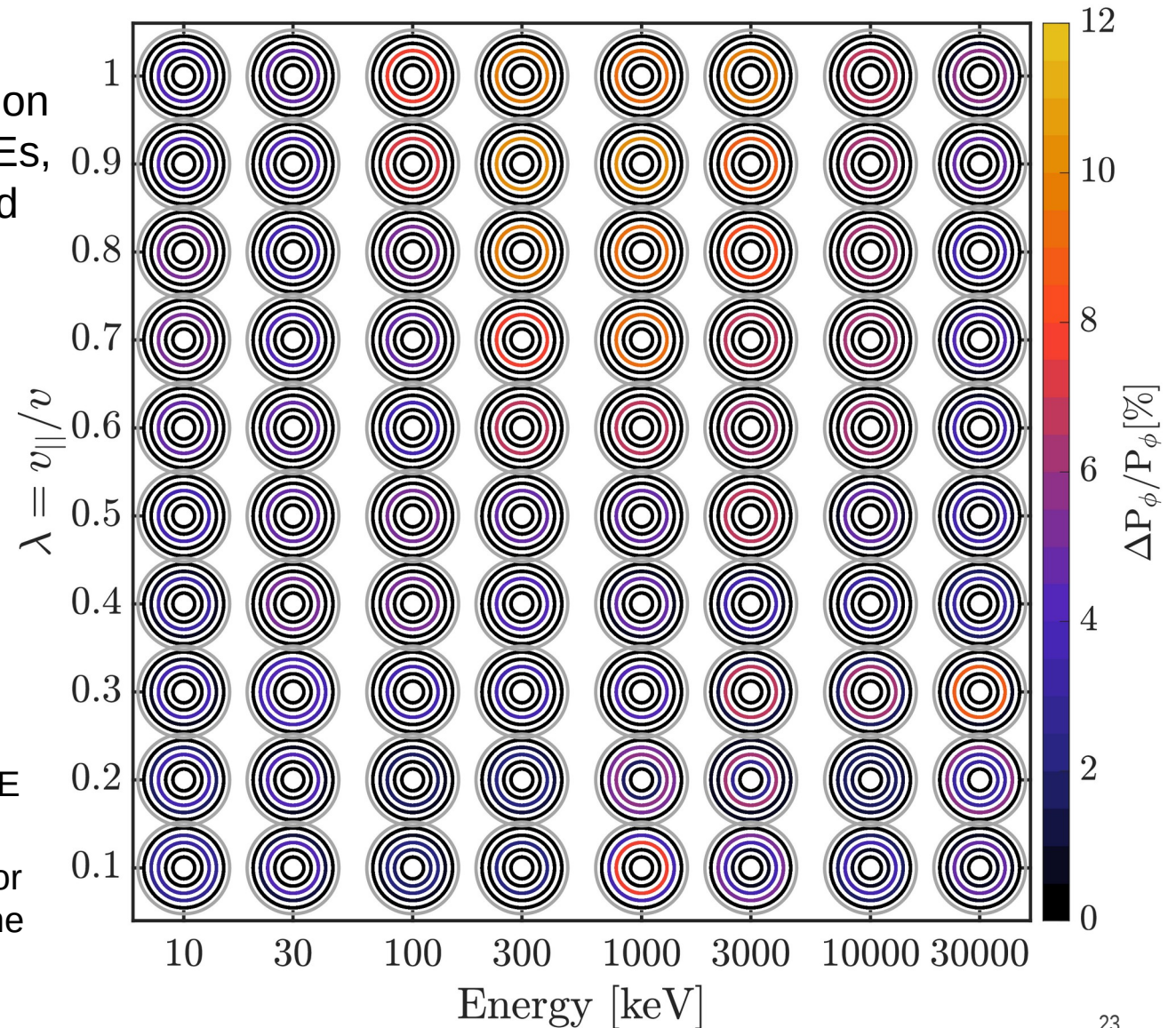
Throughout interaction measure changes to the
toroidal angular momentum

$$P_{\phi}(p_{\parallel}, \Psi_p)$$

as indication to changes of radial position.

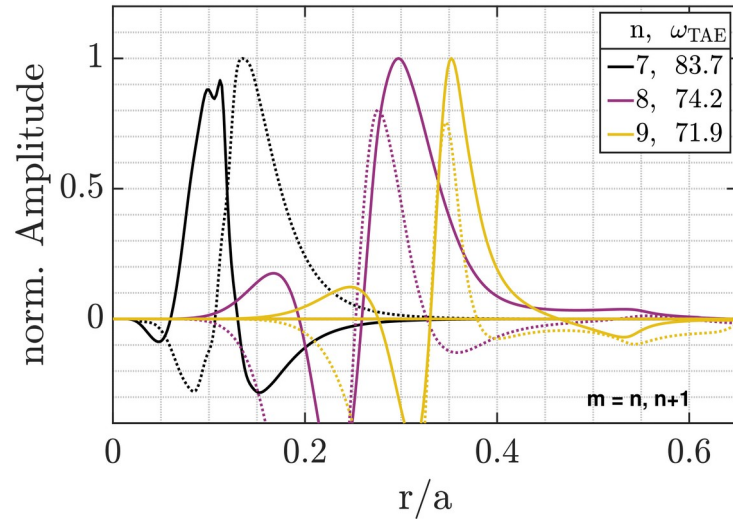
Fig.10

Ensemble-averaged change to P_{ϕ} of the RE seed as caused by the TAEs. X and y-axis show initial RE attributes, color indicates change after $t_N=2$. The Radii of the circles are the initial radial position of the particle in $r/a=[0.05-0.45]$ in steps of 0.1.



Backup – mode structures

(selected) low n branch



Mode structures for $q_0 = 1.07$

→ total structure is **core localized** along flat shear

α-pressure
core-peaked

RE-generation
core-localized

Damping strengths γ/ω [s^{-1}/s^{-1}]:

- Landau+radiative (LIGKA) $\sim 0.1\%$ ($t_N = 3$)
- Fluid damping (CASTOR¹) $\sim 1\%$ ($t_N = 8$)
- Collisional damping² $\sim 1\%$ ($t_N = 6$)

high n branch

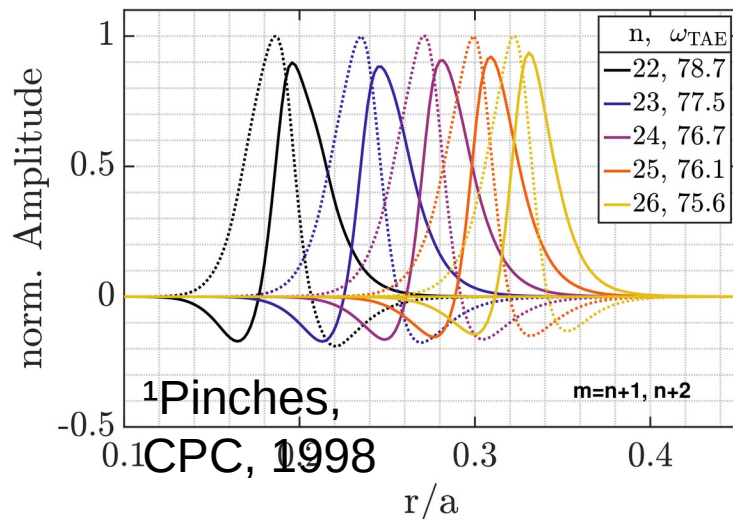


Fig.6

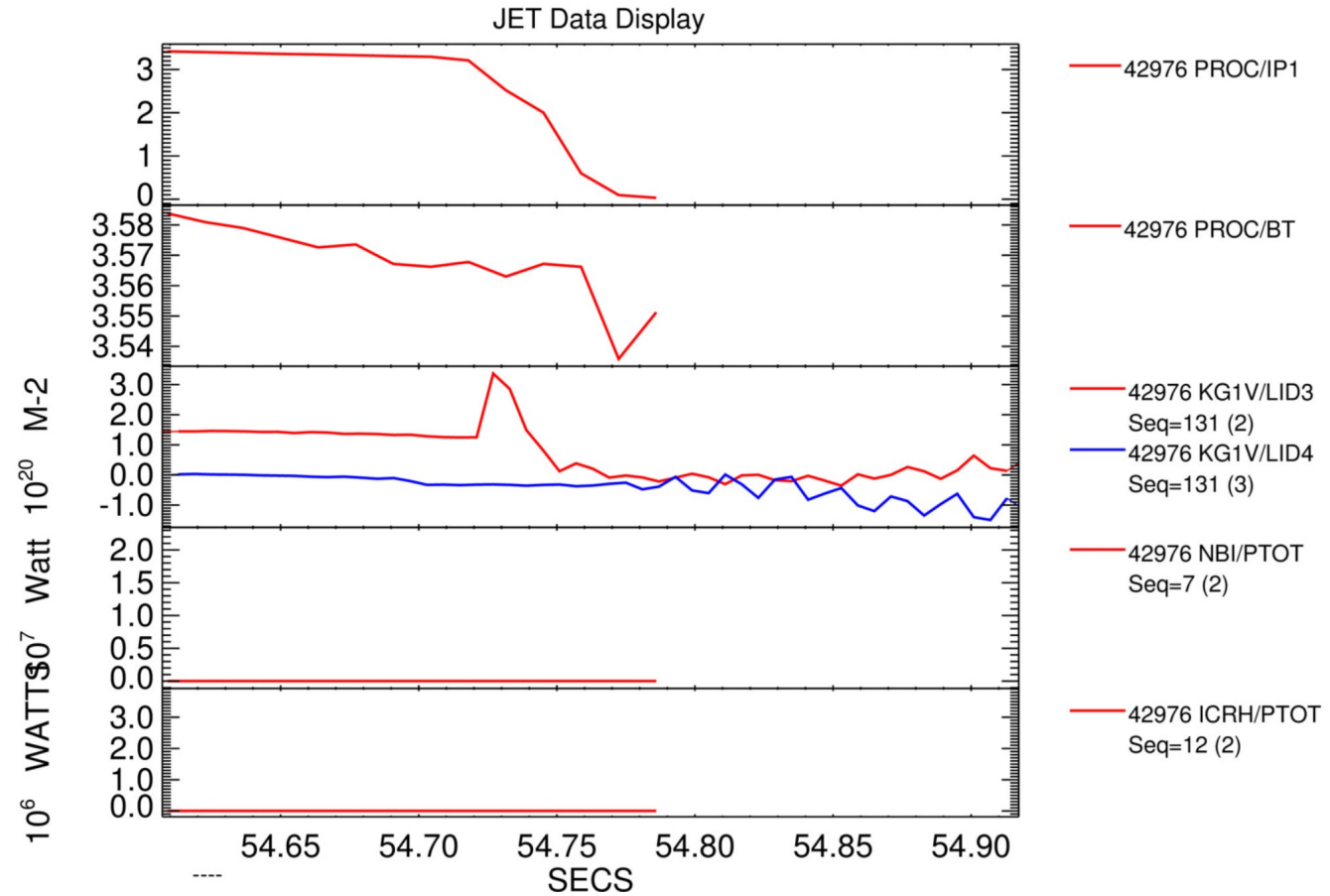
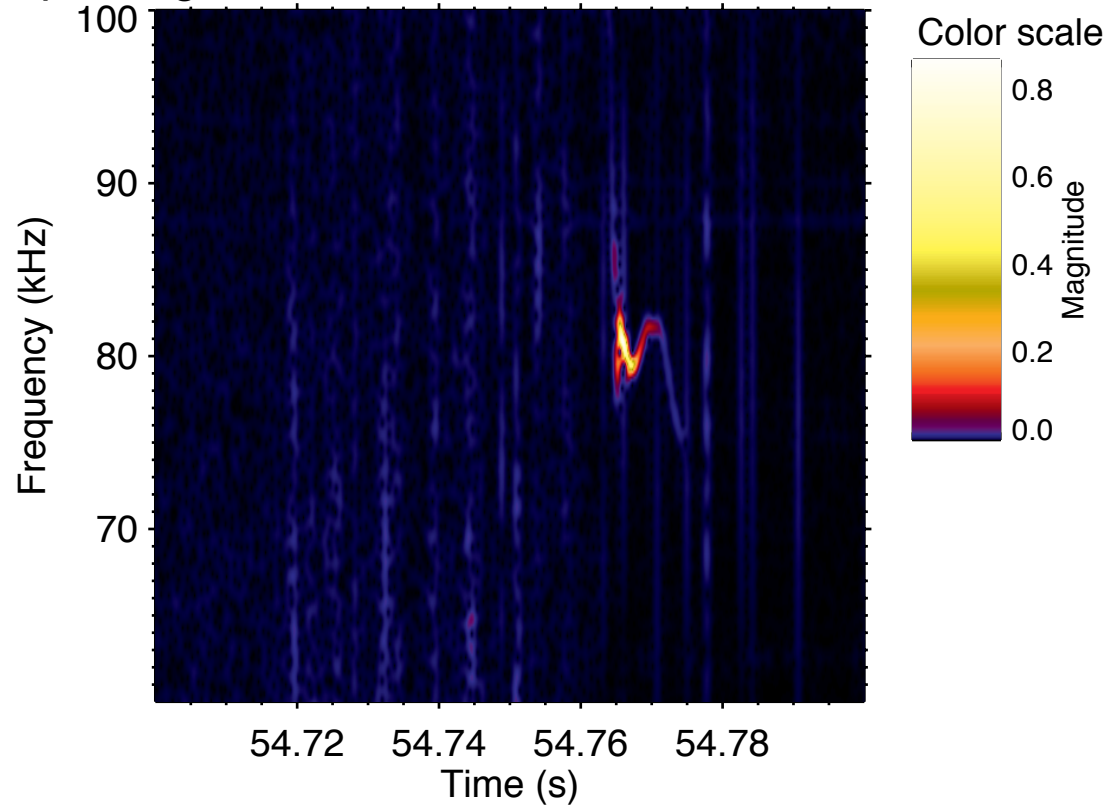
Selected TAE mode structures for $q_0=1.07$. Normalized amplitude of the real part with respective frequencies ω_{TAE} [kHz]. b) shows $m=n, n+1$ coupling and c) shows $m=n+1, n+2$ coupling.

¹[Kerner, JCP, 1998]

²[Gorelenkov, PS, 1992]

Backup – JET supershot

Spectrogram of JET 42976-DI/C1F-CHAN8/131



Backup – TEXTOR case

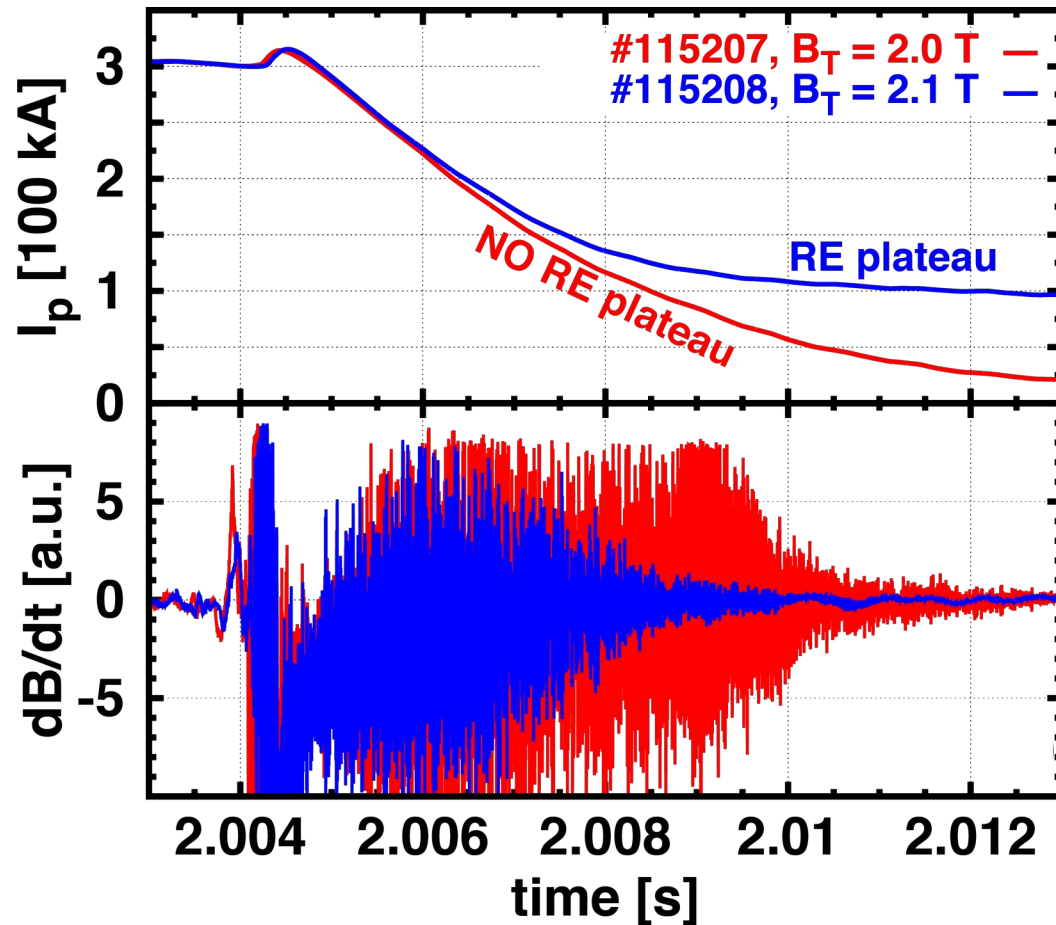


Fig.1: Plasma current evolution and Mirnov coil signals from TEXTOR shots #115207-8 (2013)

- Case of lower B_T shows strong coherent magnetic oscillations:
- Toroidal Alfvén Eigenmode (TAE) identified as the cause of runaway electron suppression
- Studies [1] estimated turbulence threshold level $dB/B \sim 0.1\%$ for suppression

↓
ITER: could alpha particles provide drive for RE-suppressing modes?

[1] Zeng, PRL, 2013

Backup – pressure profile

$$p(r, t) = n_e T_e(r, t) + n_{D, T} T_i(r, t) + p_\alpha(r, t)$$

with

$$n_e = n_{D, T} = 10^{20} [m^{-3}]$$

and

$$T(r, t) = T_f + [T(r, 0) - T_f] \exp(-t/t_0)$$

$$p_\alpha(r, t) = \frac{m_\alpha}{3} \int v^4 f_\alpha(v, r, t) dv$$



CODION is 0D in space: Each of the 100 radial points is populated by velocity distributions $f_\alpha(v, r, t)$ advancing independently in time



Assumes case of good post-disruption confinement, as is also necessary for RE beam

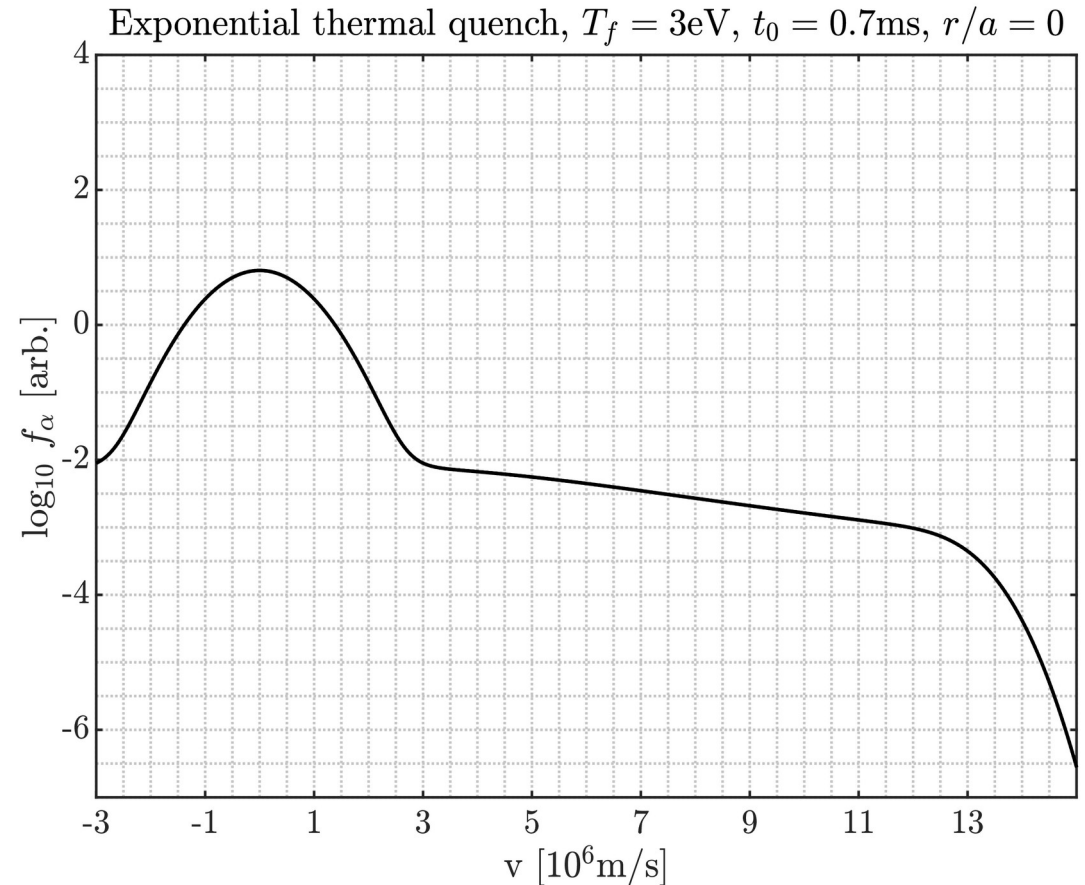
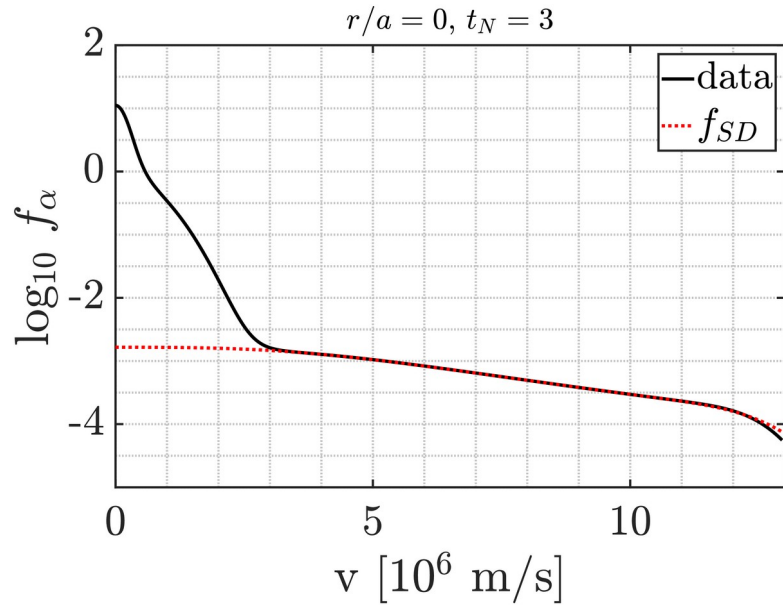


Fig.2

CODION simulation: Initial fusion born alpha particle population on axis $f_\alpha(v, 0, 0)$

Backup - active mode evolution



Energetic part of CODION obtained data is fitted with the analytic slowing down formula¹ f_{SD}

$$f_{SD}(r, v, t_N=3) = \frac{C(r)}{v_c^3(r) + v^3} \text{Erfc}\left(\frac{v - v_\alpha}{\Delta v(r)}\right)$$

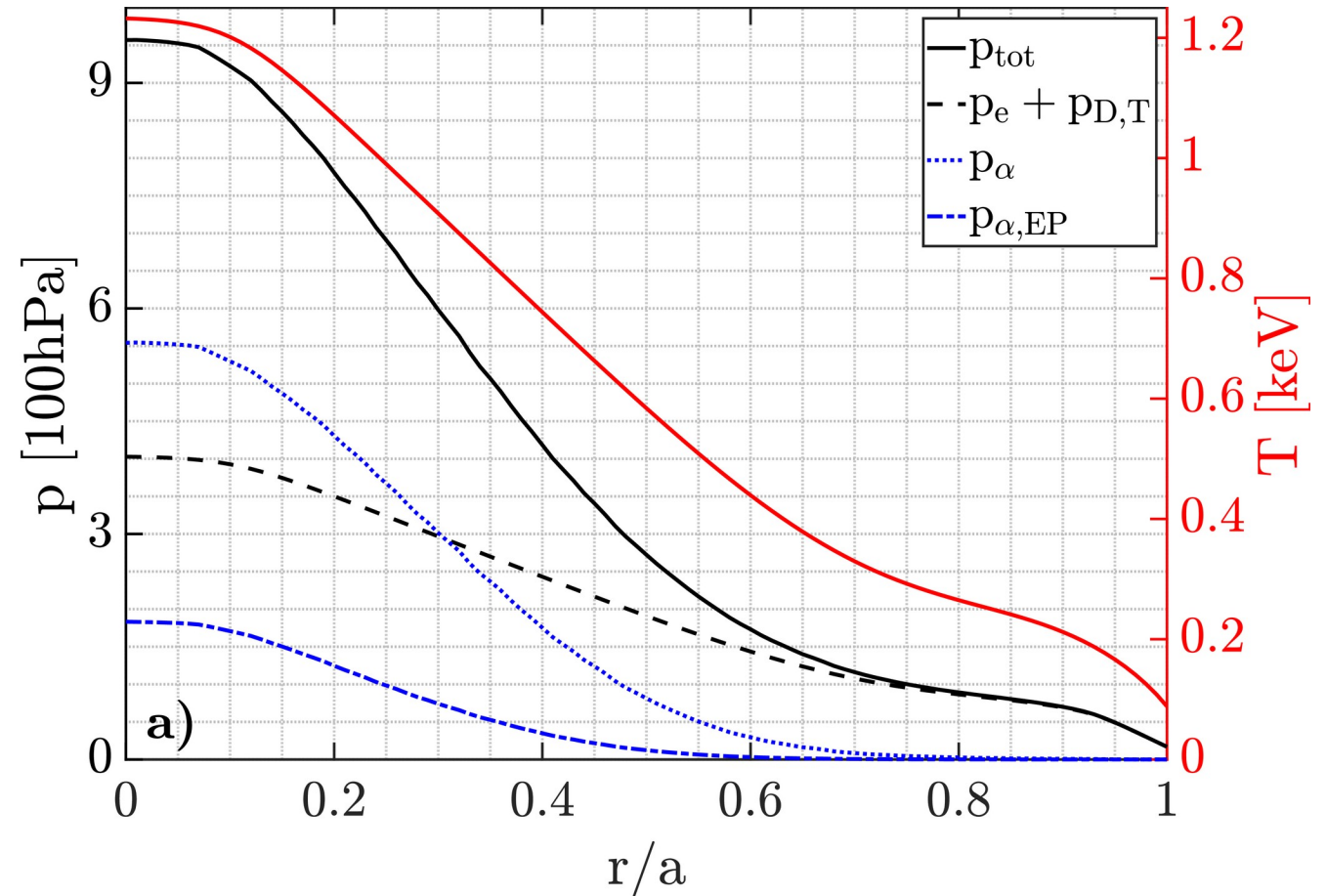


Fig.7

Pressure and temperature profiles for $t_N=3$. Note that $p_{\alpha,EP}$ is used for mode drive, not p_α , since the latter is misleading (due to CODION particle conservation)

Backup plots

