

# Density and Temperature Profiles after Low-Z and High-Z Shattered Pellet Injections on DIII-D

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with participation of

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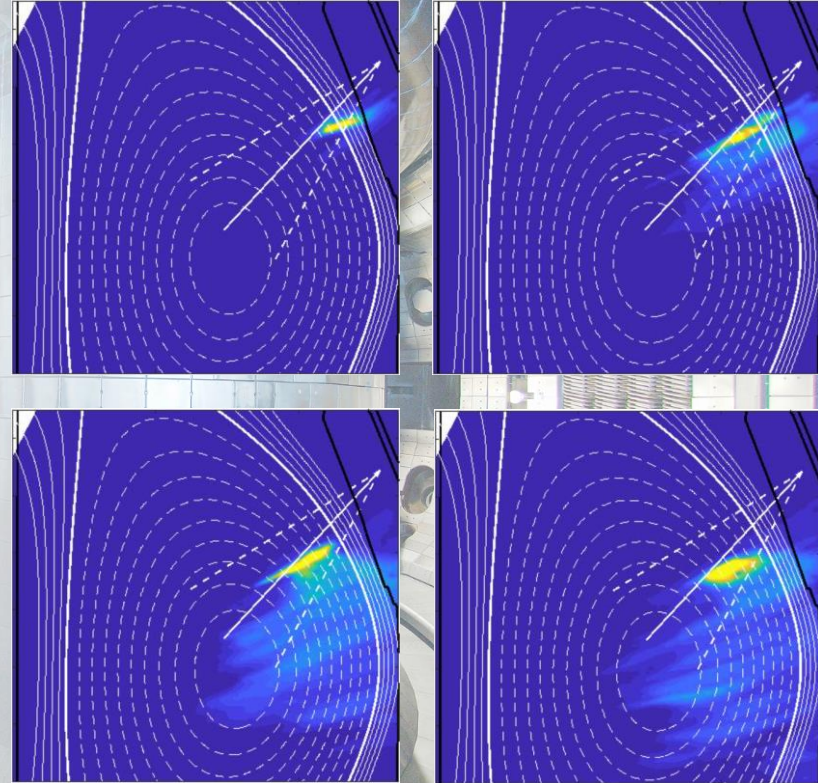
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<sup>3</sup>Oak Ridge National Laboratory, Oak Ridge, TN, USA

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**Princeton, July 19–21 2023**



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- Shattered pellet injection (SPI) is chosen for disruption mitigation on ITER

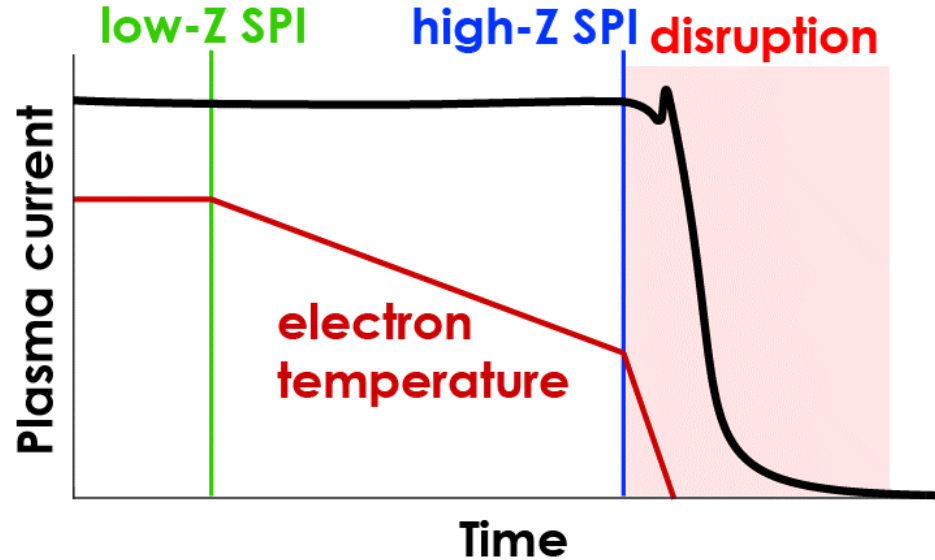
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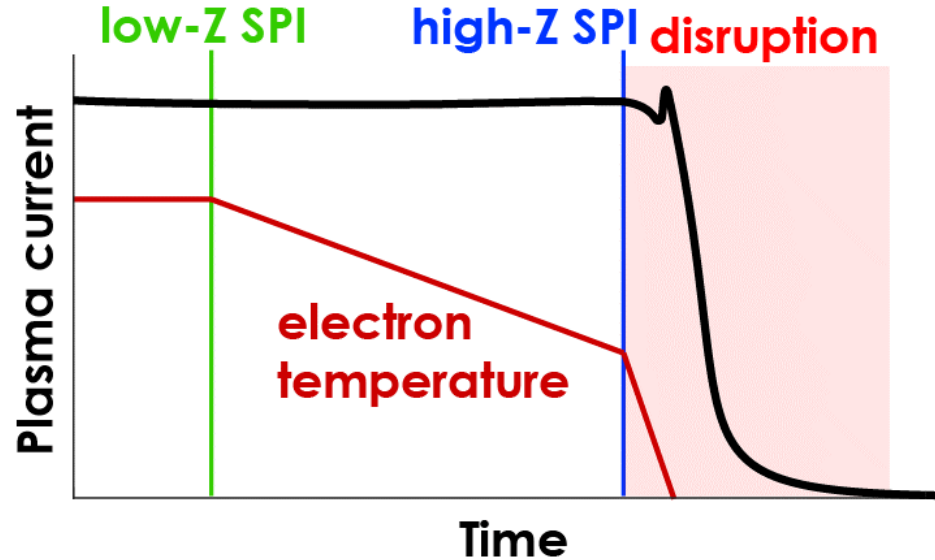


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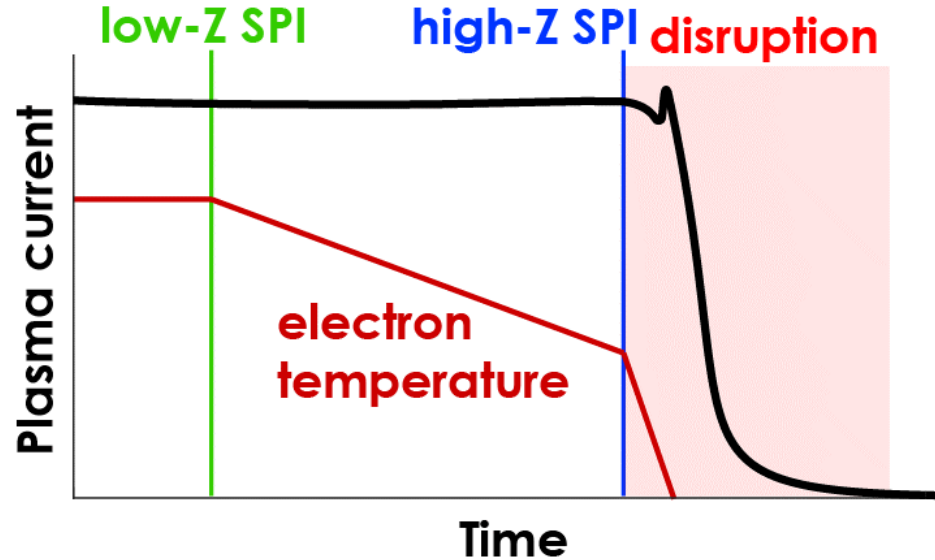


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• This talk: study of temperature and density profiles after low- and high-Z SPI on DIII-D

Introduction

# Experimental setup

Mixed Ne/D<sub>2</sub> SPI

Pure D<sub>2</sub> SPI

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Comparison with experiment

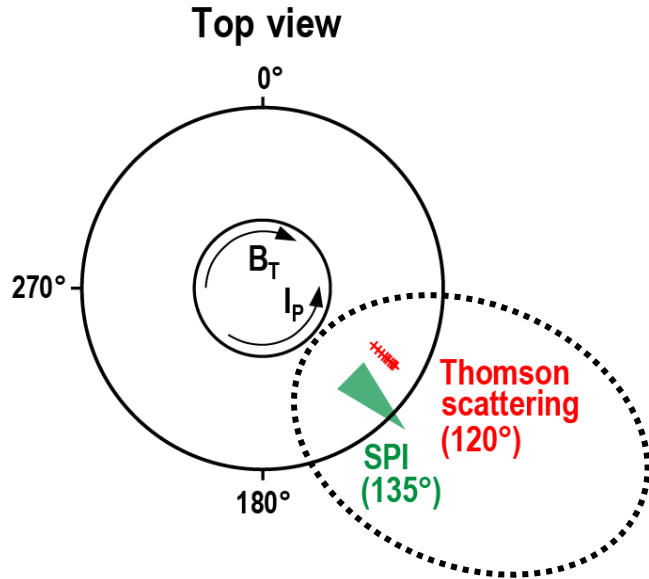
How to improve penetration

Final remarks

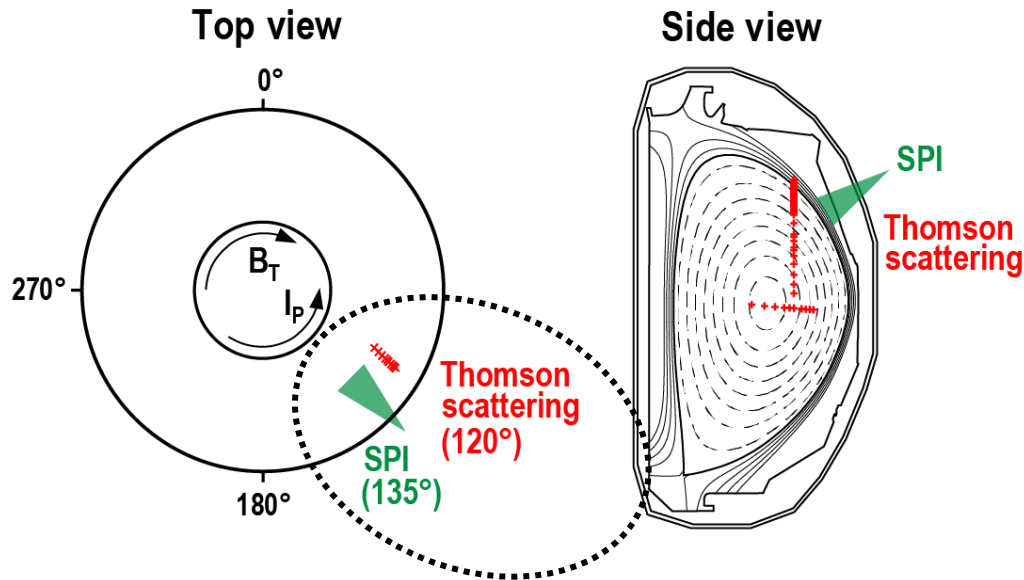


# DIII-D is well suited to study shattered pellet injection

- The SPI system closest to Thomson scattering viewpoints is analyzed in this talk

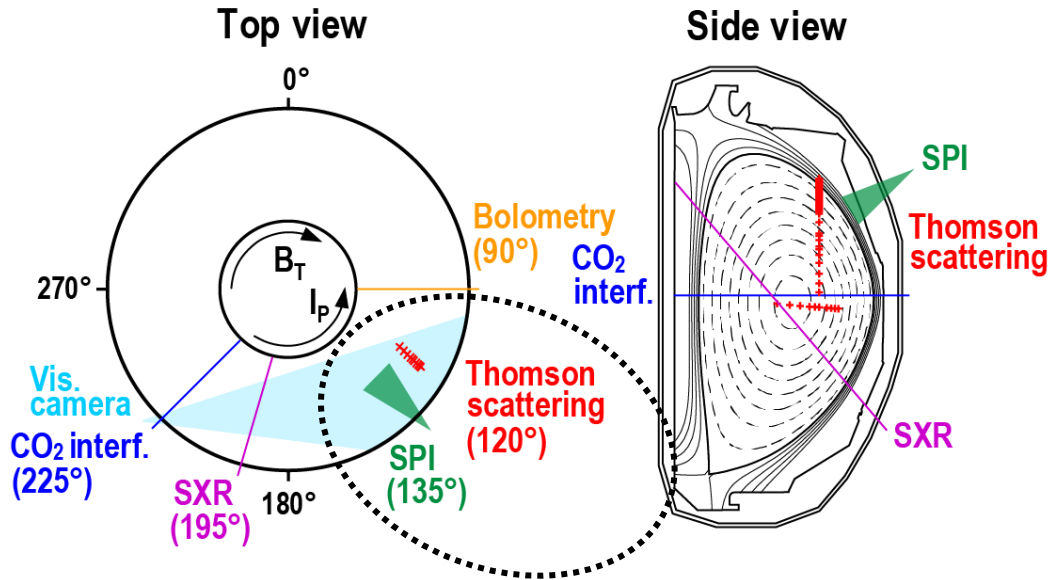


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- Thomson scattering provides plasma temperature and density profiles
  - Covers both plasma core and edge
  - Recent upgrade with narrow-bandwidth filters allows low- $T_e$  measurements
  - Can be triggered asynchronously (e.g. using ablation light signal) to reliably catch plasma cooling dynamics

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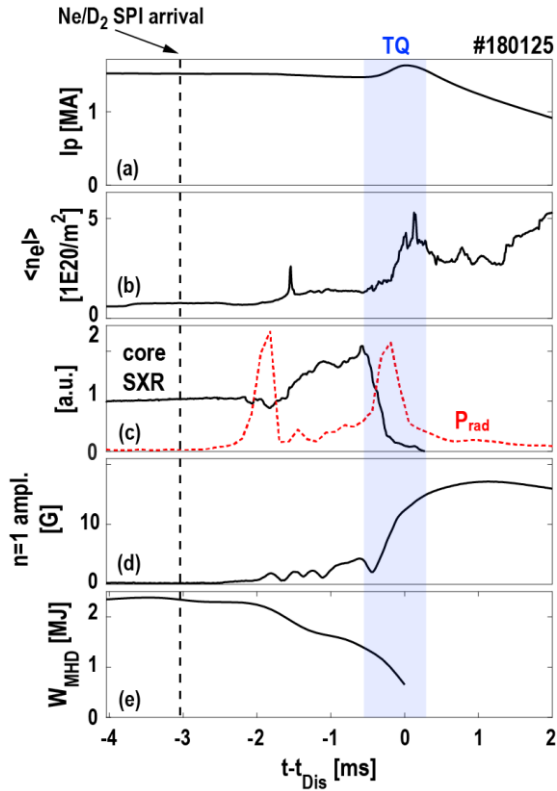
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- Other diagnostics provide **pellet imaging**, **line-integrated density**, **soft X-ray**, **total radiated power**

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**Mixed Ne/D<sub>2</sub> SPI**

Pure D<sub>2</sub> SPI  
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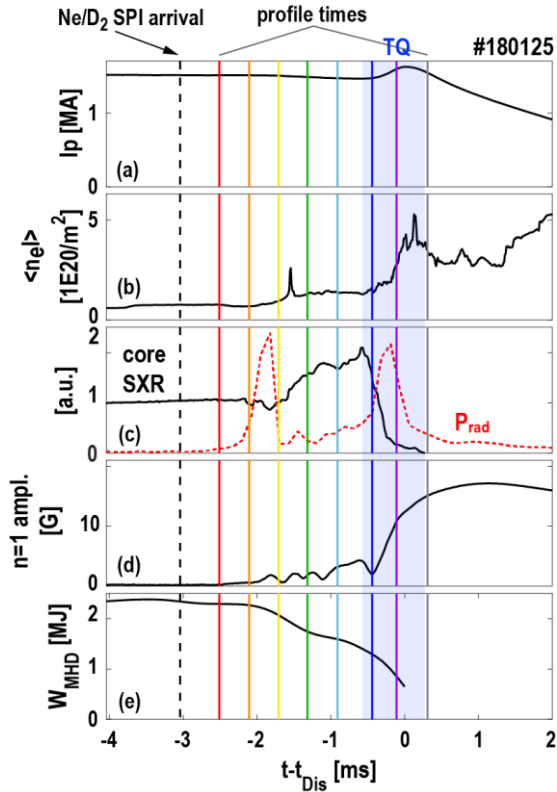
Comparison with experiment  
How to improve penetration  
Final remarks

# Mixed Ne/D<sub>2</sub> SPI quickly disrupts the plasma



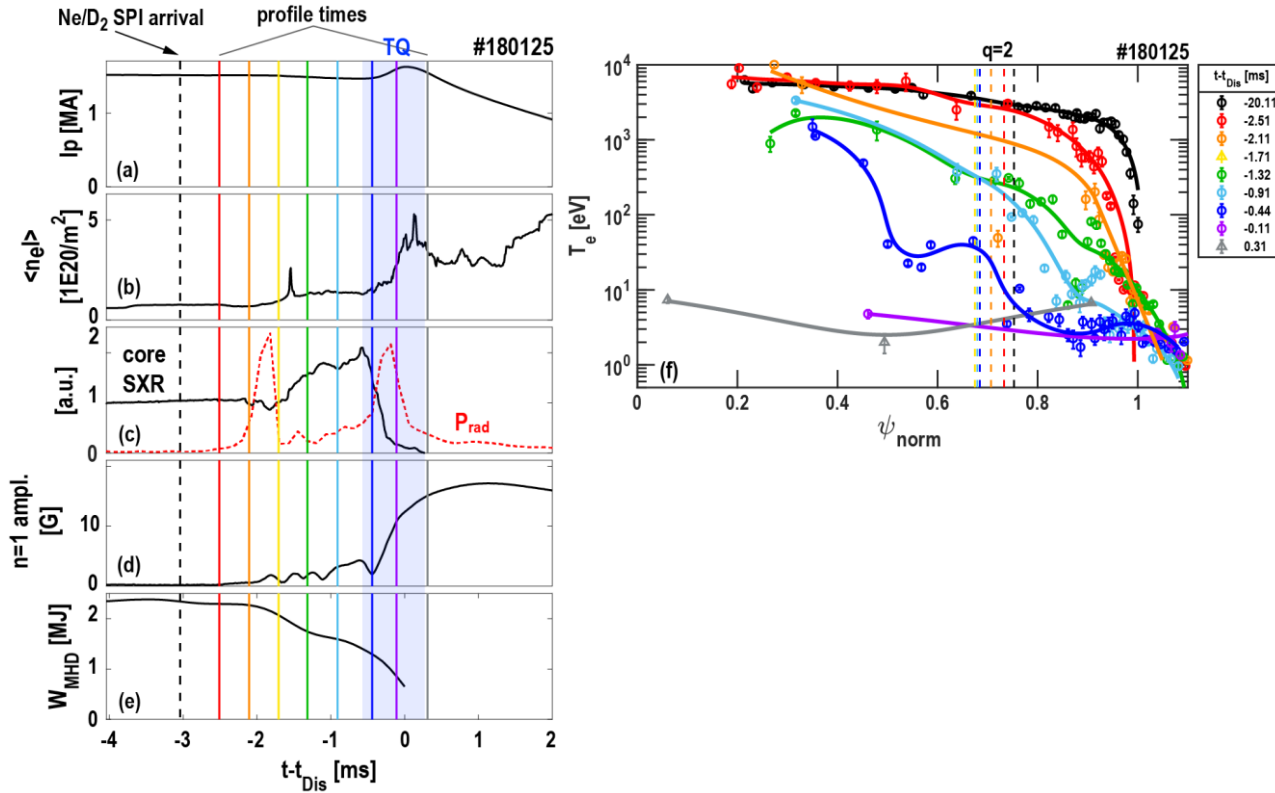
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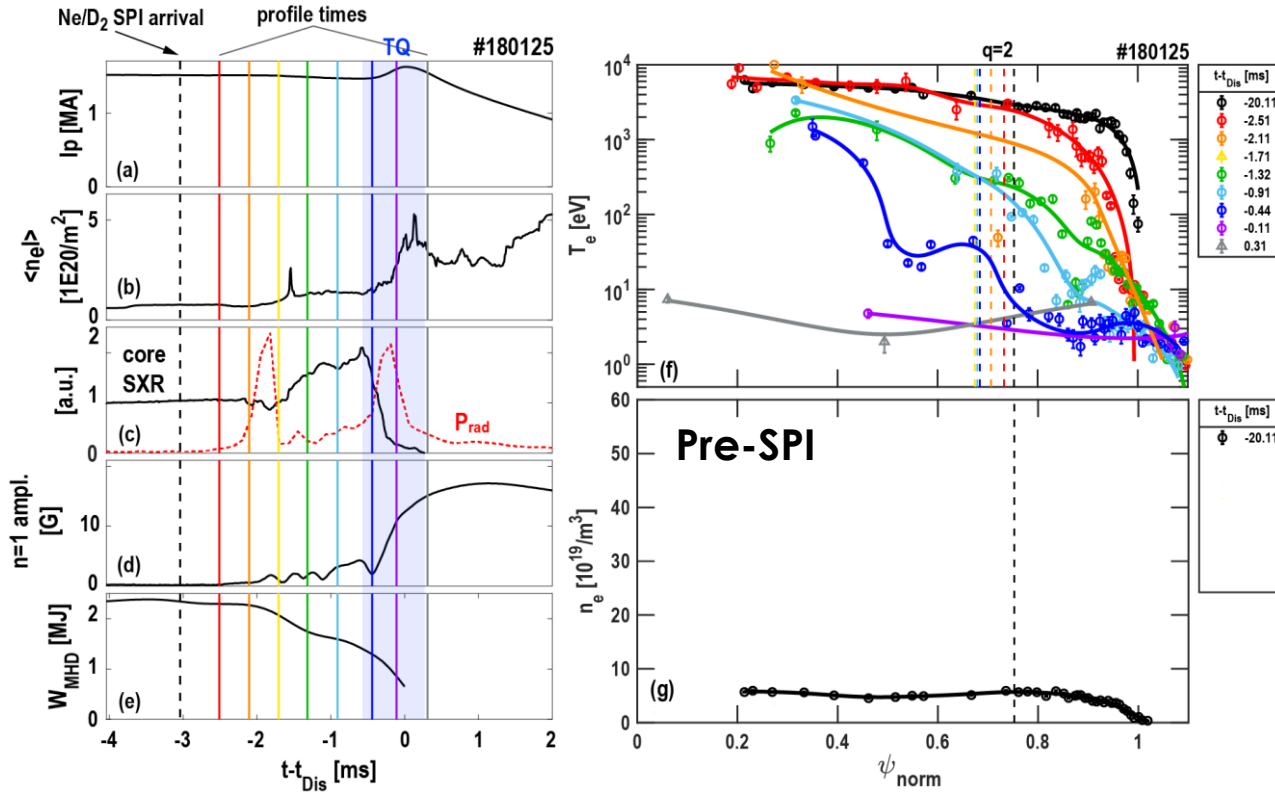
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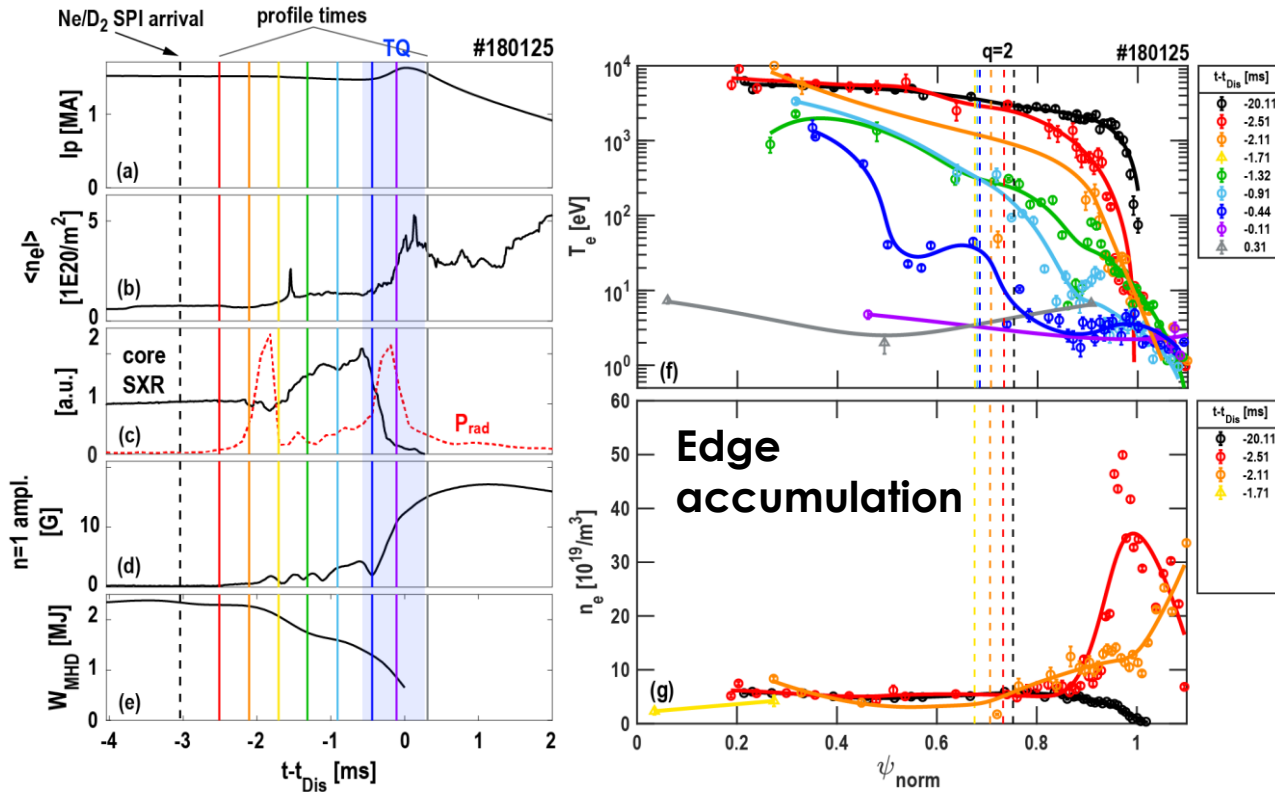
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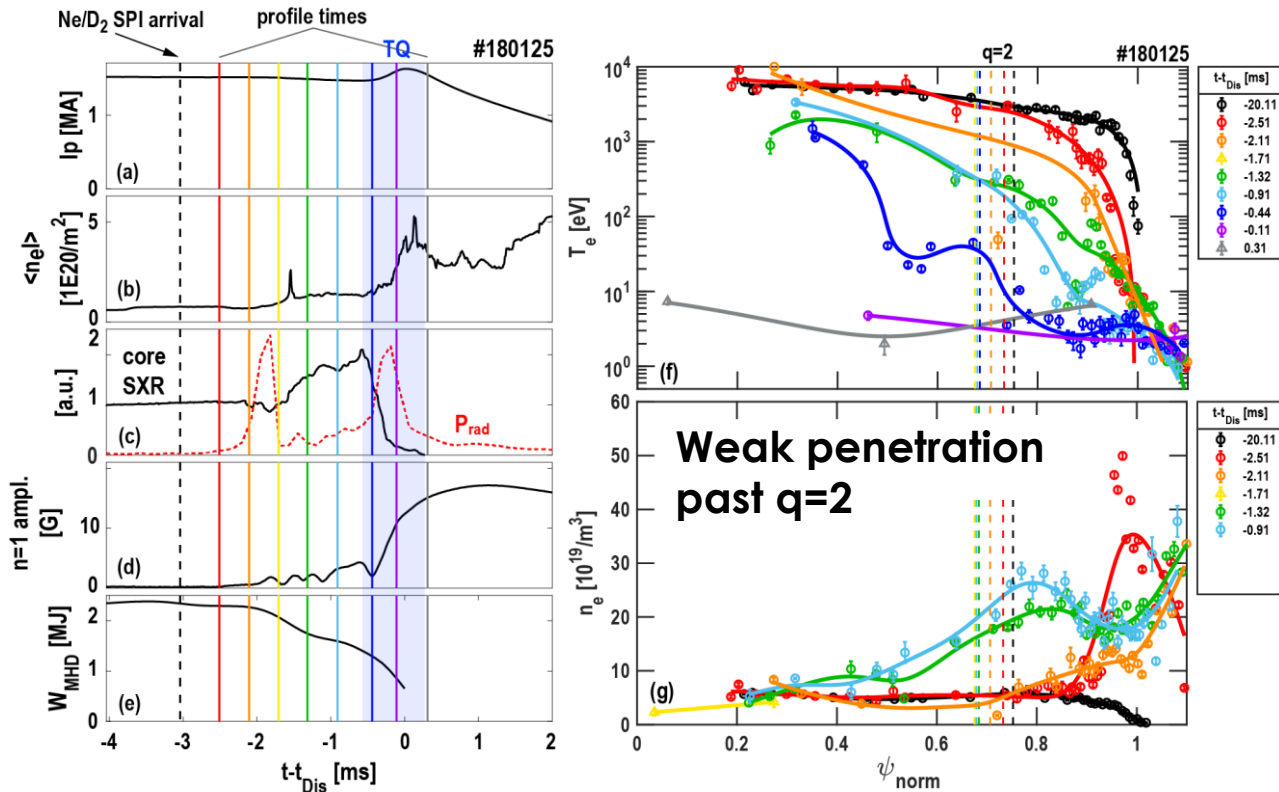


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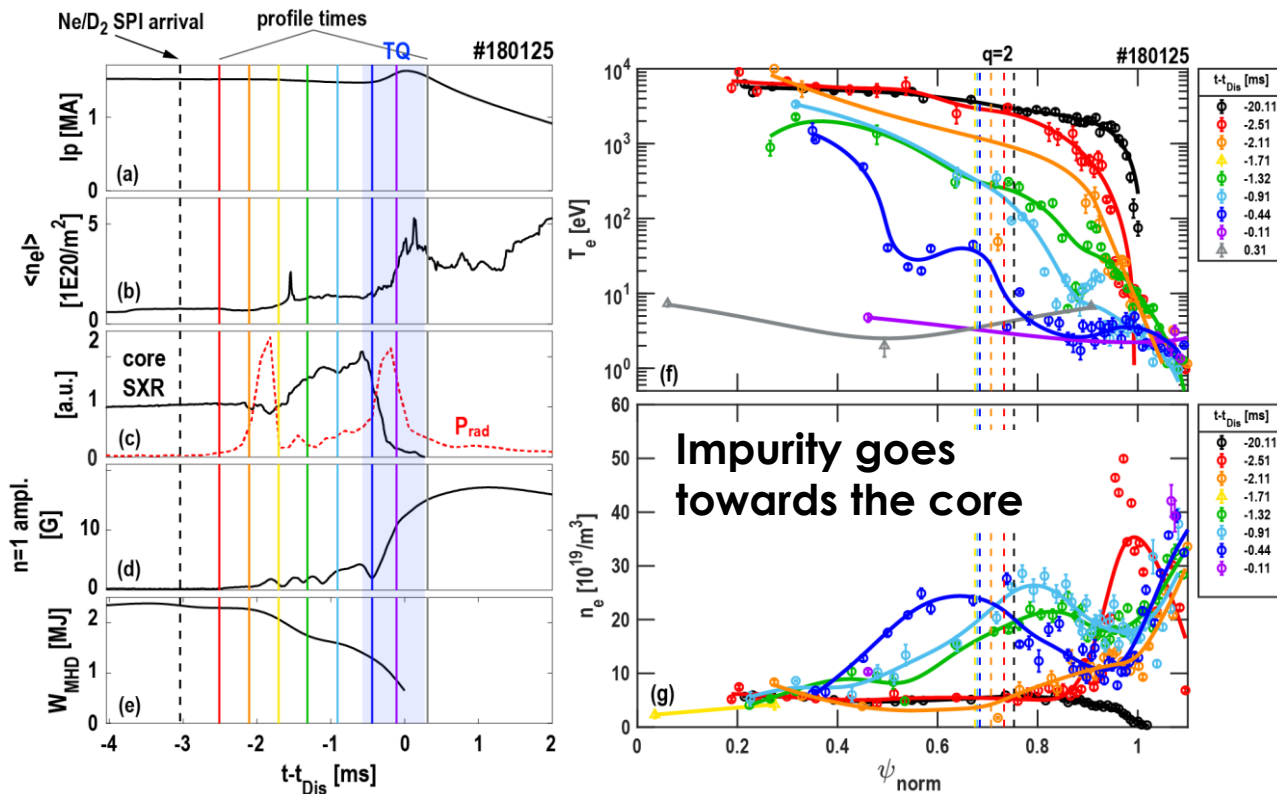
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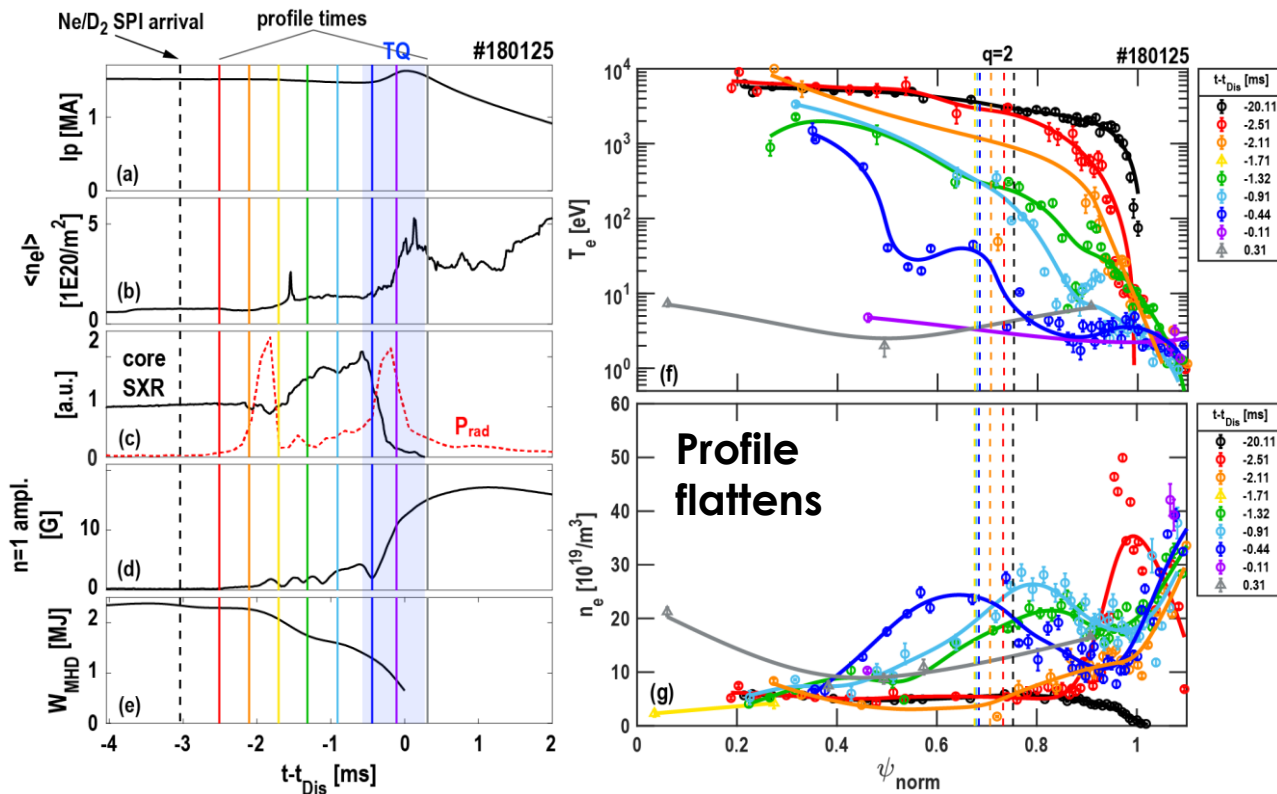
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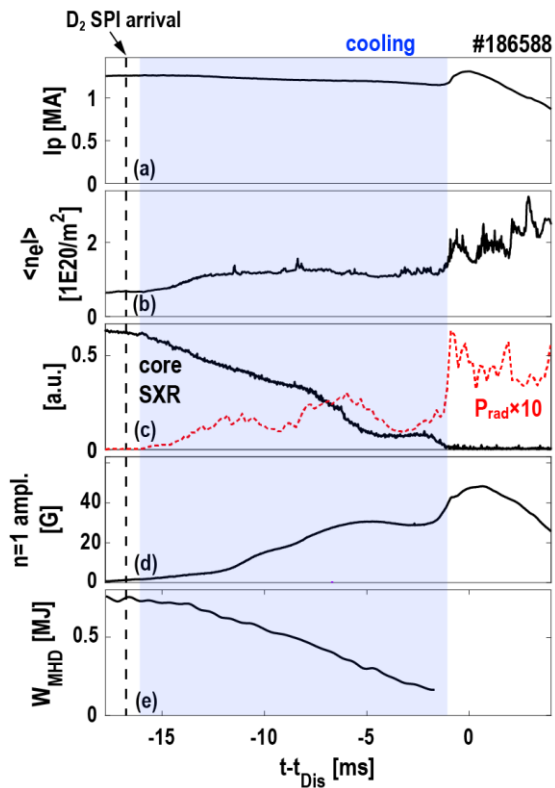


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  - Core density increases by 2x–4x after TQ

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Experimental setup  
Mixed Ne/D<sub>2</sub> SPI  
**Pure D<sub>2</sub> SPI**  
INDEX code

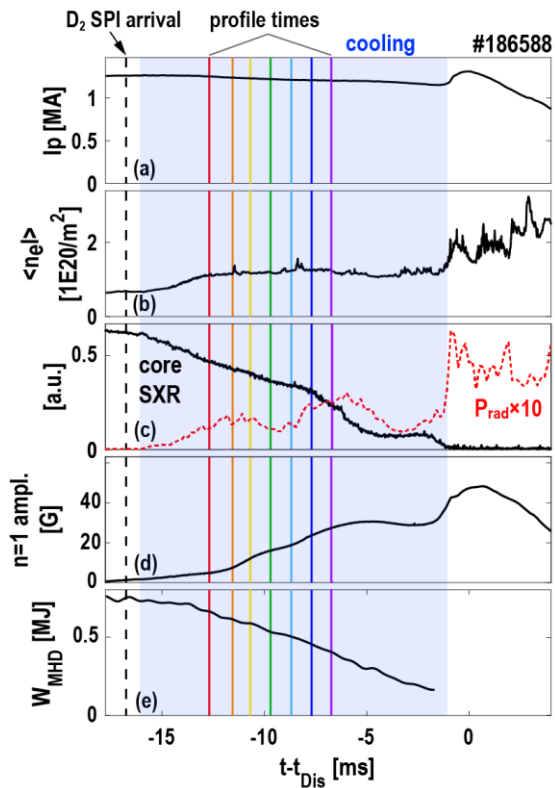
Comparison with experiment  
How to improve penetration  
Final remarks

# D<sub>2</sub> SPI slowly cools the plasma, it disrupts due to locked modes



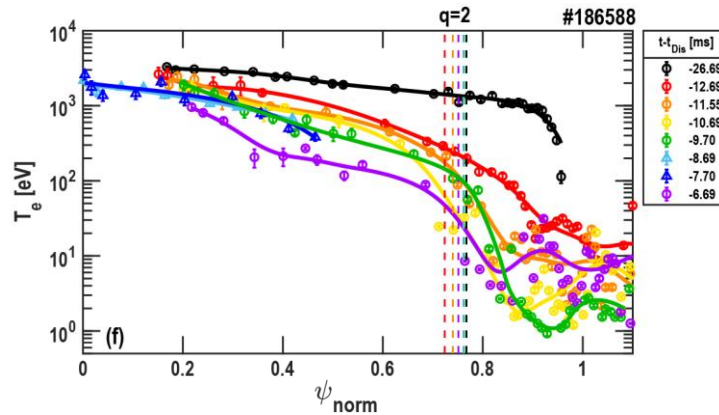
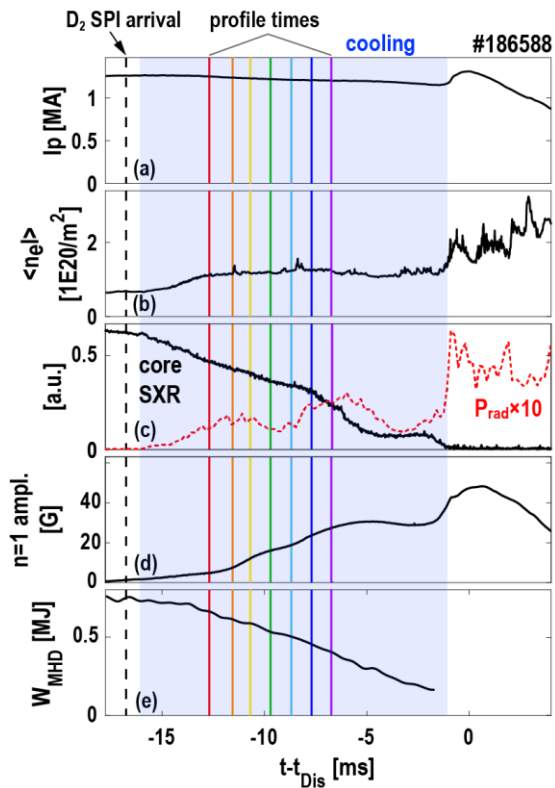
- Pure D<sub>2</sub> shattered pellet (700 Torr·L) is injected into H-mode plasma
- It slowly cools plasma via dilution and radiation over more than 15 ms
  - No classical TQ is observed
- Plasma energy decreases 4-fold by the  $I_p$ -spike time
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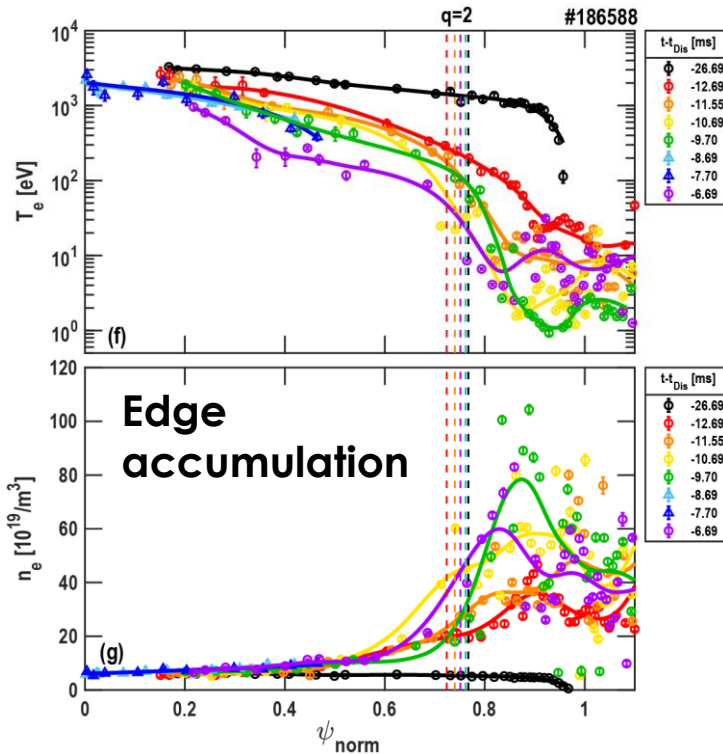
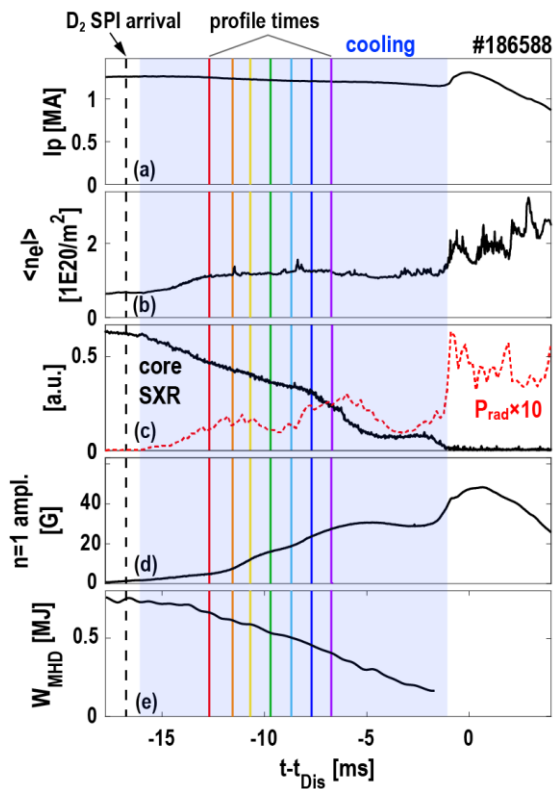
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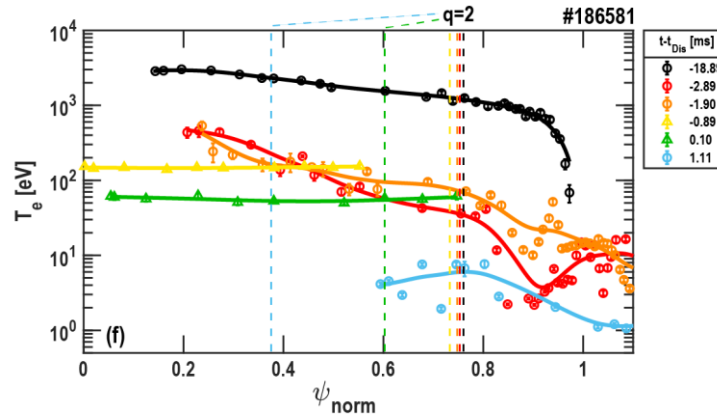
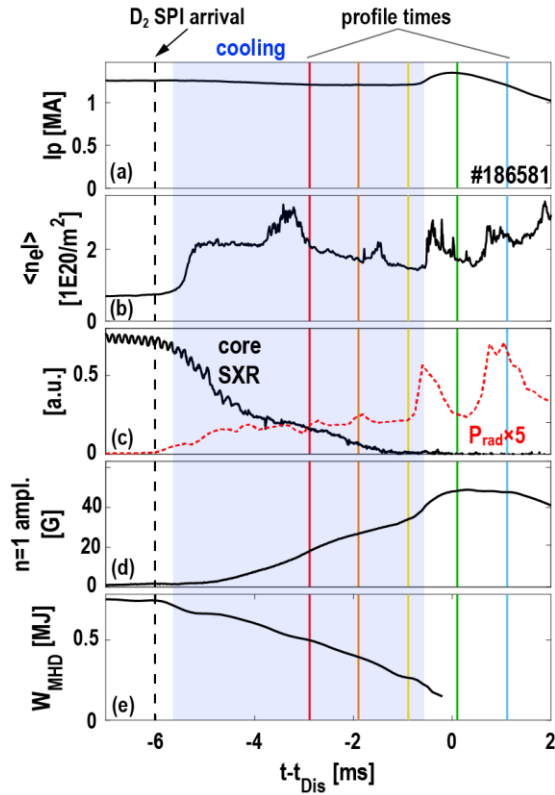


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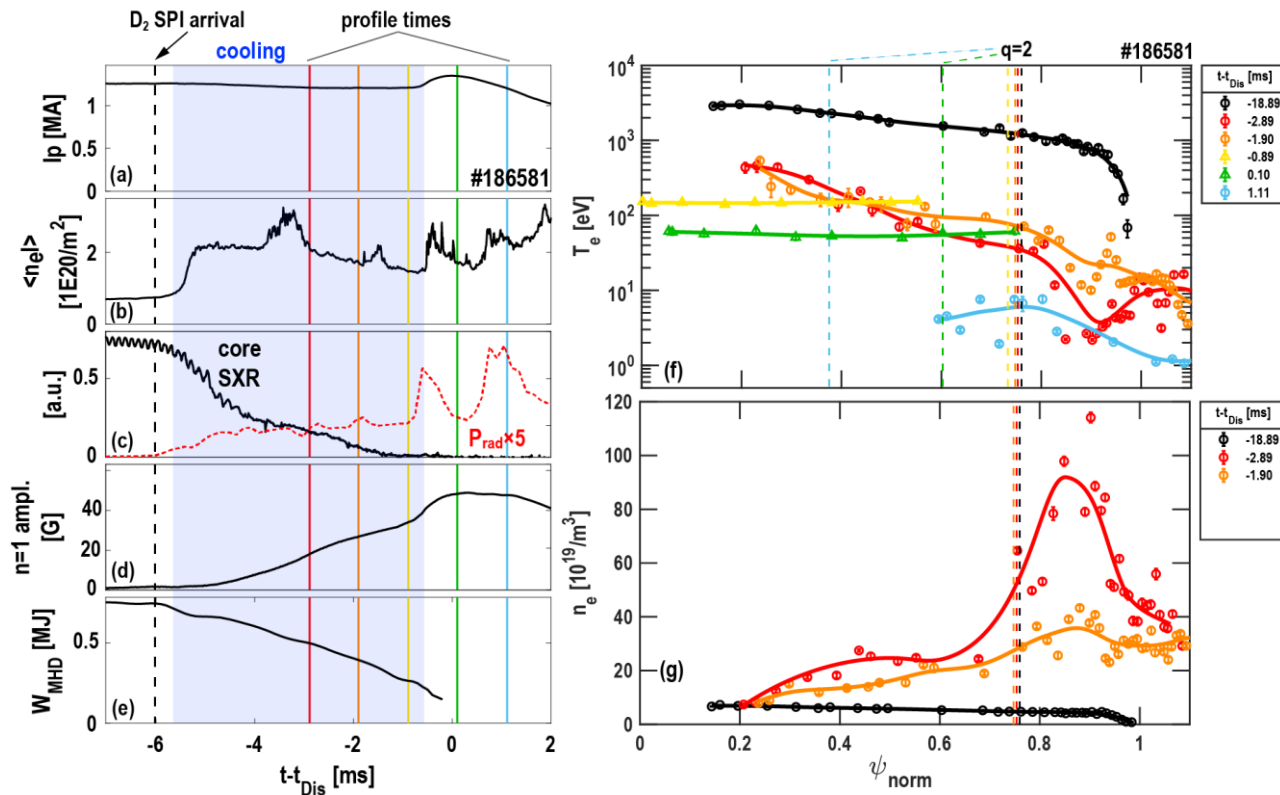
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- Very weak penetration past  $q=2$  is observed before TQ

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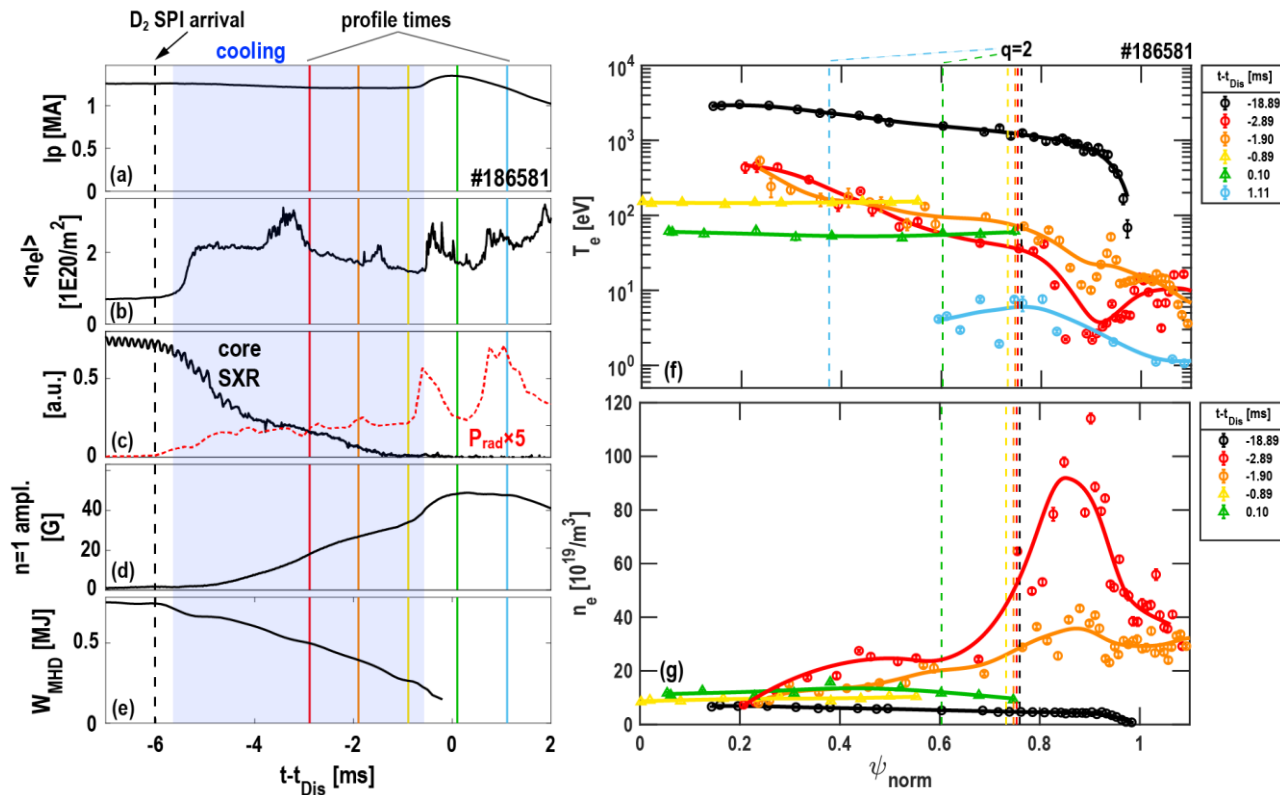
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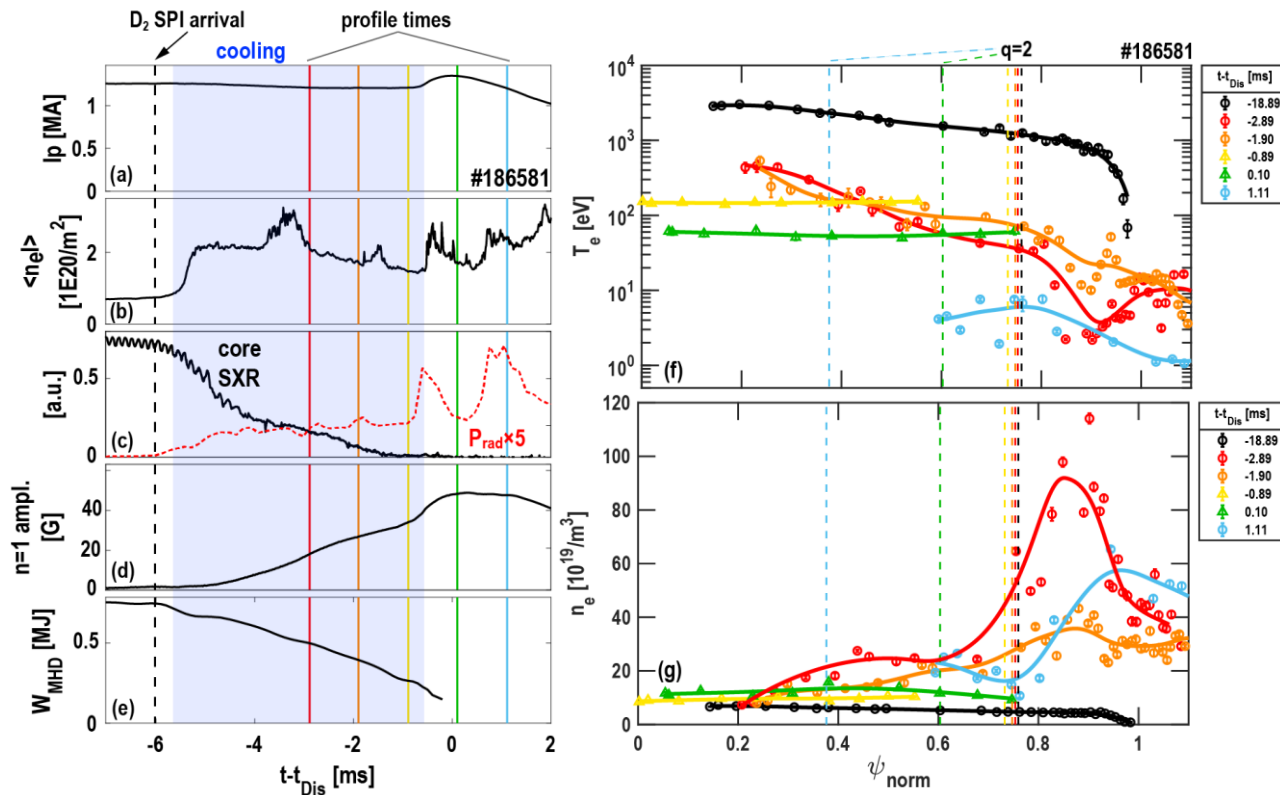
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- Density does not flatten even after the Ip-spike

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Pure D<sub>2</sub> SPI

# **INDEX code**

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# INDEX simulations are employed to interpret observed post-SPI density profiles

- INDEX is a 1.5D code self-consistently solving transport equations in magnetic flux coordinates coupled to Grad-Shafranov equilibrium calculations<sup>1,2</sup>
- It evaluates the SPI ablation rate based on the Neutral Gas Shielding model<sup>3</sup>
- The ablated particles are included as the surface-averaged neutral source in the particle balance, and the subsequent ionization processes are solved using the rate equations
- The SPI model was extended to include the effect of outward mass relocation due to the grad-B-induced (ExB) drift<sup>4-6</sup>
- Presently grad-B-induced drift is not based on first principles but employs simplified considerations following the back-averaged model<sup>7</sup>

[1] Matsuyama et al, IAEA FEC 2021

[2] Matsuyama et al, PPCF 2022

[3] Parks and Turnbull, Phys. Fluids 198

[4] Rozhansky et al, PPCF 1995

[5] Parks et al, PoP 2000

[6] Pegourie et al, NF 2006

[7] Jardin et al, NF 2000



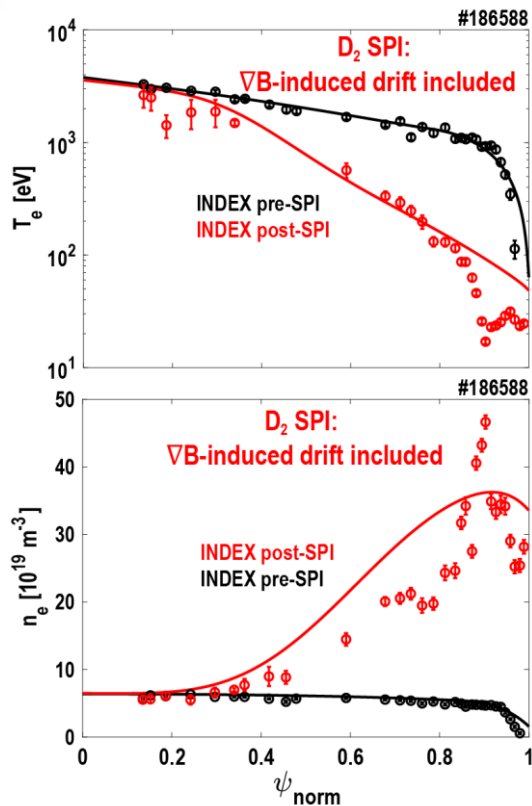
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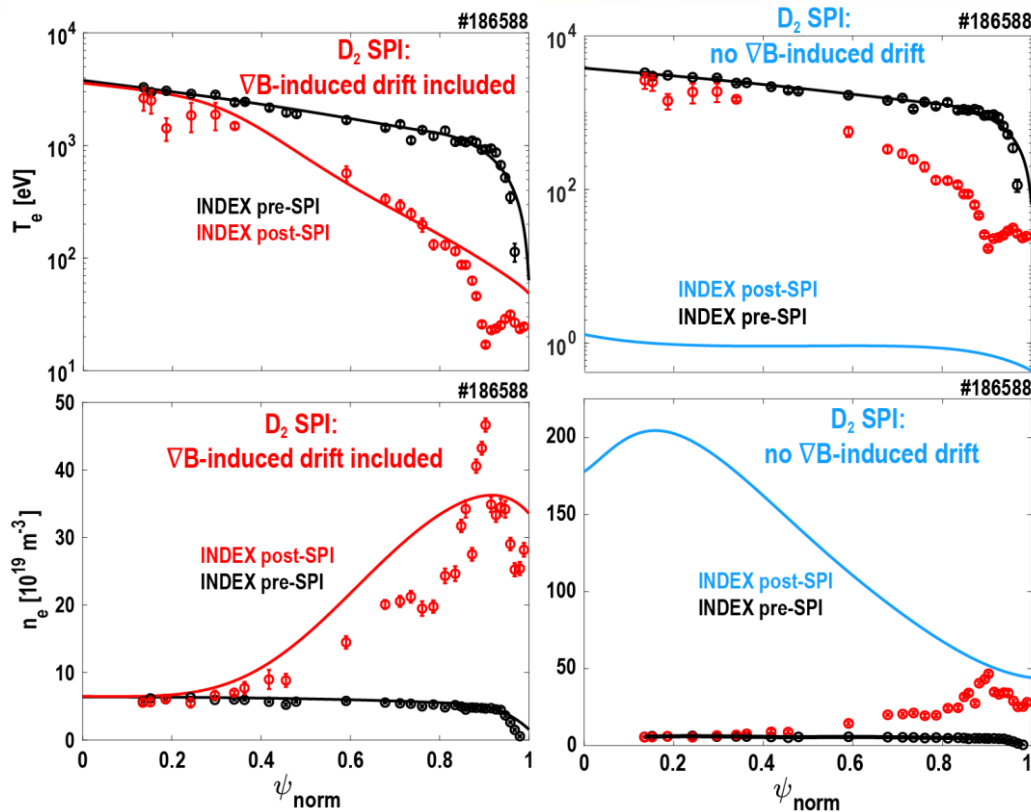


# Grad-B-induced drift of ablation cloud to the edge can explain limited core fueling after D<sub>2</sub> SPI



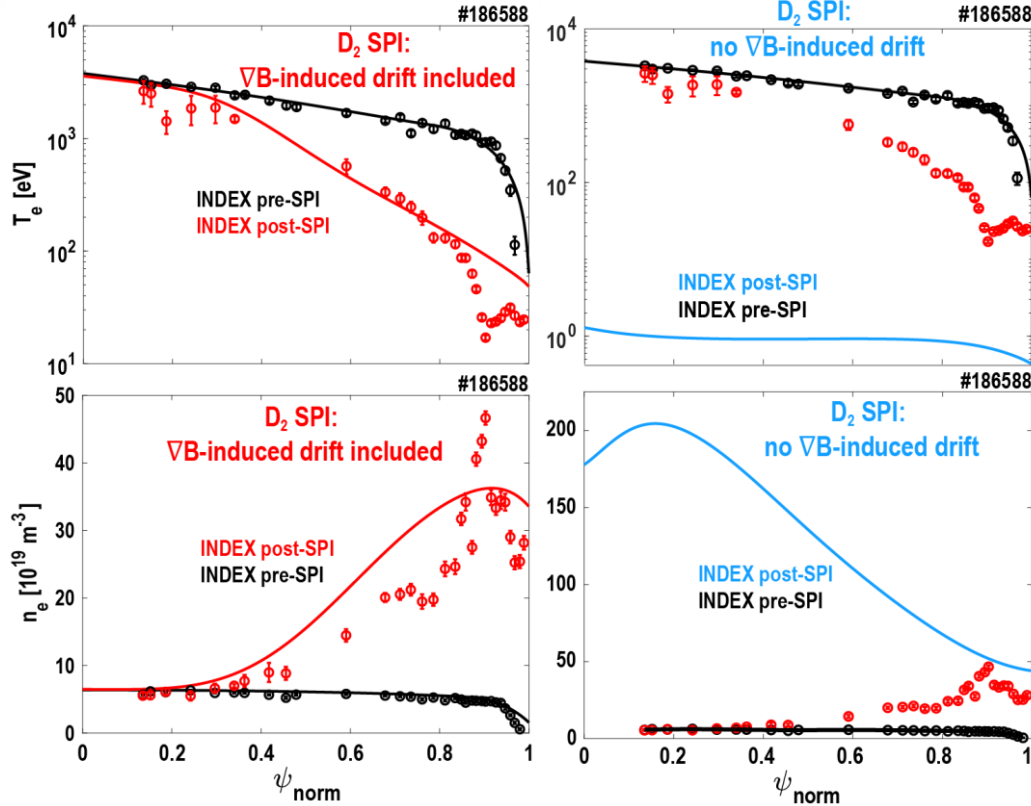
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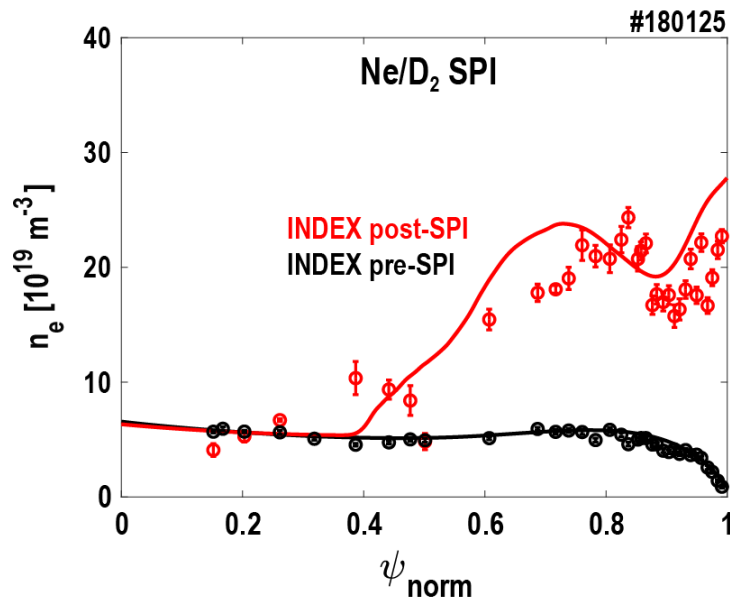
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- Estimated that only 10% of injected material is assimilated
- Density profile w/o drift is clearly overestimated
- Presumably, low line radiation of D plasmoid, heating by surrounding plasma and overpressure make the drift so significant ( $dV_{\text{drift}}/dt \propto nT$ )

# Density profile after mixed Ne/D<sub>2</sub> SPI is well reproduced without grad-B-induced drift enabled



- For comparison, no drift implication needed to reproduce the experiment with mixed Ne/D<sub>2</sub> SPI
- Though it does not simplify these simulations
- Increased radiative cooling may make surface-averaged ablation rate less accurate
  - Accounted via reduced ablation<sup>1</sup>
- Non-thermal electrons and pre-TQ MHD may enhance pellet ablation
  - Accounted via increased thermal conductivity<sup>2</sup>

# Drift of ablated material is known for fueling pellets, but experimental profiles after SPI were obtained for the first time

- Small solid fueling pellets are known to produce plasmoids drifting in inhomogeneous magnetic field<sup>1-7</sup>
- There is a bunch of theoretical models describing such experiments<sup>8-15</sup>
- However, to our knowledge, there is no experimentally measured deposition of large perturbative SPI nor models<sup>16\*-17\*</sup> interpreting experimental profiles
- In this work such observations are presented for the first time and a simplified model is developed to explain and predict density profiles after pure D<sub>2</sub> and mixed Ne/D<sub>2</sub> SPI

[1] Lang et al, PRL 1997

[2] Mueller et al, PRL 1999

[3] Baylor et al, PoP 2000

[4] Baylor et al, NF 2007

[5] Terranova et al, NF 2007

[6] Garzotti et al, NF 2010

[7] Baldzuhn et al, PPCF 2019

[8] Strauss et al, PoP 1998

[9] Strauss et al, PoP 2000

[10] Parks et al, PoP 2000

[11] Rozhansky et al, PPCF 2003

[12] Aiba et al, JPS 2004

[13] Pegourie et al, NF 2007

[14] Ishizaki et al, PPCF 2011

[15] Vallhagen et al, JPP 2023

[16\*] Samulyak et al, NF 2021

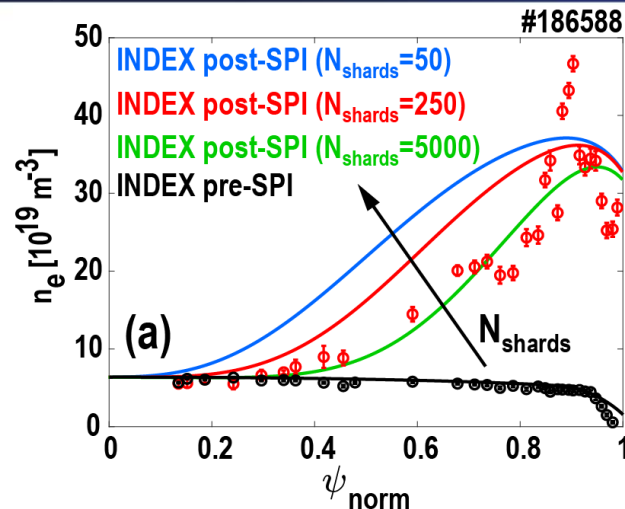
(no comparison with experiment)

[17\*] Kong et al, EPS 2022

(comparison with interferometer)

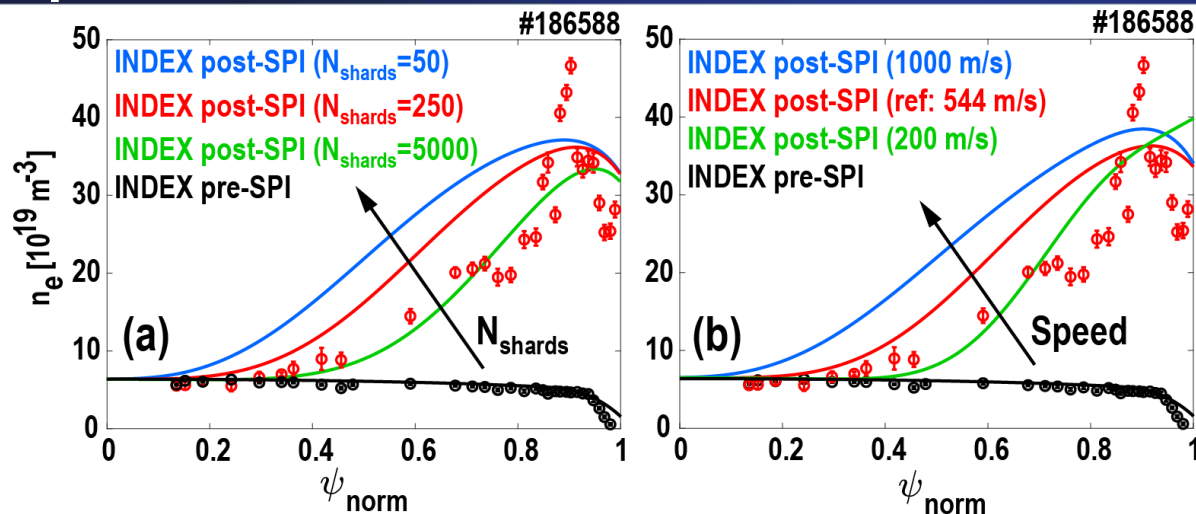
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# Modeling suggests that larger shards can improve core penetration



- Increasing shard size improves D penetration into the plasma
  - Larger shards have more time to ablate

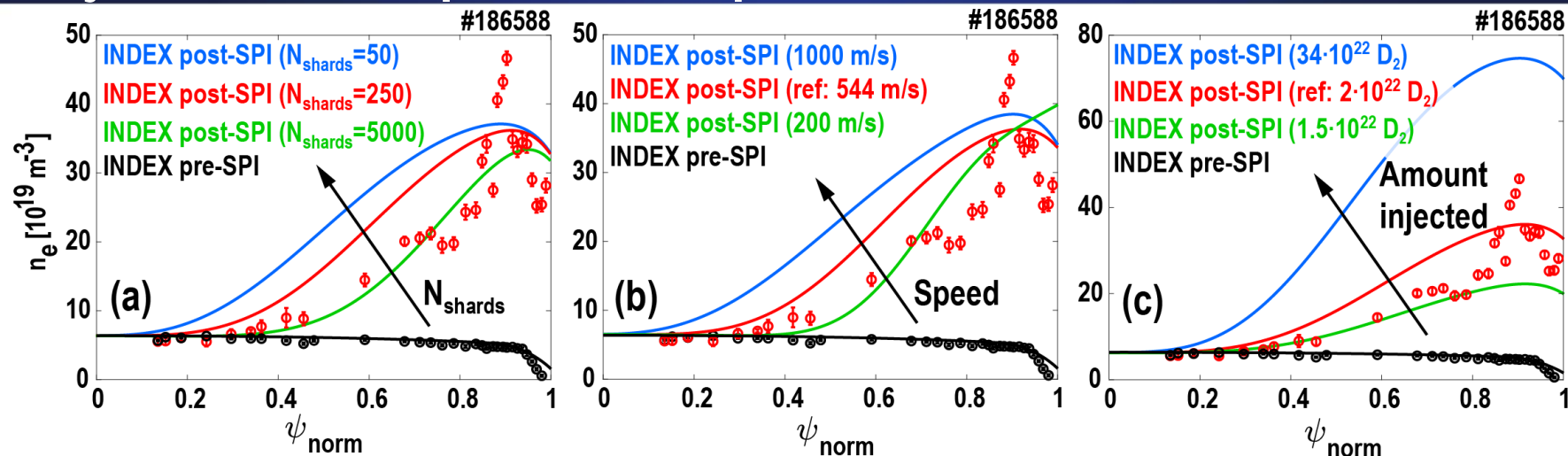
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- **Increasing mean speed of shards also improves penetration**
  - Though in a real experiment greater pellet speed would lead to smaller shards



# Modeling suggests that larger shards and greater amount injected can improve core penetration



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  - Larger shards have more time to ablate
- **Increasing mean speed of shards also improves penetration**
  - Though in a real experiment greater pellet speed would lead to smaller shards
- **Increasing amount injected can also help with core fueling**
  - Forefront shards cool the plasma reducing ablation of following shards

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# Summary and conclusions

- Staggered scheme of low-Z and high-Z SPI is expected to reduce thermal loads and provide improved RE suppression
- Experiments on DIII-D show favorable slow plasma cooling after  $D_2$  SPI but poor assimilation of  $D_2$  by core plasma even during and after TQ
- 1D modeling suggests grad-B-induced drift of deuterium ablation cloud towards the edge caused by low line radiation and heating
- Larger pellet shards and greater amount injected are predicted to improve core fueling
- Optimization of  $H_2$  SPI in ITER is necessary

Paper is to be submitted to Nuclear Fusion

# Future work

- Experimentally verify predictions of 1D model (very recent study on DIII-D, yet to be analyzed)
- Develop mature model based on first principles
- Improve model to:
  - include simulations of Ne-doped  $D_2$  SPI
  - separate effects of grad-B-induced drift, non-thermal electrons, pre-TQ MHD for mixed Ne/ $D_2$  SPI
  - study whole post-SPI dynamics and estimate total material assimilation

# Acknowledgments and Disclaimer

- This material is based upon work supported by the U.S. Department of Energy, Office of Science, Office of Fusion Energy Sciences, using the DIII-D National Fusion Facility, a DOE Office of Science user facility, under Award DE-FC02-04ER54698 and Service Contract IO/19/CT/4300001939.
- This report was prepared as an account of work sponsored by an agency of the United States Government. Neither the United States Government nor any agency thereof, nor any of their employees, makes any warranty, express or implied, or assumes any legal liability or responsibility for the accuracy, completeness, or usefulness of any information, apparatus, product, or process disclosed, or represents that its use would not infringe privately owned rights. Reference herein to any specific commercial product, process, or service by trade name, trademark, manufacturer, or otherwise, does not necessarily constitute or imply its endorsement, recommendation, or favoring by the United States Government or any agency thereof. The views and opinions of authors expressed herein do not necessarily state or reflect those of the United States Government or any agency thereof.

# Backup

# Assumptions used in the modeling

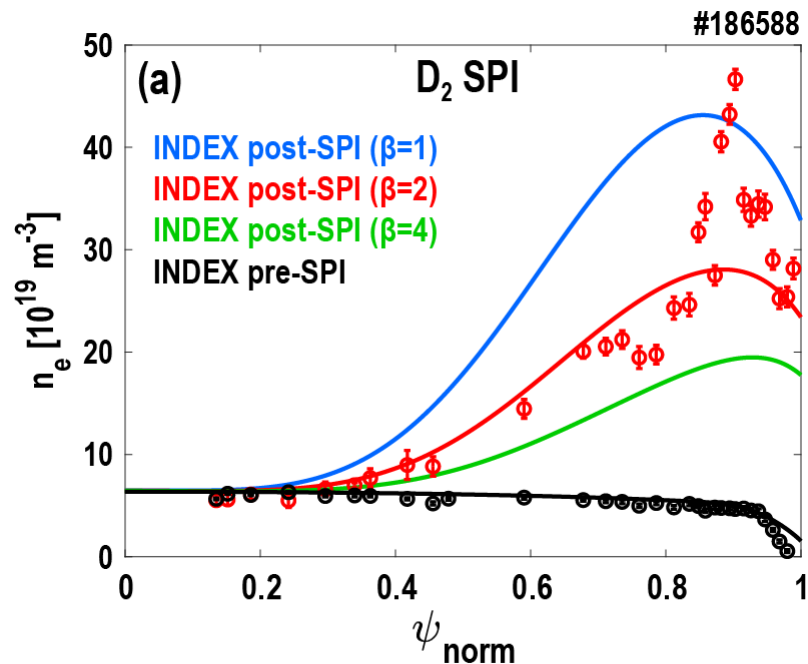
## Back-averaged density increase caused by single shards:

$$\Delta n_{\text{dep}} = \Delta N / [(1 + \beta)V_p(1) - V_p(\rho')]$$

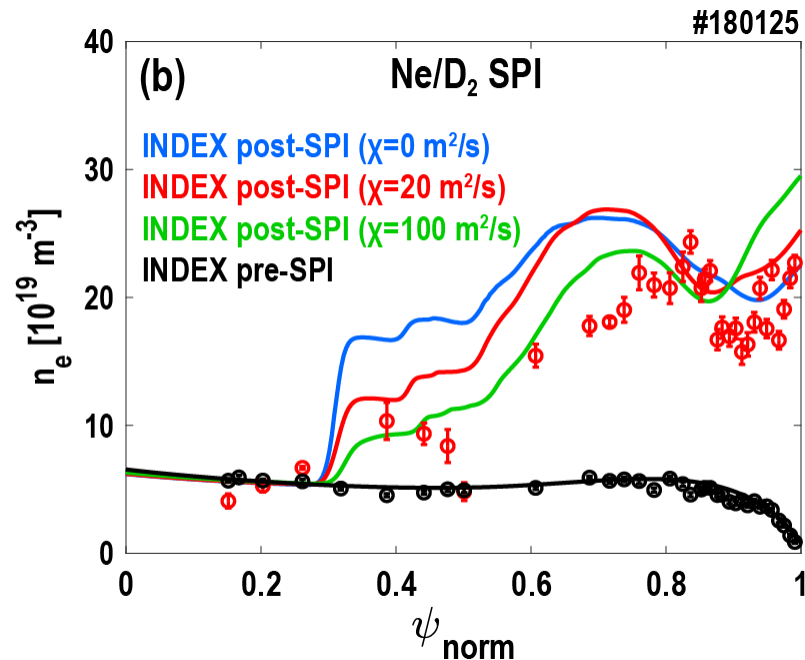
number of atoms ablated from the shard      extent of the radial drift      interior plasma volume with  $\rho < \rho'$

- Shard size distribution is based on the fragmentation model [Parks GA Rep. 2016, Hu NF 2018] with the number of shards chosen based on sensitivity studies
- The average speed is assumed to be 90–100% of the pre-shattered speed, including a small loss of the forward momentum
- The speed dispersion is estimated from the laboratory tests [Gebhart FST 2021] and taken as 50% with the mass dependence of the shard speed included so that only small shards are distributed in the front and rear of the plume
- The fraction of the pellet turned into gas during the shattering process is 30% with only 1/3 of this reaching the plasma, while the remaining 2/3 are subtracted from the injected mass as loss

# Appendix



Scan of factor  $\beta$ , assuming that 50% of injected mass does not interact with plasma



Scan of thermal transport coefficient  $\chi$



# Plasma shrinks during CQ which limits Thomson measurements in the core

